Sensitivity of the nucleon-deuteron scattering observables to N4LO contact terms of three-nucleon force

R.Skibiński, J.Golak, H.Witała



JAGIELLONIAN UNIVERSITY

LENPIC Collaboration

Jagiellonian University, Kraków Ruhr-Universität, Bochum Forschungszentrum, Jülich Bonn Universität, Ohio State University Iowa State University Technische Universität, Darmstadt Kyutech, Fukuoka IPN, Orsay TRIUMF, Vancouver



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Emulator for the Nd scattering:

Based on:

- H.Witała et al., Few-Body Syst. 62 (2021) 23 "Perturbative Treatment of Three Nucleon Force Contact Terms in Three-Nucleon Faddeev Equations."
- H.Witała et al., Eur. Phys. J. A 57 (2021) 241 "Efficient emulator for solving 3N continuum Faddeev equations with chiral 3NF comprising any number of contact terms."
- H.Witała et al., Phys. Rev. C105 (2022) 054004 "Significance of chiral 3NF contact terms for understanding of elastic nucleon-deuteron scattering"
 R.Skibiński et al., Few-Body Syst 65, 44 (2024) "Impact of the Short-Range N4LO Three-Nucleon Force Components on the Nucleon-Deuteron Spin Correlation Coefficients"
- 1. Formalism new set of equations
- 2. Proof of concept: tests and results on fixing short-range 3NF parameters uncertainties of observables related to uncertainty of parameters





Our standard approach to Nd scattering

Prepare and solve the Faddeev equation

 $T\varphi = tP\varphi + (1+tG_0)V_{123}^{(1)}(1+P)\varphi + tPG_0T\varphi + (1+tG_0)V_{123}^{(1)}(1+P)G_0T\varphi$ Compute amplitudes

$$U = PG_0^{-1} + V_{123}^{(1)}(1+P)\varphi + PT + V_{123}^{(1)}(1+P)G_0T$$
$$U_0 = (1+P)T$$

We work in the PWD scheme

3N state: $|pq\alpha\rangle \equiv |pq(ls)j(\lambda\frac{1}{2})I(jI)JM_{J}(t\frac{1}{2})TM_{T}\rangle$

what means, that PWD of appearing in this equation operators has to be performed.

 $\left\langle p'q'(l's')j'(\lambda'\frac{1}{2})I'(j'I')J'M_{J'}(t'\frac{1}{2})T'M_{T'}\right| \hat{O} \left| pq(ls)j(\lambda\frac{1}{2})I(jI)JM_{J}(t\frac{1}{2})TM_{T} \right\rangle$

After decomposing 3NF:

CPU time required for one run (i.e. one reaction energy, one NN+3NF potential) amounts from approx. 1-8 hrs., depending on number of partial waves, computer parameters, disk space available). Some hardware (GPU, fast memory) or software (parallelization) improvements are still possible but the cake's not worth the candle.





Fixing parmeters of 3NF

- Up to know, i.e. when working at N2LO there are only two free parameters cD and cE.
- Typically ³H and the ²a_{nd} or the differential Nd elastic scattering cross section at one or few energies are used. The latter requieres solving the triton many times and the Fadddeev equation 10-20 times.
- However, now we expect:
- No new 3NF free parameters at N3LO, but three new offshell LECs in the chiral NN force.
- 13 contact terms at N4LO (more precisely, due to some identities between opertors, one expect in total 13 free parameters of 3NF at N4LO).
- Thus finding an efficient emulator for solving the 3N Faddeev equation seems to be essential and of high priority.





Emulator for Nd scattering - algorithm

- The contact terms are restricted to small 3N total angular momenta and to only few partial-wave states for a given total 3N angular momentum J and parity π
- Let us spilt 3NF $V_{123}^{(1)} = V(\theta_0) + \Delta V(\theta) \equiv V(\theta_0) + \sum_{i=1}^n c_i \Delta V_i$ $\theta = \{c_1, c_2, \dots, c_n\}$ $\theta_0 = \{0, 0, \dots, 0\}$
- We divide the 3N partial-wave states into two sets:
- 1. The β set is defined by non-vanishing matrix elements of parameters dependent short-range 3NF: $\Delta V(\theta)$.
- 2. The α set comprises remaining states.
- Similarly to 3NF $T = T(\theta_0) + \Delta T(\theta)$





Emulator for Nd scattering - algorithm

Inserting this to the Faddeev equation leads to sets of equations with one equation for $T(\theta_0)$ which is a standard Faddeev equation but with $V(\theta_0)$ and

 $\langle \alpha | \Delta T(\theta) | \phi \rangle = \langle \alpha | t P G_0 \Delta T(\theta) | \phi \rangle + \langle \alpha | (1 + t G_0) V(\theta_0) (1 + P) G_0 \Delta T(\theta) | \phi \rangle$

$$\begin{split} \langle \beta | \Delta T(\theta) | \phi \rangle &= \langle \beta | (1 + tG_0) \Delta V(\theta) (1 + P) | \phi \rangle + \langle \beta | (1 + tG_0) \Delta V(\theta) (1 + P) G_0 T(\theta_0) | \phi \rangle \\ &+ \langle \beta | (1 + tG_0) [V(\theta_0) + \Delta V(\theta)] (1 + P) G_0 \Delta T(\theta) | \phi \rangle \\ &+ \langle \beta | tP G_0 \Delta T(\theta) | \phi \rangle \,. \end{split}$$

We neglect term ~ΔVΔT, which allows to separate contributions from different ΔV_i and leads to set of equations for a corresponding ΔT_i. Next, for single parameter dependent component of V: V_i=c_iV we may solve that equation separately (at c_i=1) obtaining corresponding <β|ΔT_i|φ> and build

$$\langle \beta | \Delta T(\theta) | \phi \rangle = \sum_{i=1}^{N} c_i \langle \beta | \Delta T_i | \phi \rangle$$

and find also $<\alpha |\Delta T_i|\phi>$.





Emulator for Nd scattering - algorithm

 In this way we have matrix elements of T

$$\langle \alpha | T(\theta) | \phi \rangle = \langle \alpha | T(\theta_0) | \phi \rangle + \sum_i c_i \langle \alpha | \Delta T_i | \phi \rangle,$$

$$\langle \beta | T(\theta) | \phi \rangle = \langle \beta | T(\theta_0) | \phi \rangle + \sum_i c_i \langle \beta | \Delta T_i | \phi \rangle.$$
(10)

- Let us now come back to the scattering amplitudes $U = PG_0^{-1} + V_{123}^{(1)}(1+P)\phi + PT + V_{123}^{(1)}(1+P)G_0T$ $U_0 = (1+P)T$
- They are linear in T: the dependence on the c_i parameters carries over to them

$$U = U(\theta_0) + \sum_i c_i U_i + \sum_{i,k} c_i c_k U_{ik}$$
$$U_0 = U_0(\theta_0) + \sum_i c_i U_{0i}$$

• Summarizing: one needs to solve N+1 Faddeev equations (one for $T(\theta_0)$ and N for $<\beta|\Delta T_i|\phi>$) and next build transition amplitudes for any set of c_i .





Emulator for Nd scattering – algorithm - application

- We used SMS N4LO+ NN potential at Λ=450 MeV, combined with the N2LO chiral 3NF and supplemented by all subleading N4LO 3NF contact terms from:
 1. L. Girlanda, A. Kievsky, and M. Viviani, Phys. Rev. C 84, 014001 (2011).,
 2. L. Girlanda, A. Kievsky, and M. Viviani, Phys. Rev. C 102, 019903(E) (2020).
- All terms are regulated with the non-local regulator.
- Such a Hamiltonian comprises altogether 15 short-range contributions to 3NF, two from N2LO with the strengths cD and cE, and thirteen from N4LO with the strengths E_i, i = 1, ..., 13. However, for two pairs of the E_i terms matrix elements are identical, thus finally there are 13 unknown parameters.
- We tested that taking β= {j<3} is preety good approximation: predictions from the emulator are practically indistinguishable from the full solution.





Emulator for Nd scattering – algorithm - application



N2LO D and E terms do not dominate

- $\Delta \equiv \Delta(c_i) = \frac{1}{N_{\theta}} \sum_{\theta_k} \frac{Obs(c_i, \theta_k) Obs(\theta_0, \theta_k)}{Obs(\theta_0, \theta_k)}$
- Some observables are more sensitive to specific terms, e.g. T₂₂ to E₁₀







Emulator for Nd scattering – fit to the true data at 10,70, and 135 MeV (786 data points)

TABLE III. The values of strengths c_i found in the least squares fit to the data from Table II at the three energies E = 10, 70, and 135 MeV.

		c_D	c_E	c_{E_1}	c_{E_2}	c_{E_3}	c_{E_4}	C_{E_5}	C_{E_6}	c_{E_7}	c_{E_8}	C_{E_9}	$c_{E_{10}}$	$c_{E_{13}}$
C _D	-1.49 ± 0.06													
CF	-1.27 ± 0.06	1.000	-0.122	0.068	0.202	0.040	0.107	-0.237	-0.566	0.124	0.094	-0.052	0.038	-0.066
- <u>L</u>	6.40 ± 0.33	-0.122	1.000	0.048	-0.166	0.067	-0.084	-0.059	-0.144	-0.284	-0.298	-0.127	0.128	0.191
c_{E_1}	0.40 ± 0.33	0.068	0.048	1.000	0.945	0.982	-0.843	0.087	-0.291	0.551	0.500	0.176	0.228	0.027
c_{E_2}	7.80 ± 0.30	0.202	-0.166	0.945	1.000	0.934	-0.770	-0.112	-0.356	0.629	0.581	-0.007	0.041	-0.184
C_{E_3}	6.97 ± 0.34	0.040	0.067	0 982	0.934	1,000	-0.917	0.098	-0.214	0.574	0.528	0.111	0.093	-0.012
c_{E_4}	-2.06 ± 0.13	0 107	0 001	0 9/2	0 770	0 017	1 000	0 217	0 11/	0 592	0 556	0 172	0.053	0 025
$C_{E_{\sigma}}$	-0.36 ± 0.05	0.107	-0.004	-0.045	-0.770	-0.917	1.000	-0.51/	-0.114	-0.505	-0.550	-0.1/2	0.055	-0.035
25 Ca	0.52 ± 0.03	-0.23/	-0.059	0.08/	-0.112	0.098	-0.31/	1.000	0.540	0.113	0.126	0.898	0.521	0./19
c_{E_6}	0.52 ± 0.05	-0.566	-0.144	-0.291	-0.356	-0.214	-0.114	0.540	1.000	-0.243	-0.197	0.253	-0.175	0.167
c_{E_7}	-7.40 ± 0.14	0.124	-0.284	0.551	0.629	0.574	-0.583	0.113	-0.243	1.000	0.995	0.108	-0.044	-0.205
c_{E_8}	-2.61 ± 0.05	0.094	-0.298	0.500	0.581	0.528	-0.556	0.126	-0.197	0.995	1.000	0.102	-0.068	-0.203
C_{E_9}	-4.59 ± 0.22	-0.052	-0.127	0.176	-0.007	0.111	-0.172	0.898	0.253	0.108	0.102	1,000	0.795	0.787
$c_{E_{10}}$	-0.98 ± 0.05	0 038	0 128	0 228	0 041	0 093	0 053	0 521	-0 175	-0 044	-0 068	0 795	1 000	0 746
CE	-1.14 ± 0.05	0.050	0.101	0.007	0.104	0.010	0.000	0.710	0.167	0.205	0.000	0.707	0.746	1 000
CE13	1.11 ± 0.05	(-0.000	0.191	0.02/	-0.184	-0.012	-0.035	0./19	0.10/	-0.205	-0.203	0./8/	0.746	1.000)

- Big values of c_{E1}, c_{E2}, c_{E3}, c_{E7}, c_{E9}
- Correlation coefficients close to ±1: $\rho(E_1,E_2)$, $\rho(E_2,E_3)$, $\rho(E_1,E_3)$, $\rho(E_3,E_4)$, $\rho(E_7,E_8)$
- Correlation coefficients close to 0: (c_D,c_E),(c_D,c_{Ei, beside E6}),(c_E,c_{Ei})
- ∎ χ²/data≈35





Emulator for Nd scattering – fit to the data: cross section and $A_V(N)$

NN N4LO+ + 3NF N2LO

- Data at 10, 70 and 135 MeV
 - Results at 190 and 250 MeV are predictions



NN N4LO+





NN N4LO+ + 3NF N2LO + E_i



Emulator for Nd scattering – fit to the data: iT_{11} and T_{20}

NN N4LO+ + 3NF N2LO

- Data at 10, 70 and 135 MeV
 - Results at 190 and 250 MeV are predictions



NN N4LO+



NN N4LO+ + 3NF N2LO + E_i



Emulator for Nd scattering – fit to the data: C_{yy} and $C_{yz,x}$

- Data at a) 10 and b) 135 MeV
- Results at c) 190 and d) 250 MeV are predictions



 \rightarrow A comparison with data (B. Przewoski at al., Phys. Rev. C 74, 064003 (2006)) at E = 135 MeV reveals both improvement and deterioration of the description depending on the observables.





Statistical uncertainty

Short-range parameters are correlated \rightarrow changes in their values cannot be arbitrary \rightarrow this affects possible changes in the values of observables

To study that:

- we sampled N=10⁶ sets of values c_D, c_E and 11 c_{Ei} parameters from the correlated multivariate normal distribution
- For each set, using our emulator, we computed observables (at E=70 MeV and E=250 MeV)
- In that way, for each observable and scattering angle, we have 10⁶ predictions
- Resulting standard deviation of mean of observables (due to the factor $\frac{1}{\sqrt{N(N-1)}}$) is small <0.1%
- In following I will show bands limited by max and min value of observable over all N predictions





Parameters-observables correlations

Big sample opens opportunity to study

various statistical aspects 0.1 -1.5 -1 0.2 +σ 0.18 $\rho = -0.403$ C_r -3.5 0.16 0.14 $+3\sigma$ 0.12 ം ച് щ 9 10 5 8 $\rho = 0.187$ -3σ C_{E5} -5.5 -6^L₋₂ -1.8 -1.6 -1.4 -1.2 -1 5 6 8 9 CD \mathbf{E}_{1} -0.3 0 06 5e+04 1e+05 $\bar{D} = -0.287 C$ 9e+04 0.16 4e+04 8e+04 0.14 0.12 7e+04 3e+04 0.1 6e+04 -2.4 -6 -2.7 5e+04 0.2 2e+04 4e+04 0.18 $C_{E13} \rho = 0.169$ 3e+04 0.16 1e+04 2e+04 0.14 1e+04 0.12 0e+00 0.10.072 0e+000.074 0.12 0.13 0.14 C_{YY} (Θ=130 deg, E=250 MeV) 0.11 0.15 0.16 $d\sigma/d\Omega$ [mb sr⁻¹] (Θ =130 deg, E=250 MeV)







Statistical uncertainty (min-max over N=10⁶)













Summary

- We constructed and tested an efficient and accurate emulator for solving 3N Faddeev equation.
- Our emulator allows us to fix free parameters of all short-range terms in the 3NF. We found that even at low energies some observables are sensitive to N4LO 3NF contact terms.
- In general, sensitivity of predictions to N4LO 3NF contact terms depends on observable, energy and scattering angle.
- Usually, we observe improvements in data description, but very likely above
 ≈200 MeV 3NF is not sufficient to explain discrepancies with the data.
- Theoretical uncertainty related to uncertainty of stddev of short-range 3NF parameters remains small, but not negligible at higher energies, and comparable with other types of theoretical uncertainties. Thus, it should be included in the uncertainty budget.
- 3NF at N3LO must be included for final conclusions.





18