

Application of chiral two- and three-baryon interactions to light hypernuclei

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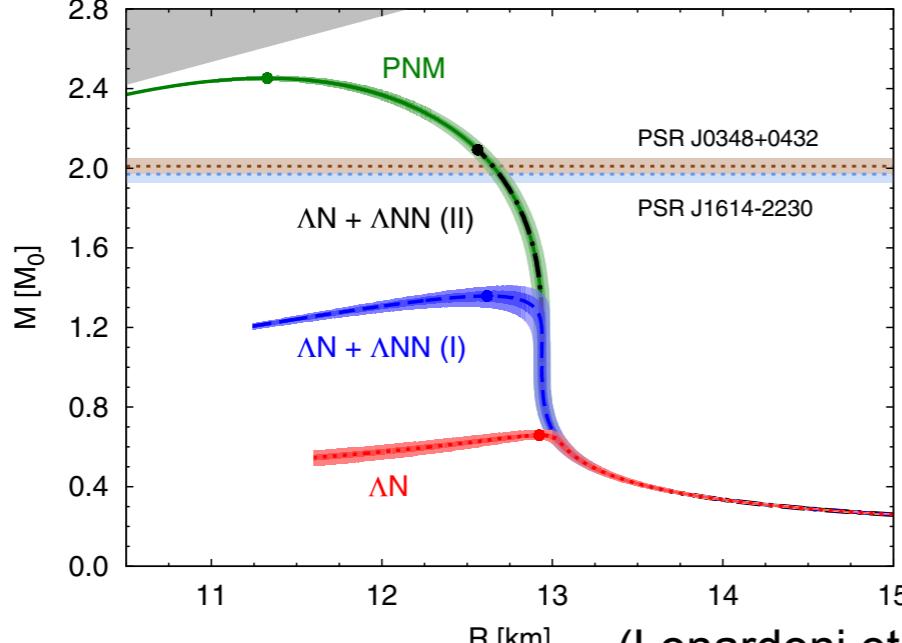
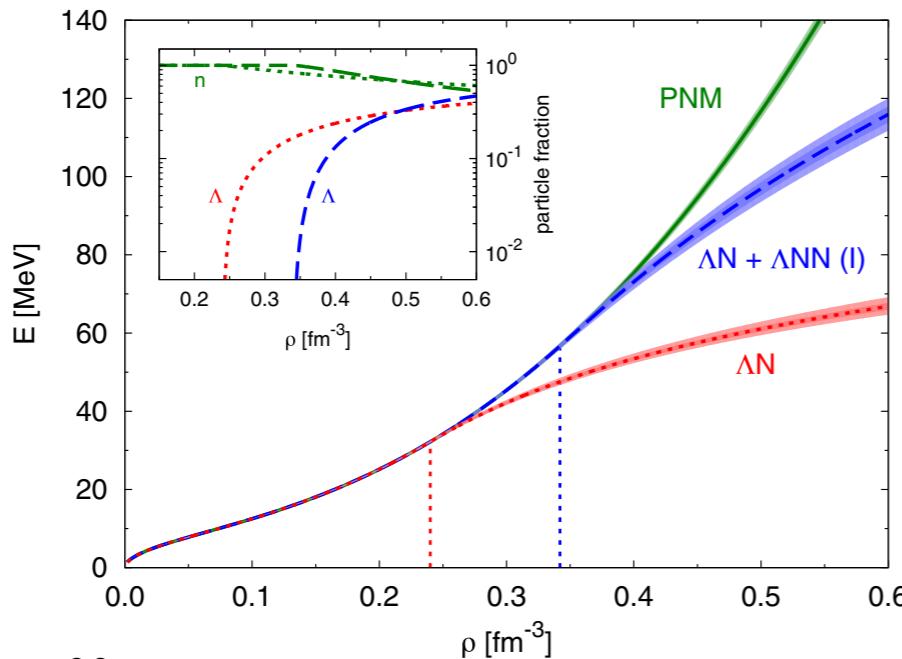
- Motivation
- YN interactions
- SRG evolution of (hyper-)nuclear interactions
- Determination of CSB contact interactions and Λn scattering length
- Application to $A = 7$ and 8 hypernuclei
- Uncertainty of Λ separation energies and size of chiral 3BF contributions
- Chiral YNN interactions
- Conclusions & Outlook

in collaboration with Johann Haidenbauer, Hoai Le, Ulf Meißner

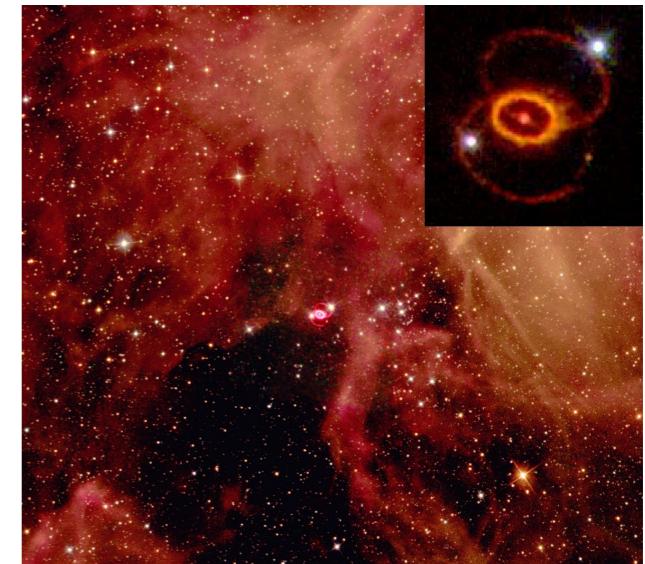
Hypernuclear interactions

Why is understanding hypernuclear interactions interesting?

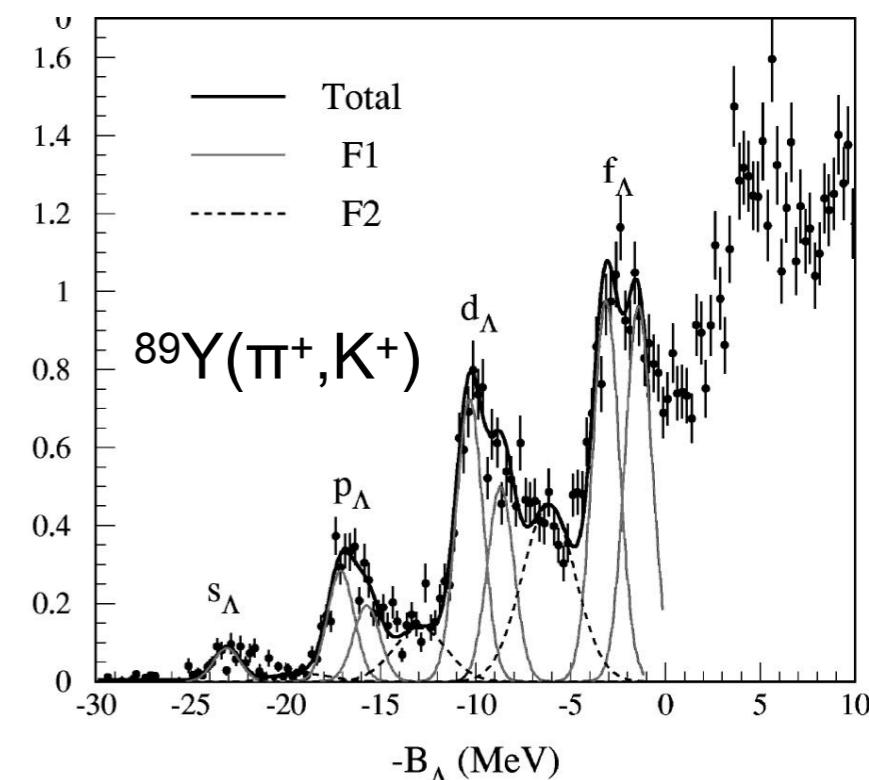
- *hyperon contribution to the EOS, neutron stars, supernovae*
- *"hyperon puzzle"*
- *Λ as probe to nuclear structure*
- *flavor dependence of baryon-baryon interactions*



(Lonardoni et al. (2015))



(SN1987a, Wikipedia)



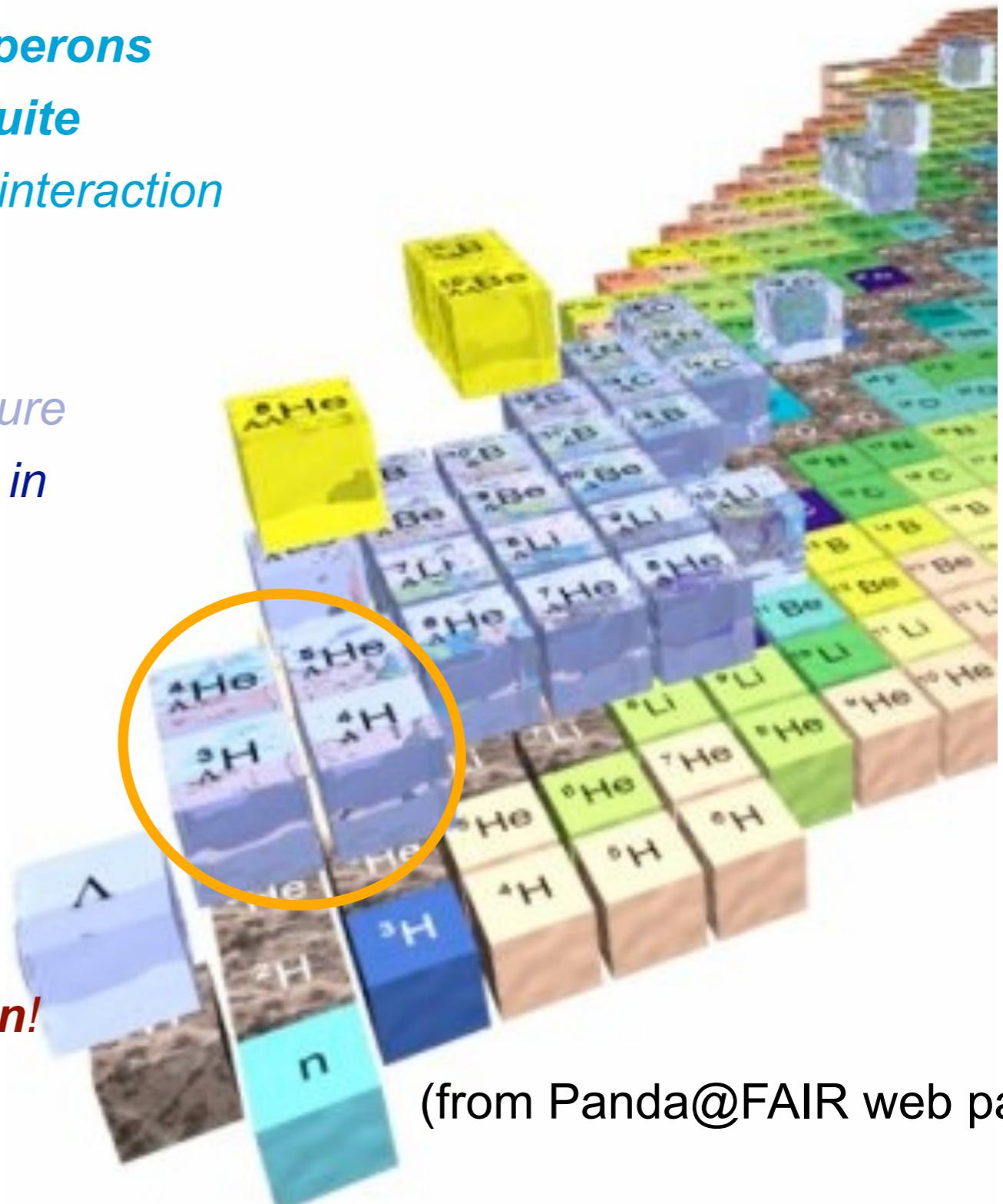
(Hotchi et al. (2001))

Hypernuclei



Only few YN data. Hypernuclear data provides additional constraints.

- ΛN interactions are generally weaker than the NN interaction
 - naively: **core nucleus + hyperons**
 - „separation energies“ are quite independent from $NN(+3N)$ interaction
- no Pauli blocking of Λ in nuclei
 - good to study nuclear structure
 - even light hypernuclei exist in several spin states
- non-trivial constraints on the YN interaction even from lightest ones
- size of YNN interactions?
need to include Λ - Σ conversion!



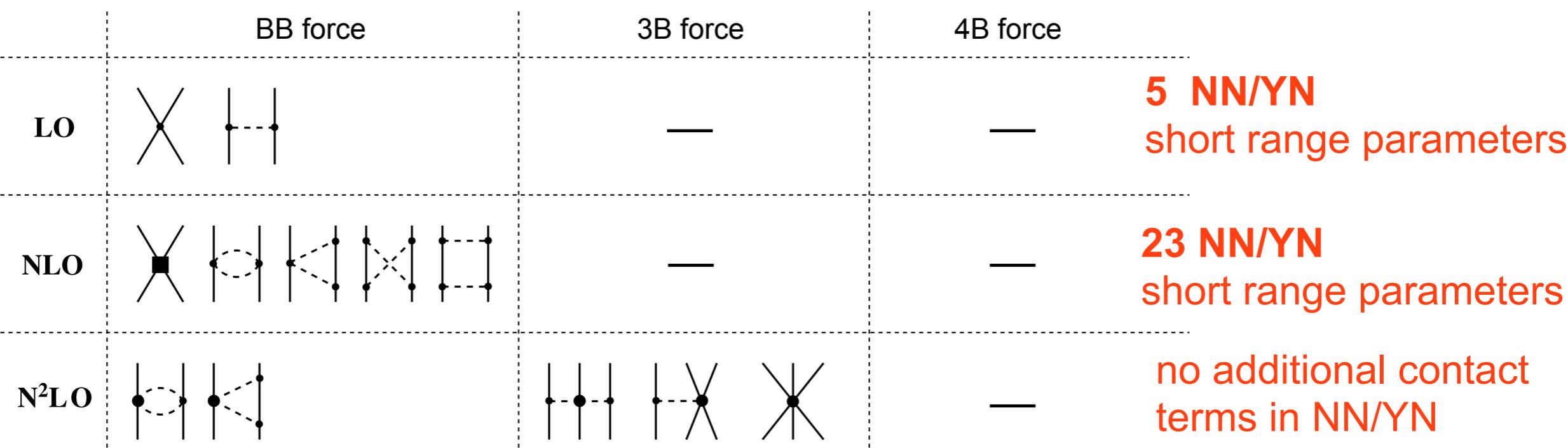
Chiral NN & YN interactions



EFT based approaches

Chiral EFT implements **chiral symmetry of QCD**

- symmetries constrain exchanges of Goldstone bosons
- relations of two- and three- and more-baryon interactions
- breakdown scale $\approx 600 - 700 \text{ MeV}$
- Semi-local momentum regularization (SMS) up to N²LO (for YN)



(adapted from Epelbaum, 2008)

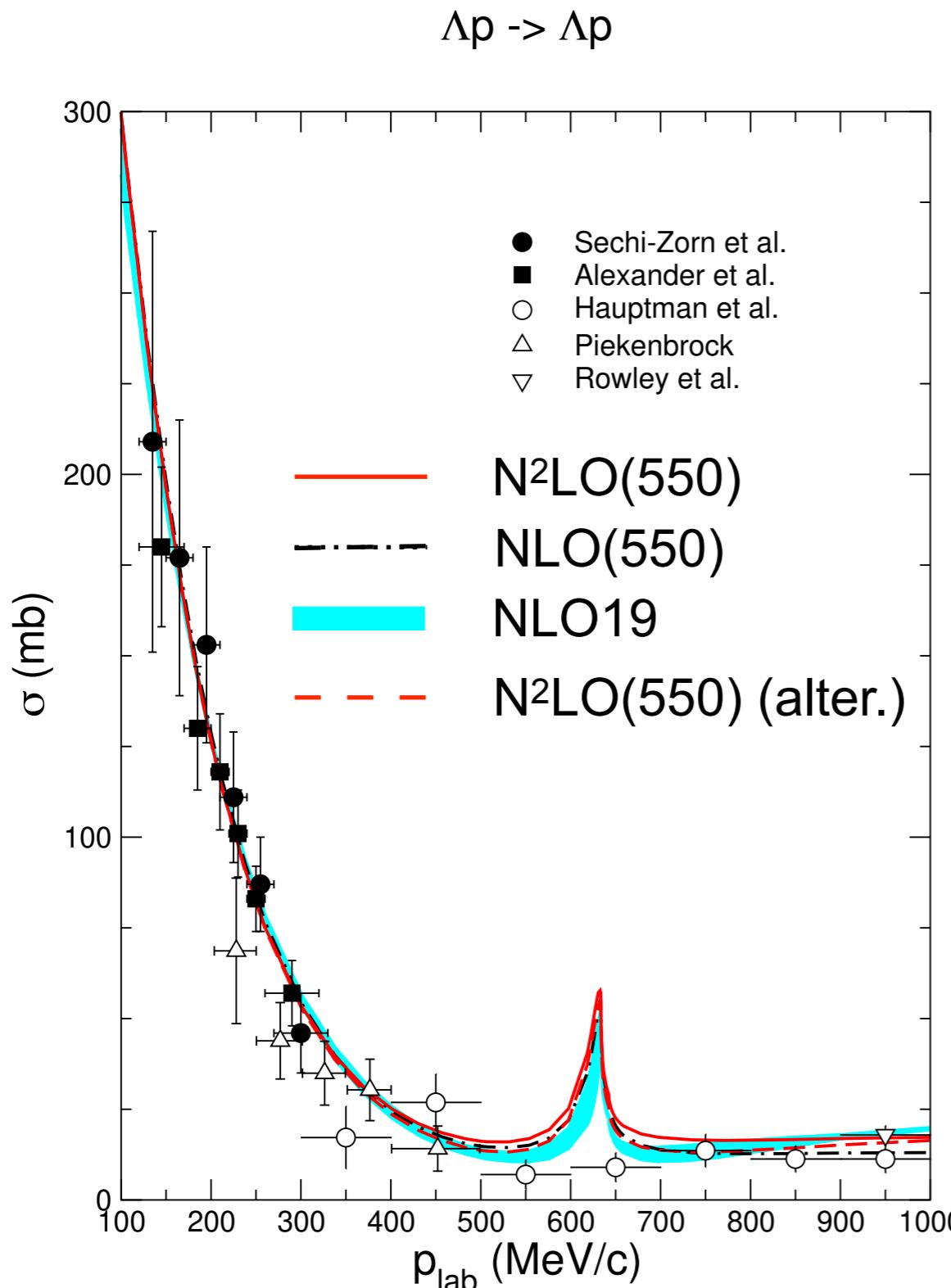
Retain flexibility to adjust to data due to counter terms

Regulator required — cutoff/different orders often used to estimate uncertainty

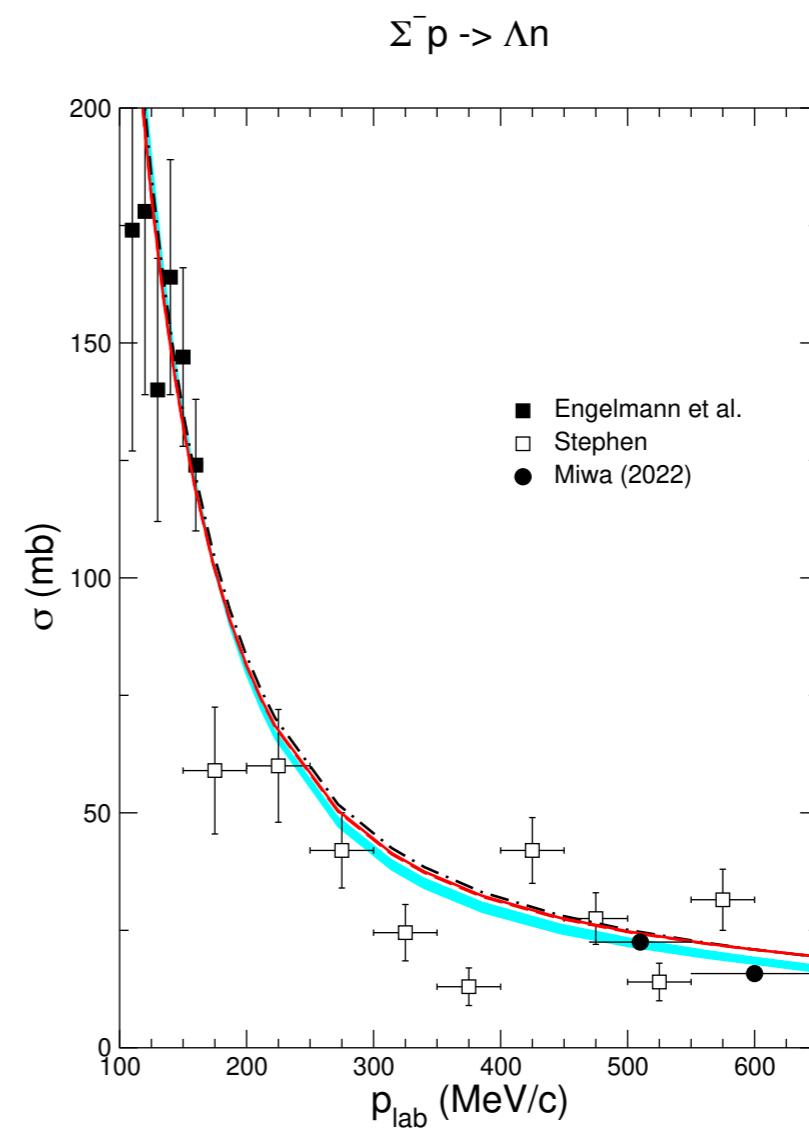
$\Lambda - \Sigma$ conversion is explicitly included (3BFs only in N²LO)

SMS NLO/N²LO interaction

Selected results (show $\Lambda = 550$ MeV, others are very similar in quality)



- most relevant cross sections very similar in NLO and N²LO
- similar to NLO19
- alternative fit (see later)

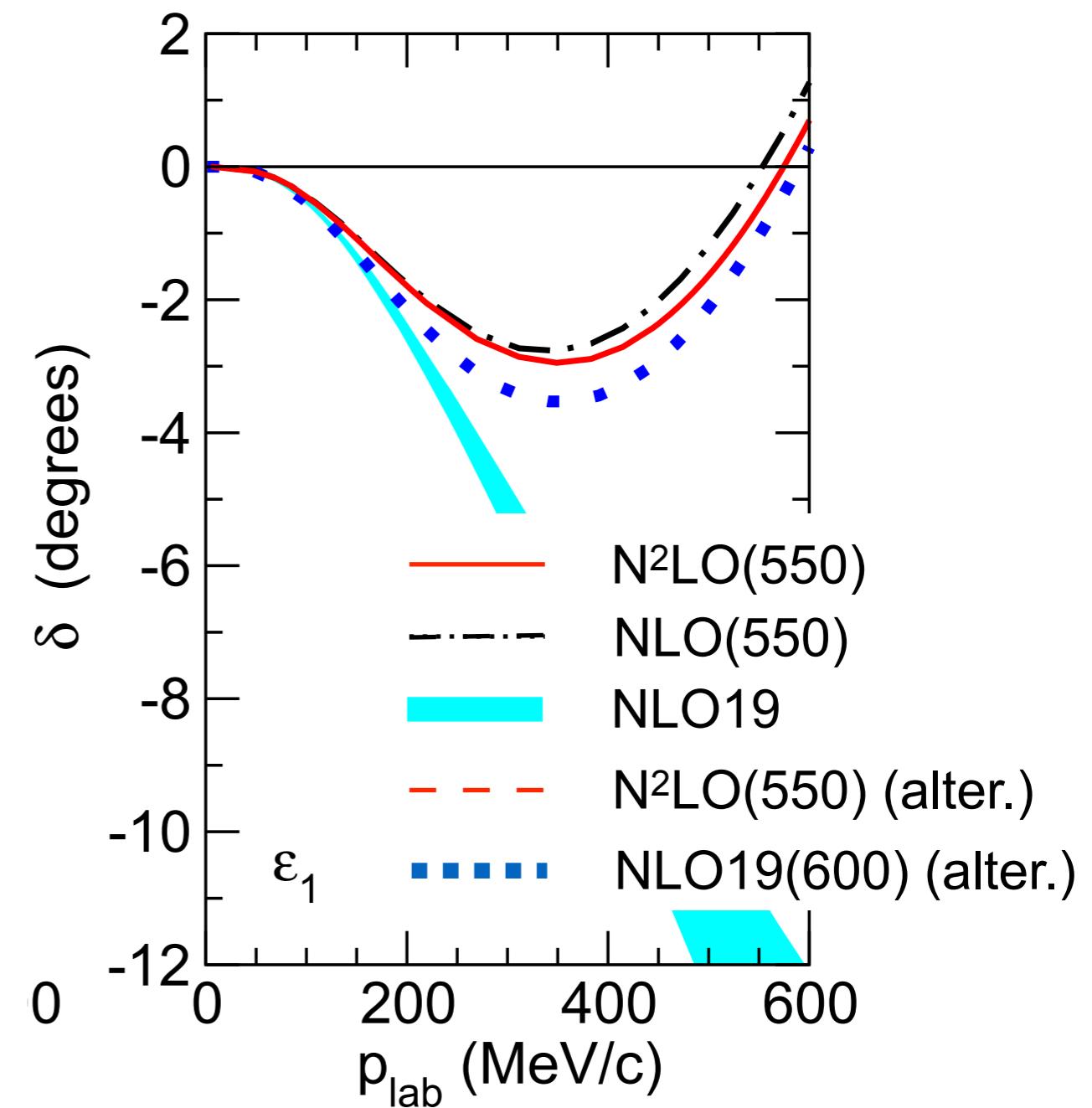
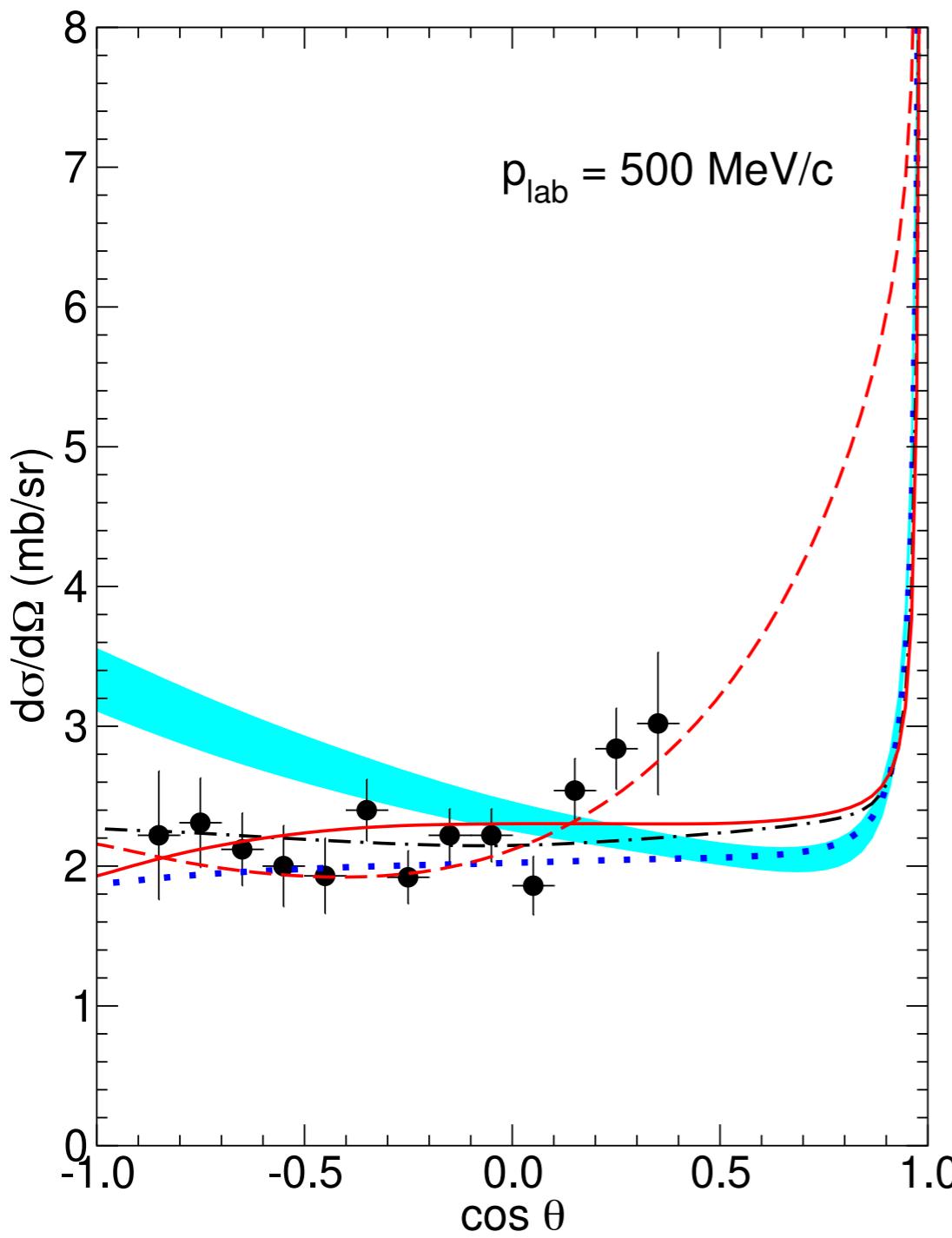


SMS NLO/N²LO interaction



new data (Miwa(2022)) at higher energies provides new constraints!

$$\Sigma^+ p \rightarrow \Sigma^+ p$$



J. Haidenbauer et al. EPJ A 59, 63 (2023).

Similarity renormalization group is by now a standard tool to obtain soft effective interactions for various many-body approaches (NCSM, coupled-cluster, MBPT, ...)

Idea: perform a unitary transformation of the NN (and YN interaction) using a cleverly defined "generator"

$$\frac{dH_s}{ds} = [\underbrace{[T, H(s)]}_{\equiv \eta(s)}, H(s)] \quad H(s) = T + V(s)$$

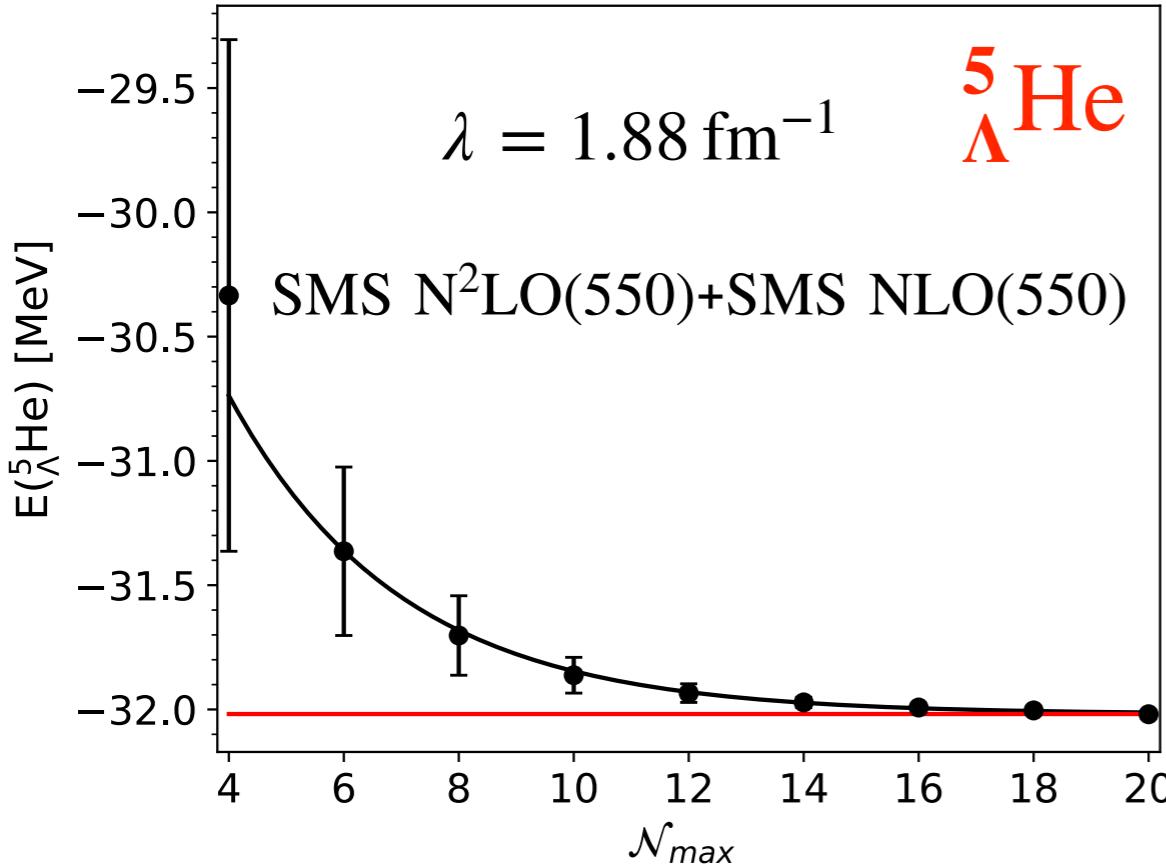
this choice of generator drives $V(s)$ into a diagonal form in momentum space

- $V(s)$ will be **phase equivalent** to original interaction
- short range $V(s)$ will change towards **softer interactions**
- Evolution can be restricted to **2-,3-, ... body level** (approximation)
- $\lambda = \left(\frac{4\mu_{BN}^2}{s} \right)^{1/4}$ is a measure of the width of the interaction in momentum space
- **dependence** of results on λ or s is a measure for **missing terms**

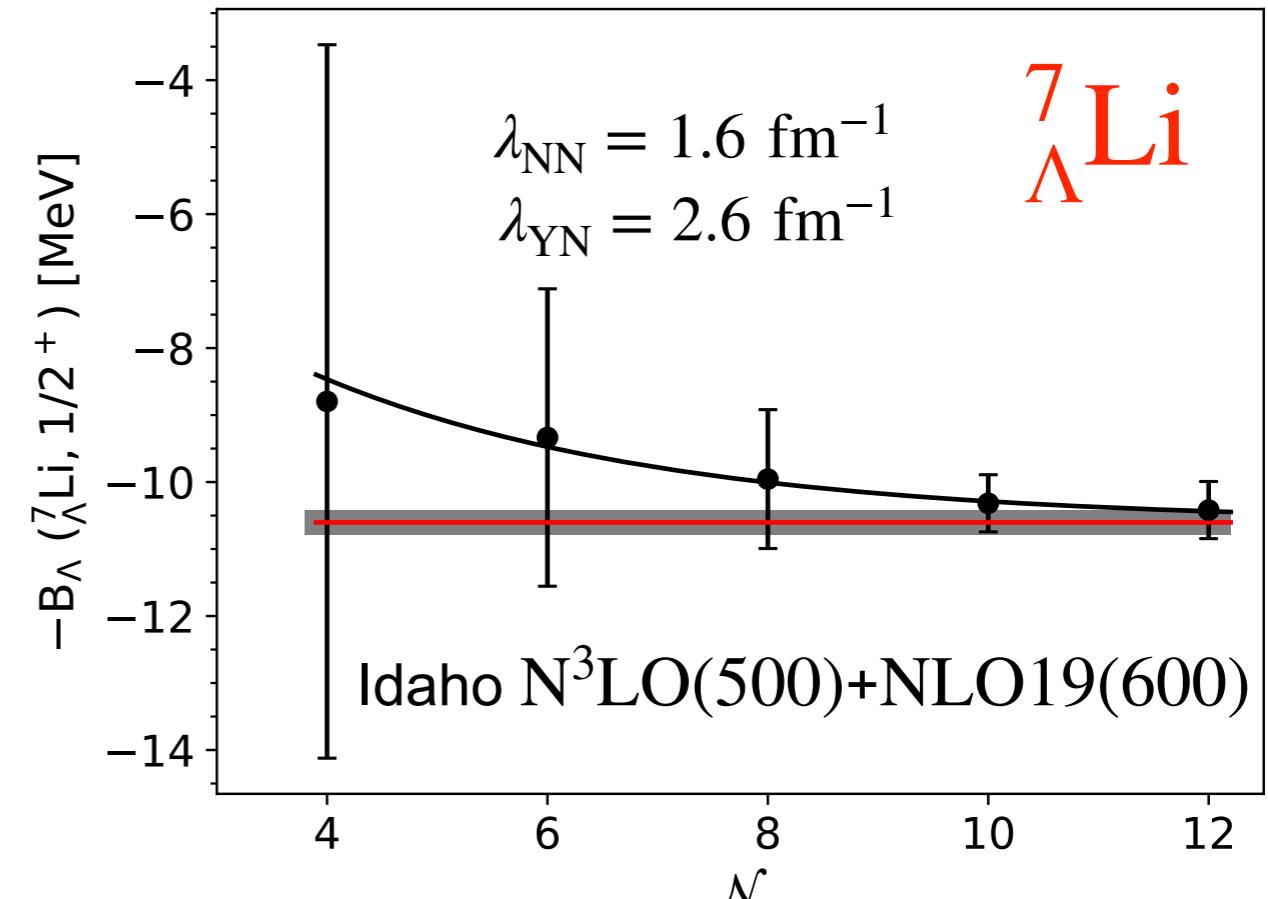
(Bogner et al., 2007)

J-NCSM convergence

SRG evolution improves convergence



$$E(^5_\Lambda \text{He}) = -32.018 \pm 0.001 \text{ MeV}$$



$$E_\Lambda(^7_\Lambda \text{Li}) = 10.6 \pm 0.2 \text{ MeV}$$

- for light nuclei and hypernuclei, the numerical uncertainty is negligible.
- for p-shell nuclei/hypernuclei, the uncertainty is visible
- extrapolation of separation energy can reduce uncertainty of this quantity

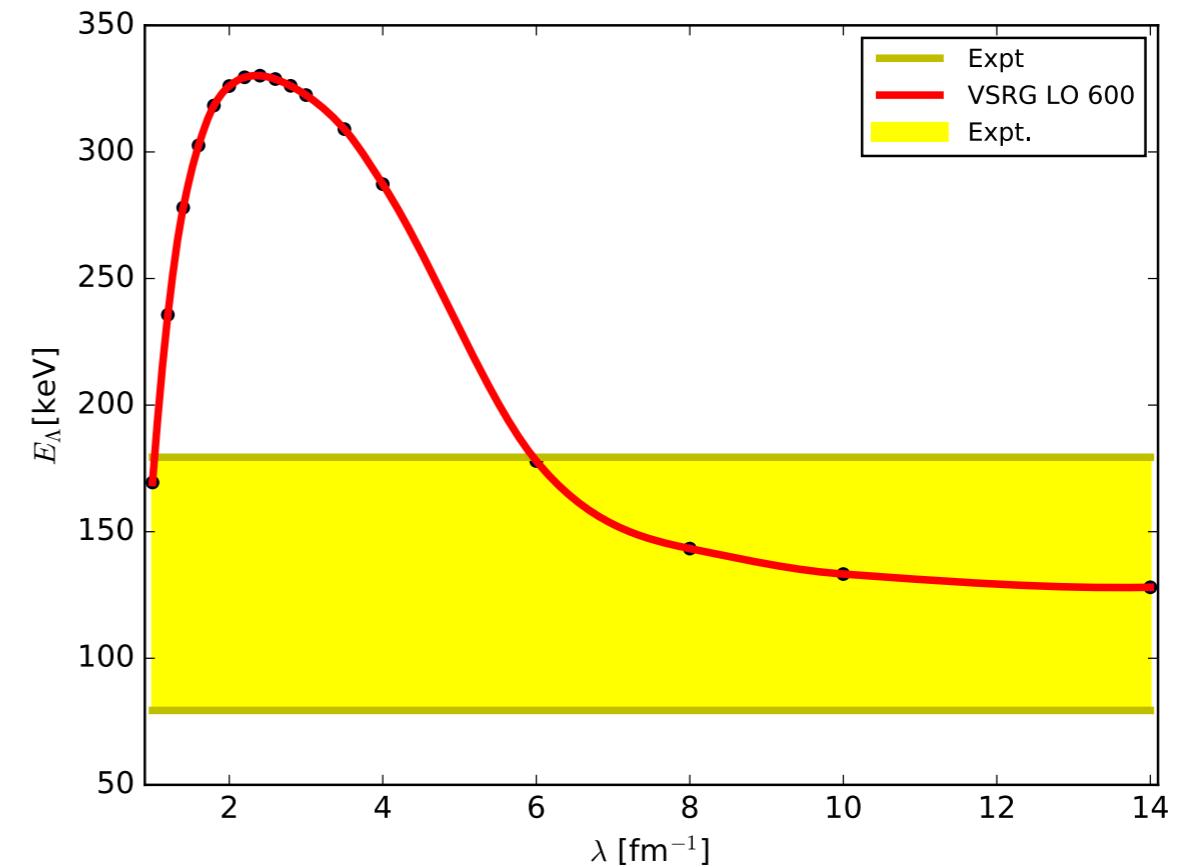
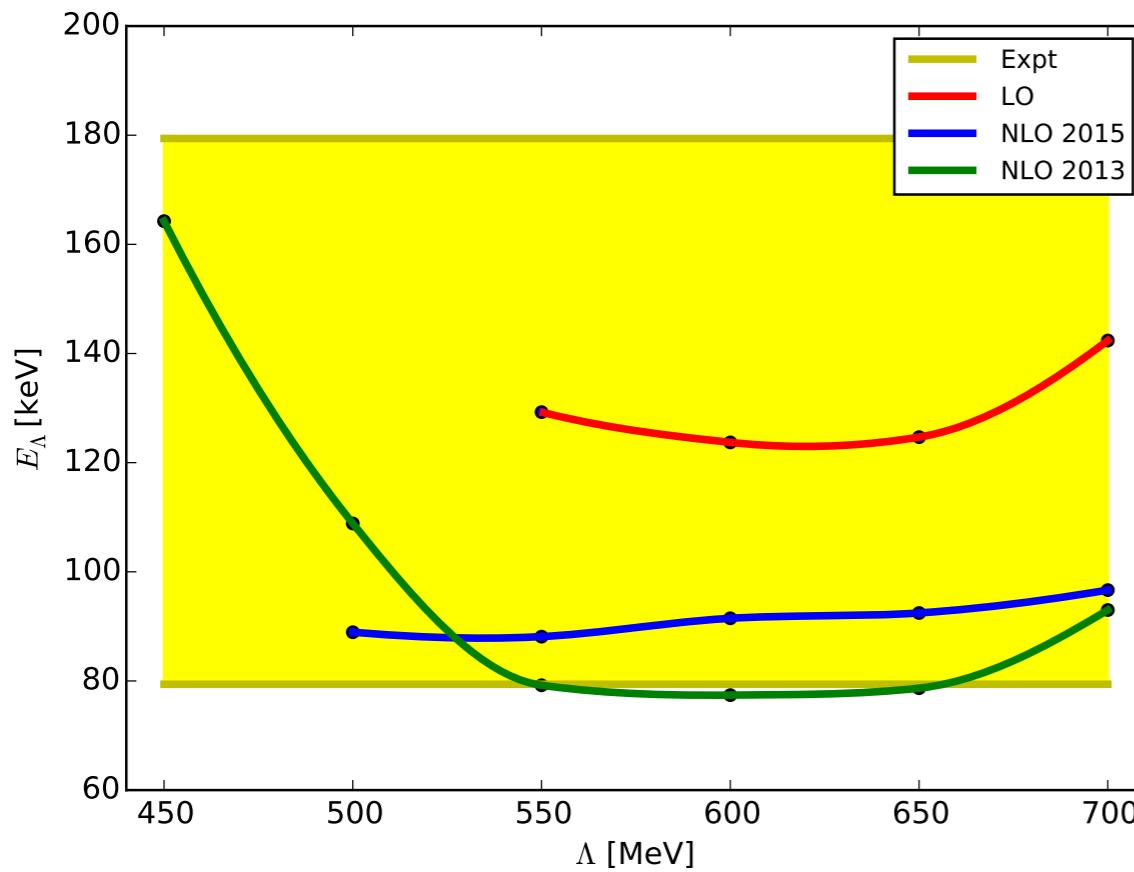
Induced 3BF ...



SRG parameter dependence is significant when NN and YN interactions are evolved

→ missing 3N and YNN interactions

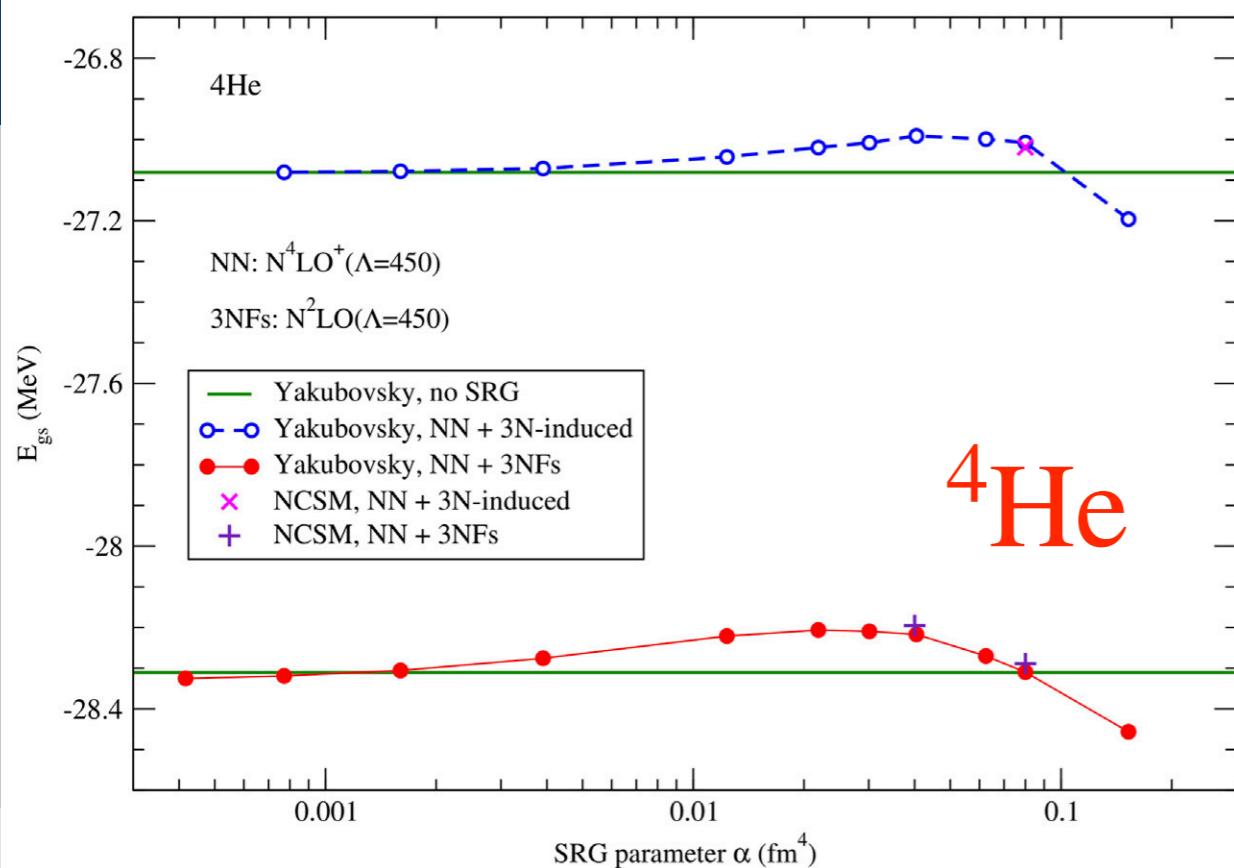
- 3NF is comparable to chiral 3NF
- YNN is larger than chiral YNN



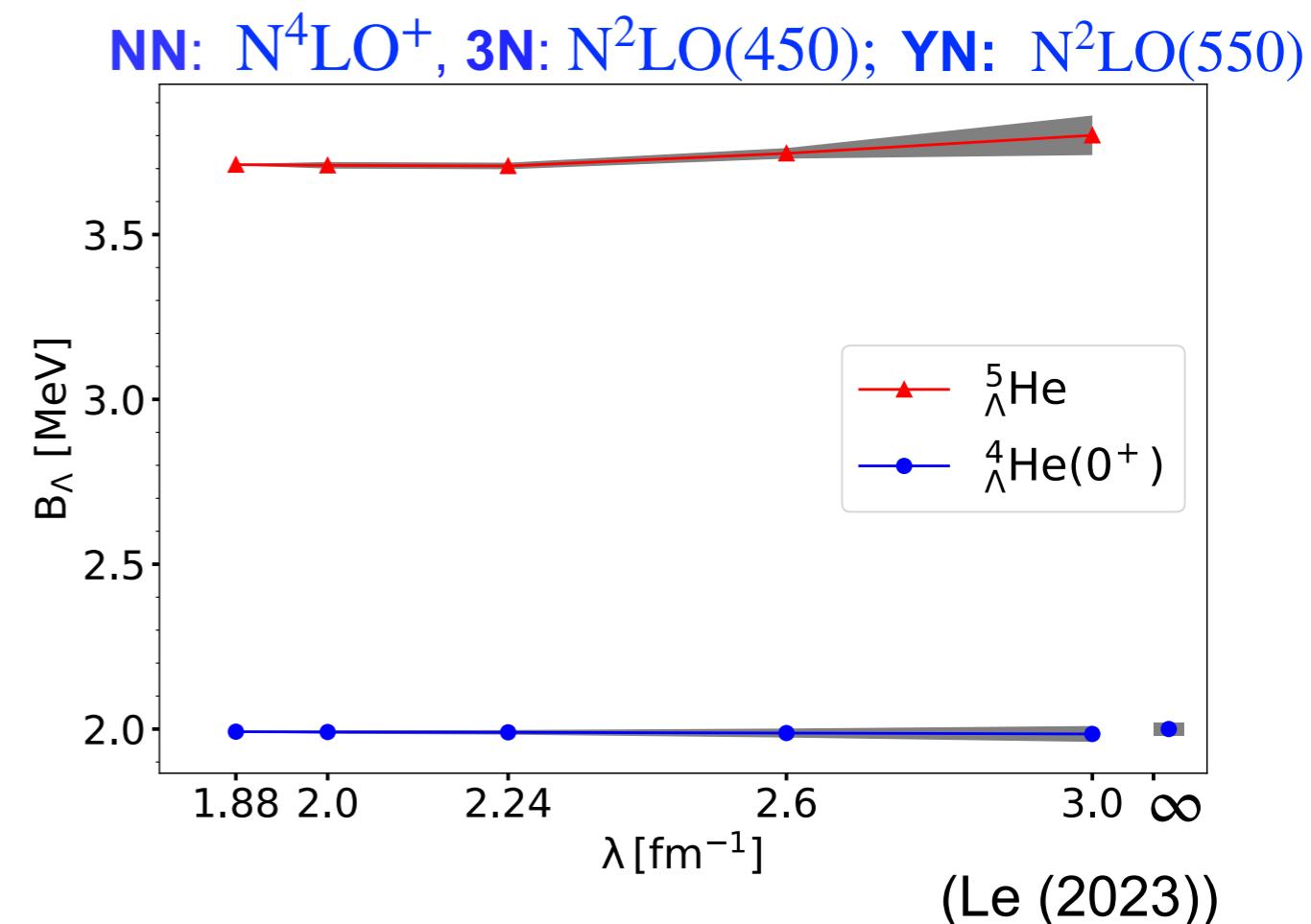
SRG dependence of results



- SRG-induced 3N and YNN interactions
- ^4He binding energies varies by $\approx 100 - 200$ keV (relevant in the future?)
- separation energies are even less dependent (YNNN forces small)



(Maris, Le, Nogga, Roth, Vary (2023))



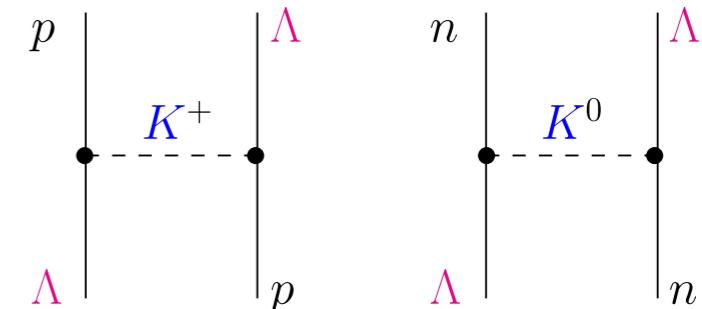
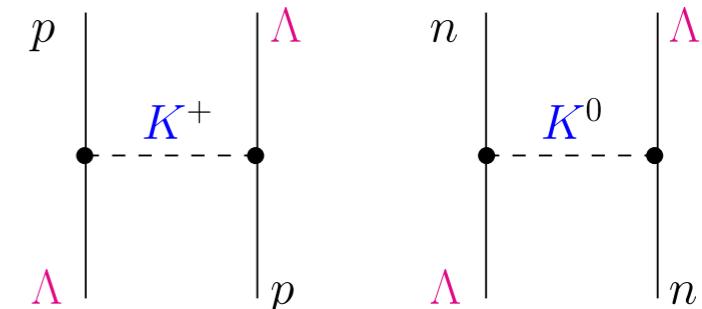
(Le (2023))

For **hypernuclei**, calculations based on SRG induced BB and 3B interactions are sufficiently accurate!

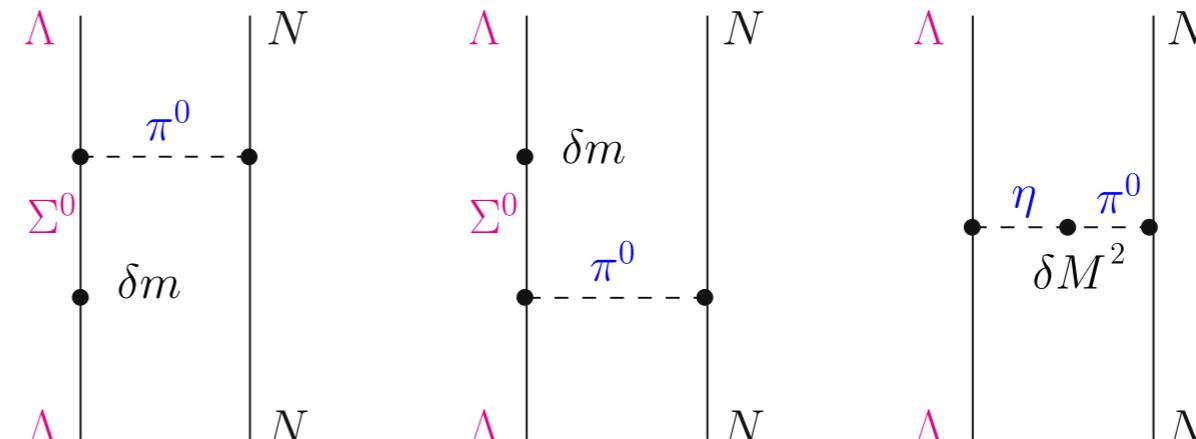
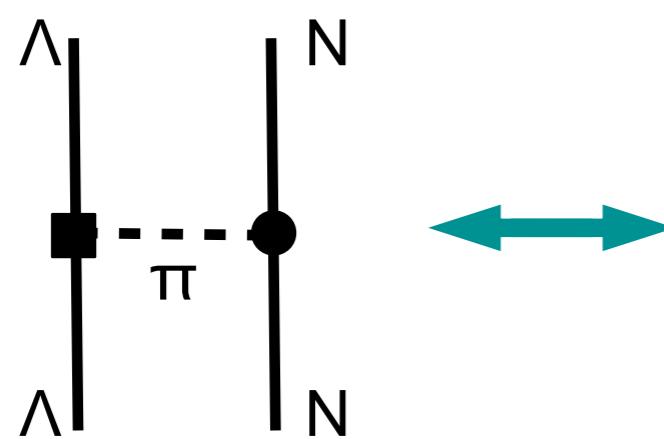
CSB contributions to YN interactions



- **formally leading** contributions:
Goldstone boson mass difference
 - very small due to the small relative difference of kaon masses
- **subleading but most important**
 - effective CSB $\Lambda\Lambda\pi$ coupling constant (Dalitz, van Hippel, 1964)

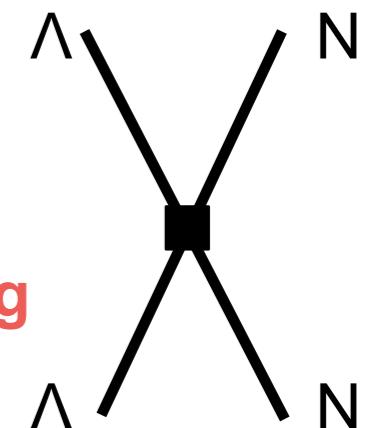


$$f_{\Lambda\Lambda\pi} = \left[-2 \frac{\langle \Sigma^0 | \delta m | \Lambda \rangle}{m_{\Sigma^0} - m_\Lambda} + \frac{\langle \pi^0 | \delta M^2 | \eta \rangle}{M_\eta^2 - M_{\pi^0}^2} \right] f_{\Lambda\Sigma\pi} \approx (-0.0297 - 0.0106) f_{\Lambda\Sigma\pi}$$

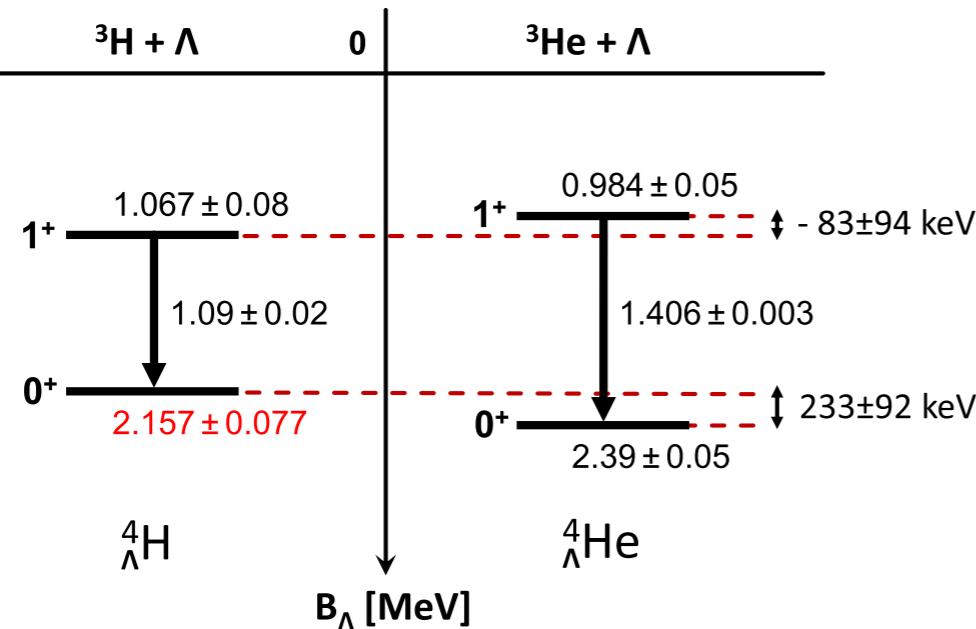


- **so far less considered but necessary for proper renormalization**
 - CSB contact interactions (for singlet and triplet)

Aim: determine the two unknown CSB LECs and predict Λn scattering



Fit of contact interactions



(Schulz et al., 2016; Yamamoto, 2015)

- Adjust the two CSB contact interactions to one main scenario (**CSB1**)

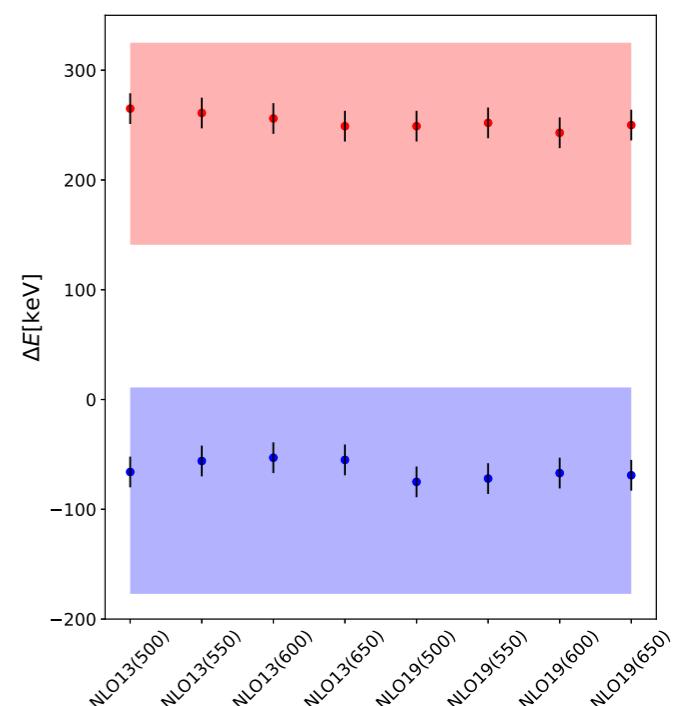
Λ	NLO13		NLO19	
	C_s^{CSB}	C_t^{CSB}	C_s^{CSB}	C_t^{CSB}
500	4.691×10^{-3}	-9.294×10^{-4}	5.590×10^{-3}	-9.505×10^{-4}
550	6.724×10^{-3}	-8.625×10^{-4}	6.863×10^{-3}	-1.260×10^{-3}
600	9.960×10^{-3}	-9.870×10^{-4}	9.217×10^{-3}	-1.305×10^{-3}
650	1.500×10^{-2}	-1.142×10^{-3}	1.240×10^{-2}	-1.395×10^{-3}

The values of the LECs are in 10^4 GeV^{-2}

- Size of LECs as expected by power counting

$$\frac{m_d - m_u}{m_u + m_d} \left(\frac{M_\pi}{\Lambda} \right)^2 C_{S,T} \approx 0.3 \cdot 0.04 \cdot 0.5 \cdot 10^4 \text{ GeV}^{-2} \propto 6 \cdot 10^{-3} \cdot 10^4 \text{ GeV}^{-2}$$

- Problem: large experimental uncertainty of experiment
- here only **fit to central values** to test theoretical uncertainties

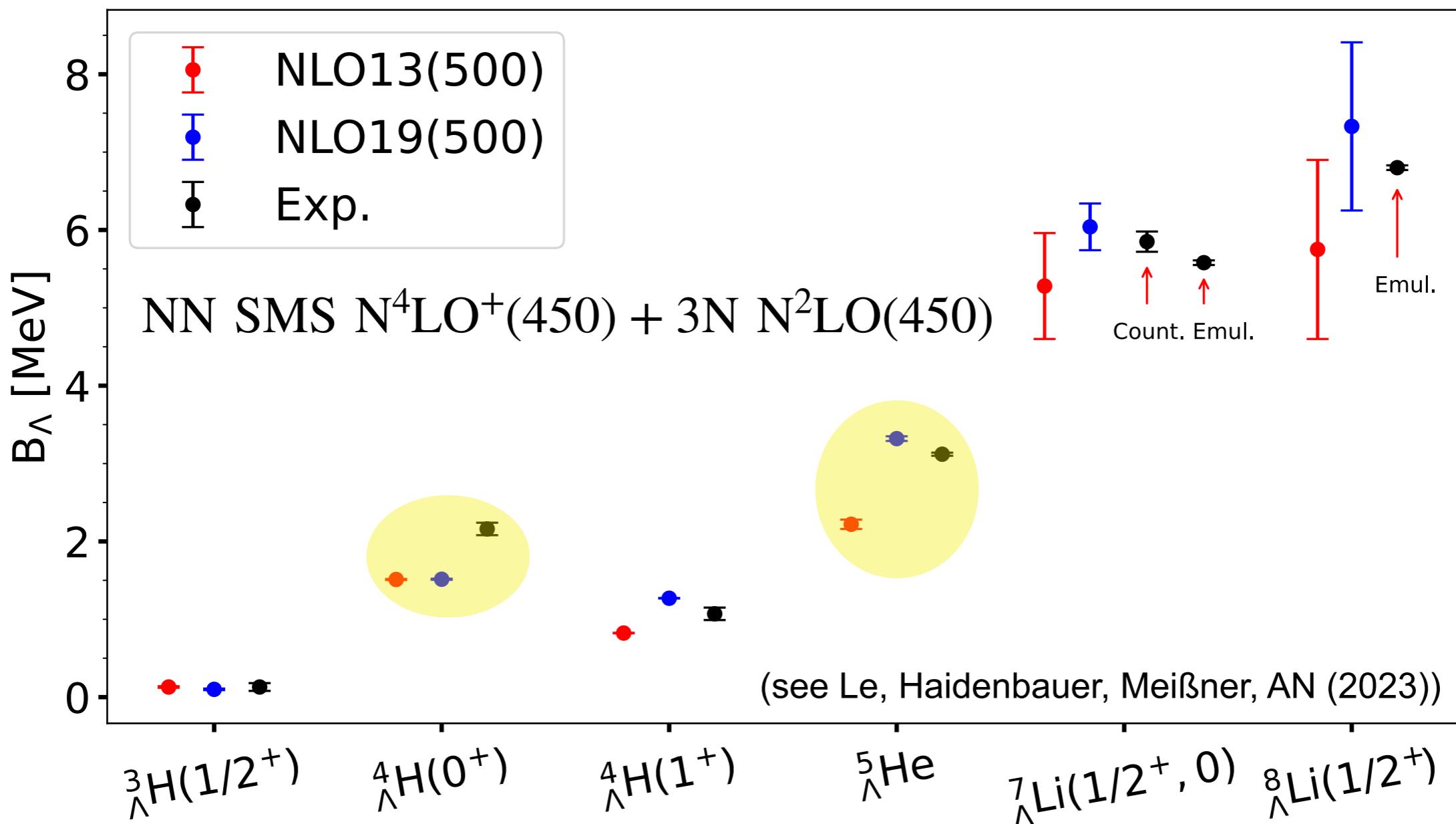


(see Haidenbauer, Meißner, AN (2021))

Application to $A = 7$ and 8



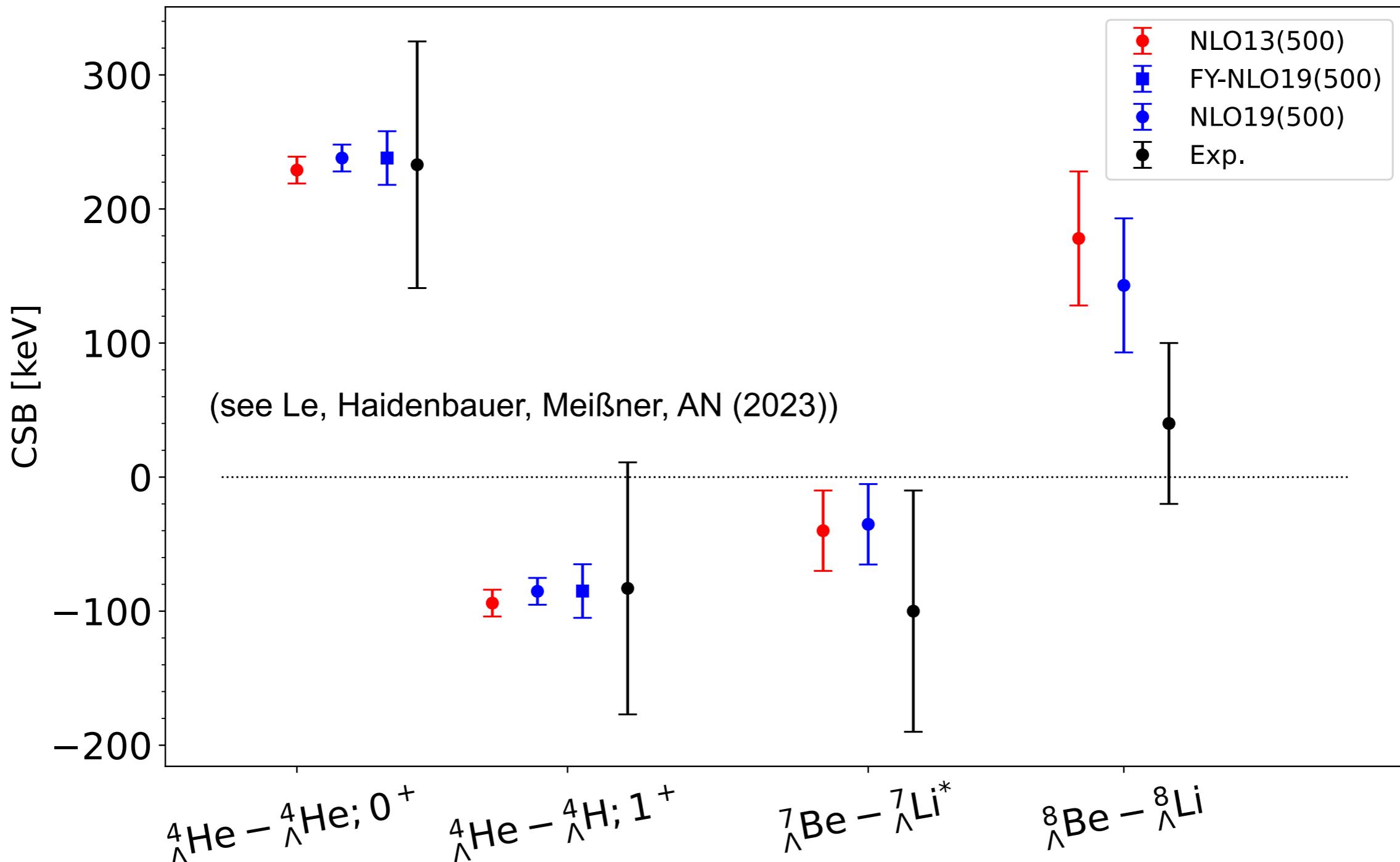
- YN interaction adjusted to the hypertriton — YNN is small
- based only on YN interactions: splitting for ${}^4_{\Lambda}\text{H}$ is not well reproduced — YNN(?)
- NLO19 gives better results for ${}^5_{\Lambda}\text{He}$ and heavier hypernuclei
— accidentally small YNN interaction?
- uncertainties are numerical — no estimate of chiral uncertainties yet



Application to $A = 7$ and 8



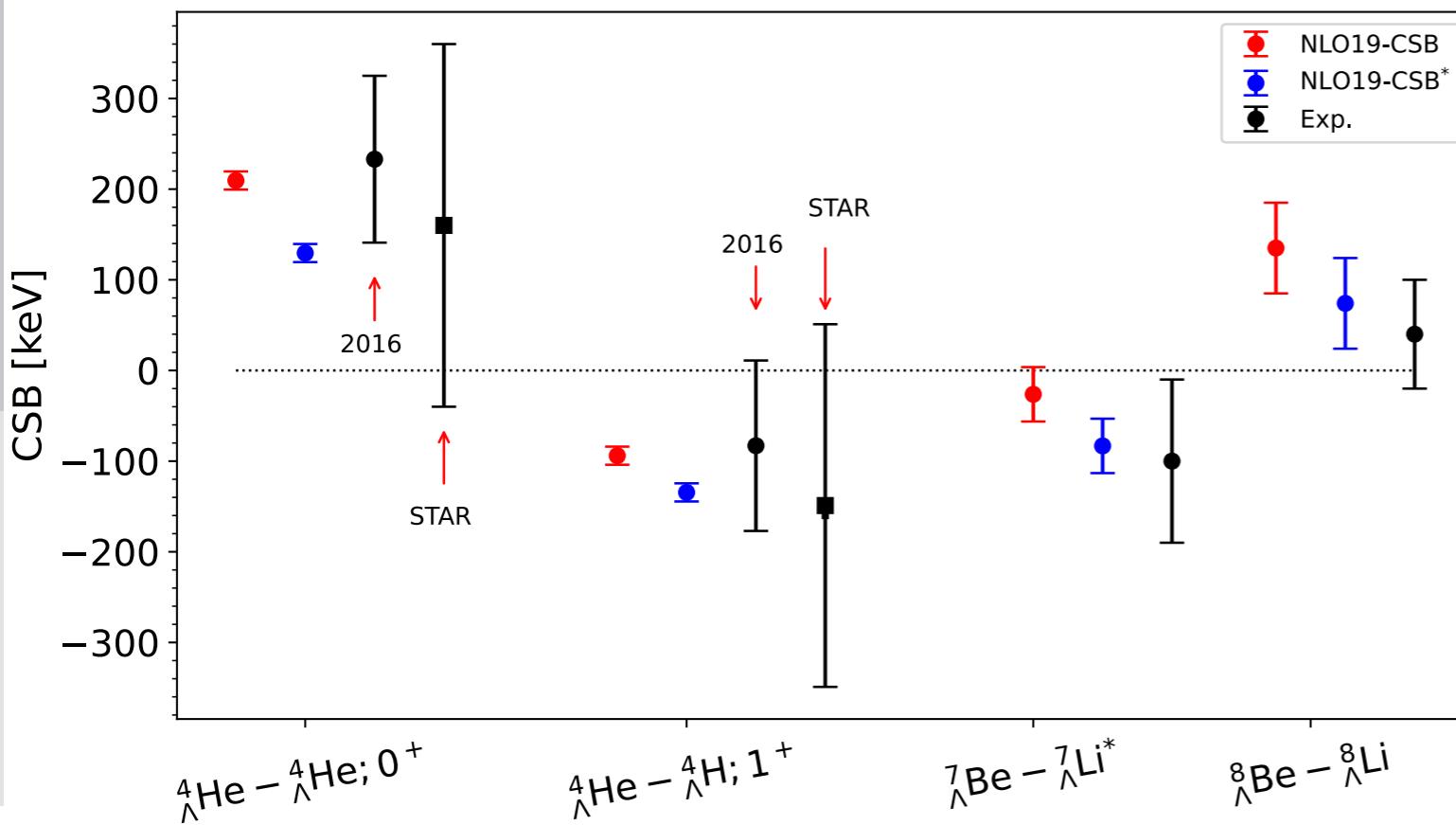
- CSB of singlet and triplet states interferes differently
- CSB still not fixed — experimental uncertainty is large
- scenario studied here is only **marginally consistent** with CSB in $A = 8$



New STAR data for $A = 4$ CSB



- fit to STAR data only
- only slight adjustment required
- improves description to p-shell CSB
- higher experimental accuracy is desirable
- good example of using hypernuclei to determine YN interactions



	NLO19(500)	CSB	CSB*
a_s^{Ap}	-2.91	-2.65	-2.58
a_s^{An}	-2.91	-3.20	-3.29
δa_s	0	0.55	0.71
a_t^{Ap}	-1.42	-1.57	-1.52
a_t^{An}	-1.41	-1.45	-1.49
δa_t	-0.01	-0.12	-0.03

(see Le, Haidenbauer, Meißner, AN (2023))

Uncertainty analysis to $A = 3$ to 5

Order N²LO requires combination of chiral NN, YN, 3N and YNN interaction

Results for **different orders** enable uncertainty estimate:

Ansatz for the order by order convergence:

$$X_K = X_{ref} \sum_{k=0}^K c_k Q^k \quad \text{where } Q = M_\pi^{eff}/\Lambda_b \quad (X_{ref} \text{ LO, exp., max, ...})$$

Bayesian analysis of the uncertainty following Melendez et al. 2017,2019

Extracting c_k for $k \leq K$ from calculations

→ probability distributions for c_k

$$\delta X_K = X_{ref} \sum_{k=K+1}^{\infty} c_k Q^k$$

Uncertainty due to missing higher orders is more relevant

than numerical uncertainty! (for light nuclei)

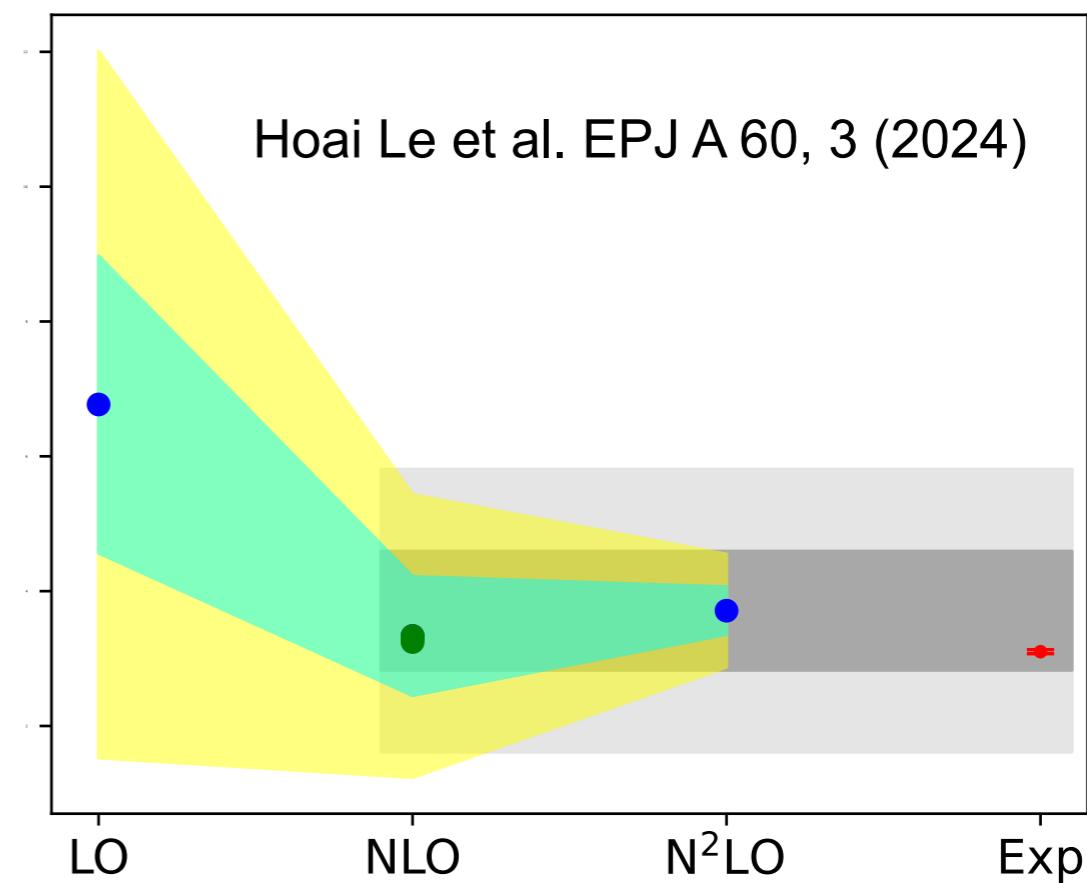
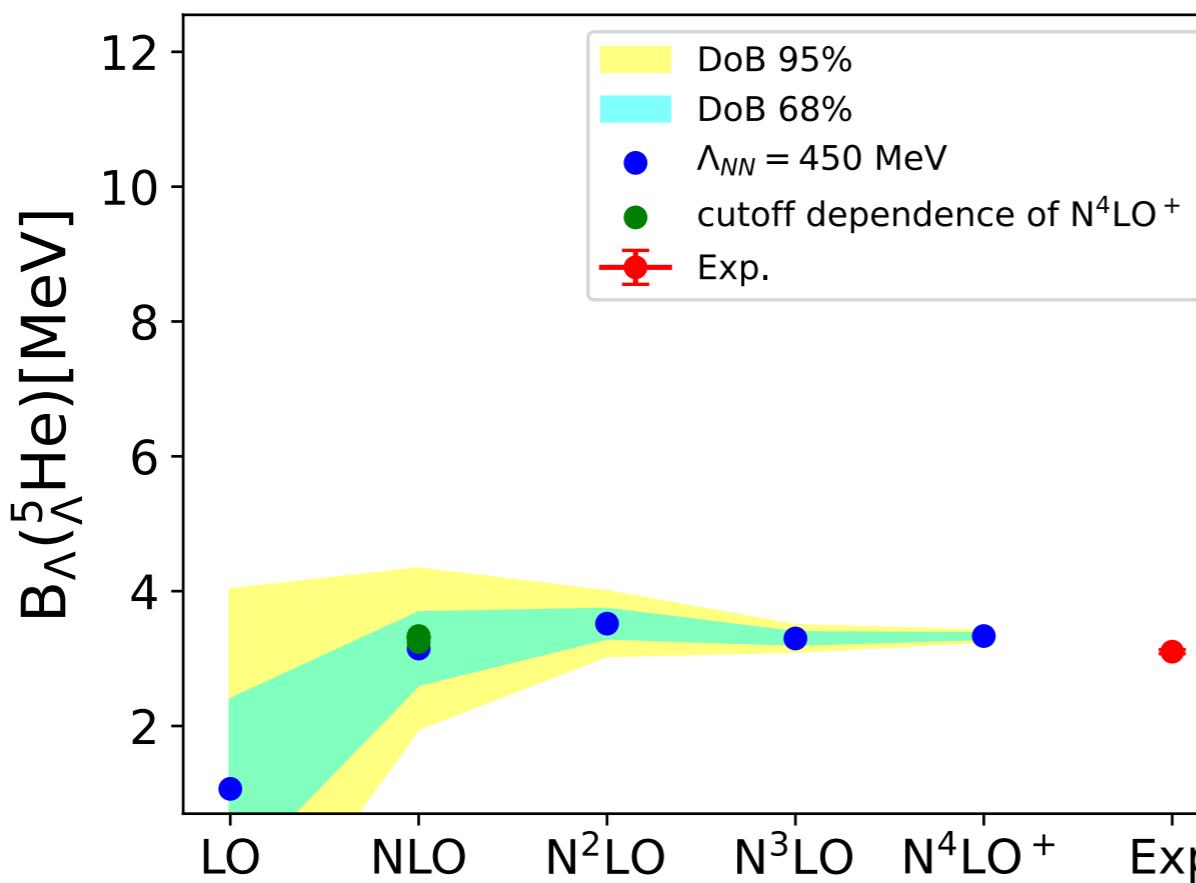
Application to $^5_{\Lambda}\text{He}$ and summary



- without YNN: sizable uncertainties at $A = 4$ and 5
- $A = 3$ sufficiently accurate
- NN/YN dependence small at least for $A = 3$

nucleus	$\Delta_{68}(NN)$	$\Delta_{68}(YN)$
$^3_{\Lambda}\text{H}$	0.011	0.015
$^4_{\Lambda}\text{He} (0^+)$	0.157	0.239
$^4_{\Lambda}\text{He} (1^+)$	0.114	0.214
$^5_{\Lambda}\text{He}$	0.529	0.881

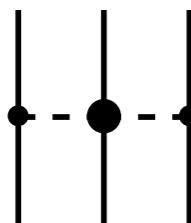
→ at the same time: estimate of YNN !



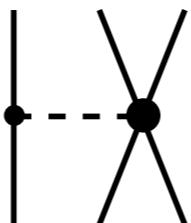
YNN (Λ NN) interactions

Leading 3BF with the usual topologies (see Petschauer et al., 2016 & 2017)

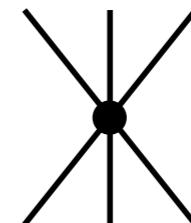
ChPT  all octet mesons contribute  **only take π** explicitly into account



2 LECs in Λ NN
(up to 10)



2 LECs in Λ NN
(up to 14)



3 LECs in Λ NN
5 LECs in Σ NN + 1 Λ - Σ transition

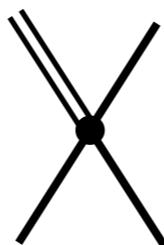
only few data  need to keep the **# of LECs** small

Decuplet baryons (Σ^* ...) might enhance YNN partly to NLO (see Petschauer et al., 2017)

By decuplet saturation all LECs can be related to the following leading octet-decuplet transitions (Petschauer et al. , 2020)



$$\propto C = \frac{3}{4}g_A$$

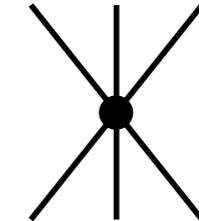
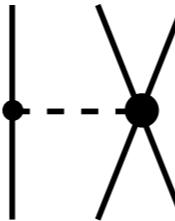
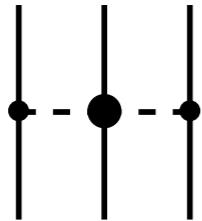


$$\propto G_1, G_2$$

 reduction to 2 LECs

YNN (Λ NN) interactions

Decuplet saturation relates all LECs to G_1 and G_2



$$\propto C^2$$

For Λ NN: $\propto C^2$

$$\propto CG_1, CG_2$$

$$\propto C(G_1 + 3G_2)$$

$$\propto (G_1)^2, (G_2)^2, G_1 G_2$$

$$\propto (G_1 + 3G_2)^2 \quad 1 \text{ LEC}$$

→ density dependent BB interactions (Petschauer et al., 2017)

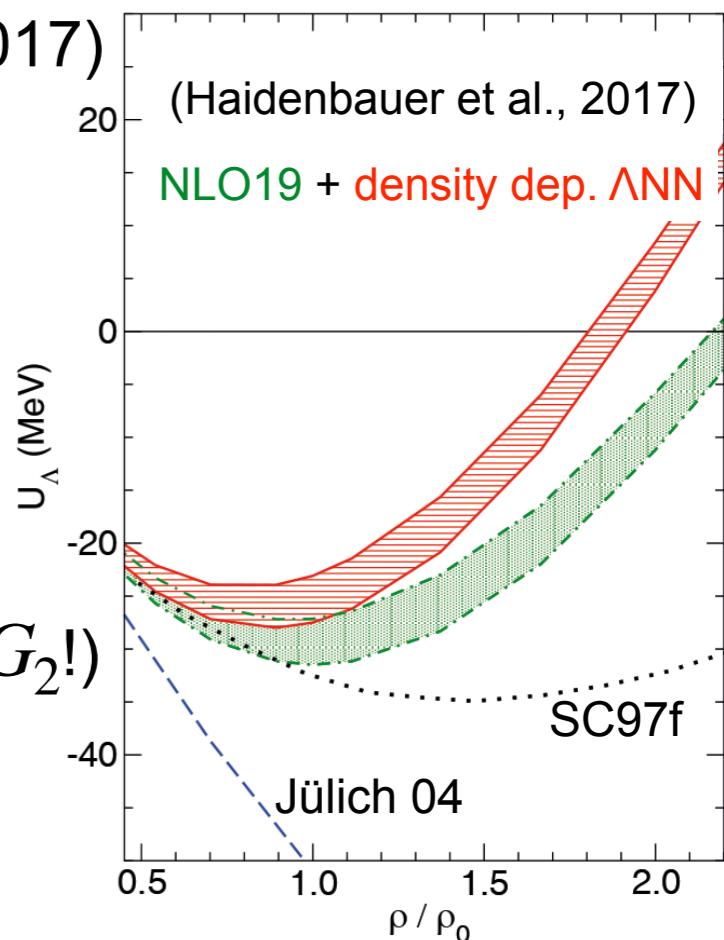
→ application to nuclear matter (Haidenbauer et al., 2017)

neutron stars (Logoteta et al., 2019)

- contribution on the single particle potentials can be large
- realistic results seem to require partly cancellations of 2π and 1π exchange (fixes sign of $G_1 + 3G_2$!)

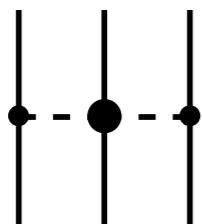
Recently: successful benchmark of matrix elements:

Hoai Le et al. [arXiv:2407.02064v1](https://arxiv.org/abs/2407.02064v1)

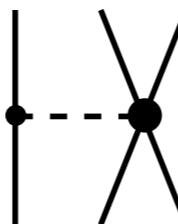


YNN (Λ NN) interactions in practice

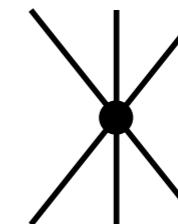
Decuplet approximation in YNN



$$\propto C^2$$



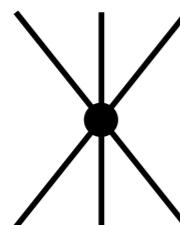
$$\propto CG_1, CG_2$$



$$\propto (G_1)^2, (G_2)^2, G_1 G_2$$

is not sufficient to fix spin dependence

→ + Λ NN contact terms **without decuplet constraints**



$$\Lambda\text{NN} \propto C'_1, C'_2, C'_3$$

ad hoc choice: alter C_2 :

$$C'_1 = C'_3 = \frac{(G_1 + 3G_2)^2}{72\Delta}$$
$$C'_2 = 0$$

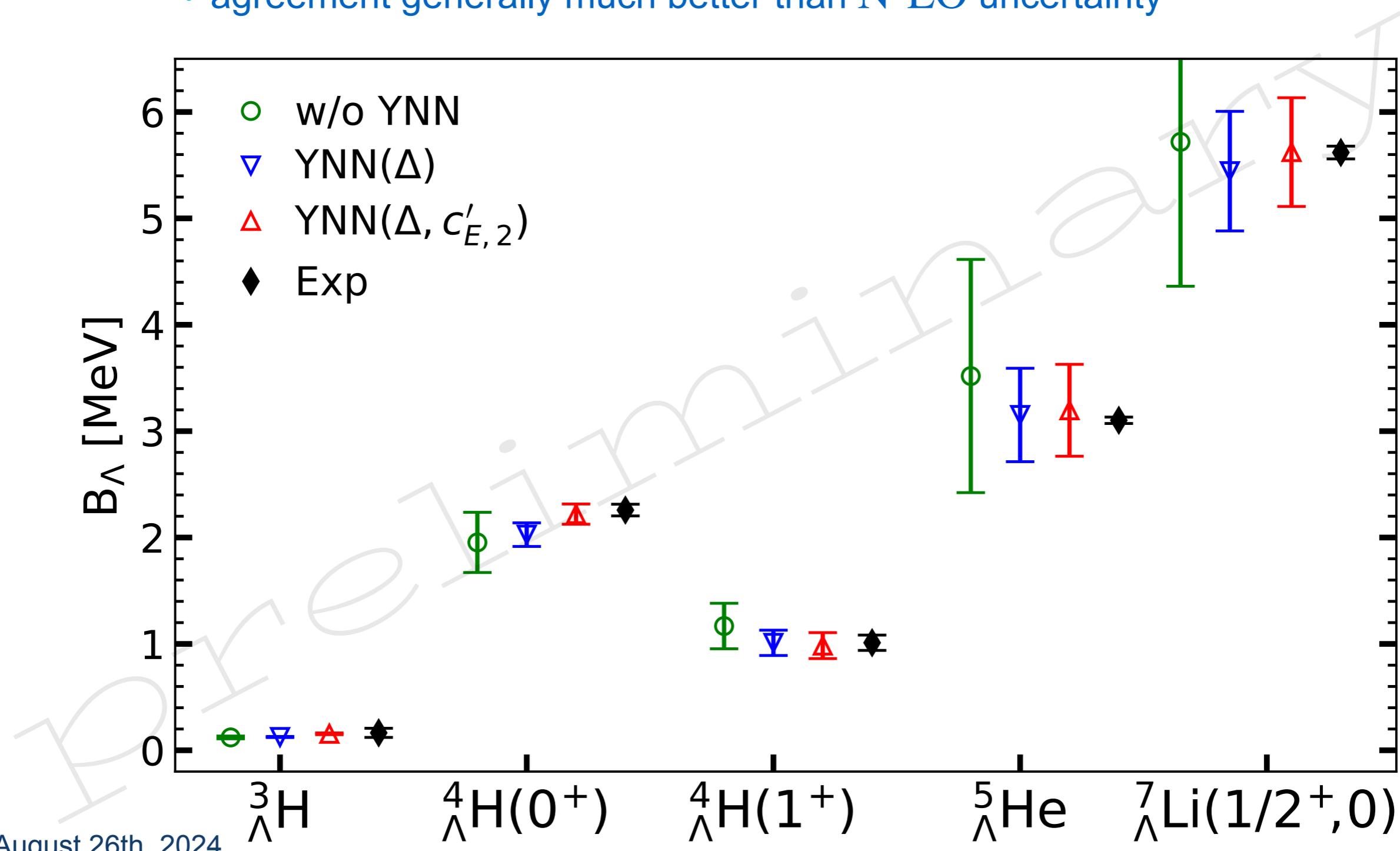


$$V_{\Lambda\text{NN}} = C'_2 \vec{\sigma}_1 \cdot (\vec{\sigma}_2 + \vec{\sigma}_3) (1 - \vec{\tau}_2 \cdot \vec{\tau}_3)$$
$$C'_2 = G_3$$

C'_2 introduces a spin dependent interaction in the most relevant particle channel

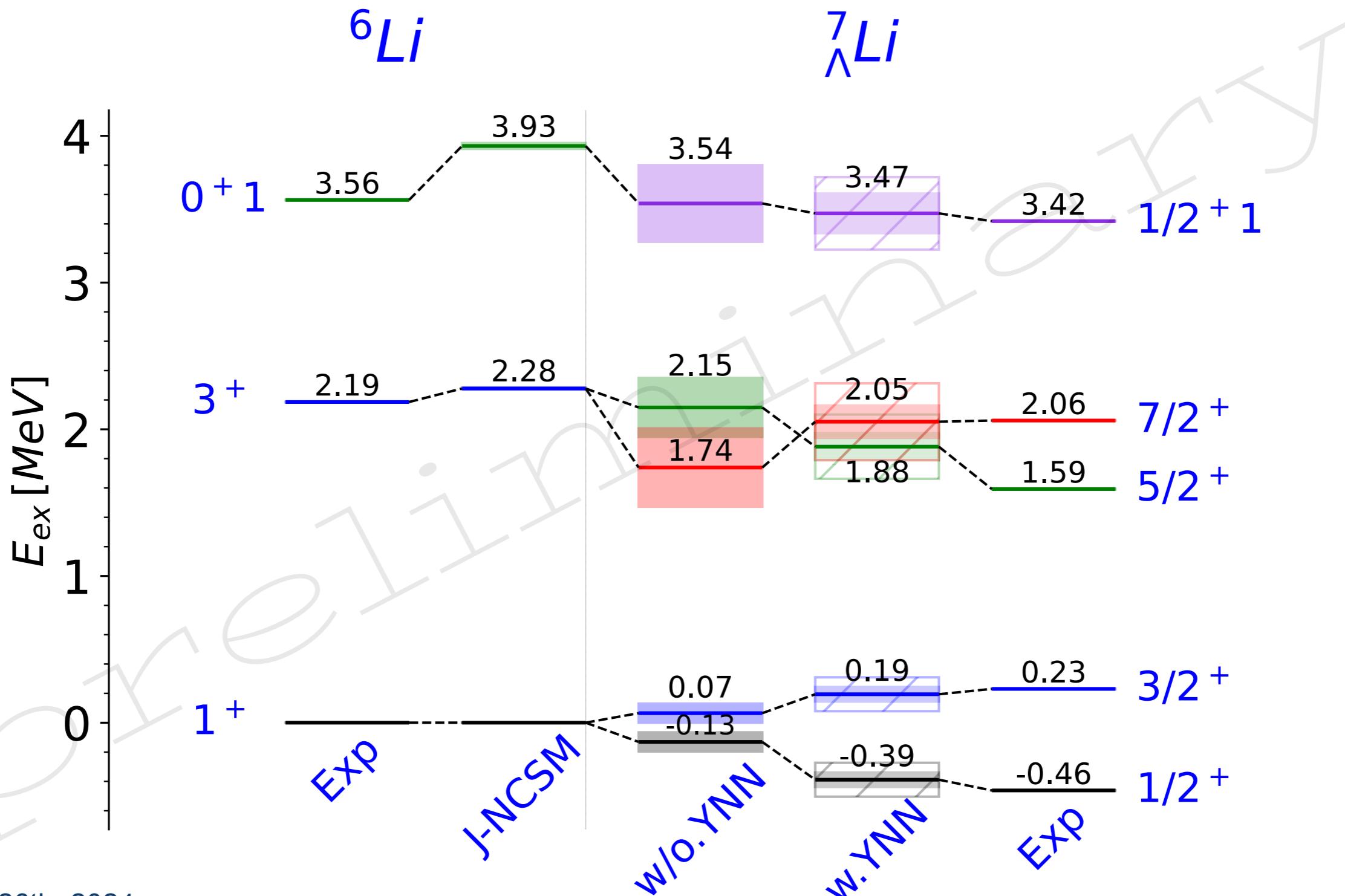


- Fit to 0^+ and 1^+ state of ${}^4_\Lambda\text{He}$ and/or ${}^5_\Lambda\text{He}$
- spin-dependence in $A=4$ not well explained by decuplet saturation
- C'_2 term improves 0^+ of ${}^4_\Lambda\text{He}$ and $1/2^+$ of ${}^7_\Lambda\text{Li}$
- agreement generally much better than N^2LO uncertainty



YNN prediction for $^7\Lambda$ Li

- good agreement
- C'_2 term included, but not very important (not shown)
- higher states have significant uncertainty



Conclusions & Outlook

- YN interactions not well understood
 - scarce YN data
 - more information necessary to solve "hyperon puzzle"
- Hypernuclei provide important constraints
 - CSB of ΛN scattering & ${}^4_\Lambda\text{He} / {}^4_\Lambda\text{H}$
 - ${}^3_\Lambda\text{H}$ is used to constrain the spin dependence
 - new experiments & analyses planned at J-PARC, MAMI, J-Lab, FAIR, ...
- New SMS YN interactions
 - give an accurate description low energy YN data
 - order LO, NLO and $N^2\text{LO}$ allow uncertainty quantification
 - have a **non-unique** determination of contact interactions (data necessary)
- Chiral 3BF need to be included
 - chiral 3BFs are now available — non-local and SMS regularization
 - but: decuplet saturation alone does not improve spin dependence
 - spin-dependent ΛNN leads to further improvement
- however: uncertainty estimate in $N^2\text{LO}$ of incomplete $N^2\text{LO}$ YNN force?
- study cutoff dependence / application to more p-shell hypernuclei