

# Semileptonic Kaon Decays and the Precise Determination of V<sub>us</sub>

#### **Chien-Yeah Seng**

#### University of Washington and FRIB, Michigan State University

seng@frib.msu.edu

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CYS, Galviz and Meißner, 2020 JHEP; CYS, Feng, Gorchtein, Jin and Meißner, 2020 JHEP; Ma, Feng, Gorchtein, Jin and CYS, 2021 PRD; CYS, Galviz, Gorchtein and Meißner, 2021 PLB, 2021 JHEP, 2022 JHEP; CYS, Galviz, Marciano and Meißner, 2022 PRD

#### In collaboration with:

Xu Feng Daniel Galviz Mikhail Gorchtein Luchang Jin Peng-Xiang Ma William Marciano Ulf-G. Meißner

# 1. First-row CKM unitarity

## **Charged weak decays and CKM unitarity**



Massive quarks ==> Generation mixing

$$\psi_{d,f} = \begin{pmatrix} d \\ s \\ b \end{pmatrix}_{f} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}_{m}$$

#### Cabibbo-Kobayashi-Maskawa (CKM) matrix

# Three generations of quarks and leptons





Weak interaction universality ==> Unitarity of the measured CKM matrix

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} V_{ud}^* & V_{cd}^* & V_{td}^* \\ V_{us}^* & V_{cs}^* & V_{ts}^* \\ V_{ub}^* & V_{cb}^* & V_{tb}^* \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$
  
**First row:**  
$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$
  
~10<sup>-5</sup>

Can be tested at 0.01% level! Probes new physics at the scale:

$$\left(\frac{v_{\rm H}}{\Lambda_{\rm BSM}}\right)^2 \sim 0.01\% \implies \Lambda_{\rm BSM} \sim 20 \,{\rm TeV}$$

**Primary avenues to extract V**<sub>ud</sub> and V<sub>us</sub>

V<sub>ud</sub>: Superallowed nuclear decays  $\phi_i(0^+) \rightarrow \phi_f(0^+) \ e^+ \nu_e$ V<sub>us</sub>: Semileptonic kaon decays (K<sub>12</sub>)

$$K \to \pi \ell^+ \nu_\ell$$

 $V_{us}/V_{ud}$ : K/ $\pi$  leptonic decays (K<sub>µ2</sub>/ $\pi_{µ2}$ )

$$K/\pi \to \mu^+ \nu_\mu$$

2. Kaon semileptonic decays (K<sub>13</sub>) and the long-distance RC



Six channels:  $K_{e3}^L$  :  $K^L \to \pi^- e^+ \nu_e$  $K^L_{\mu3}$  :  $K^L \to \pi^- \mu^+ \nu_\mu$  $K_{e^3}^S : K^S \to \pi^- e^+ \nu_e$  $K^S_{\mu3}$  :  $K^S \to \pi^- \mu^+ \nu_\mu$  $K_{e3}^+$ :  $K^+ \to \pi^0 e^+ \nu_e$  $K^{+}_{\mu3} : K^{+} \to \pi^{0} \mu^{+} \nu_{\mu}$ 



$$\Gamma_{K_{\ell 3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{\rm EW} |f_+^{K^0 \pi^-}(0)|^2 I_{K\ell}^{(0)} \left(1 + \delta_{\rm EM}^{K\ell} + \delta_{\rm SU(2)}^{K\pi}\right)$$



## **Theory inputs:**

- Universal electroweak correction
- $K\pi$  form factor at zero momentum transfer
- Phase space factor
- Isospin-breaking correction



$$\Gamma_{K_{\ell3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} \sum_{\text{EW}} |f_+^{K^0 \pi^-}(0)|^2 I_{K\ell}^{(0)} \left(1 + \delta_{\text{EM}}^{K\ell} + \delta_{\text{SU}(2)}^{K\pi}\right)$$

#### Universal electroweak correction:

$$S_{\rm EW} = 1 + \frac{2\alpha}{\pi} \left( 1 - \frac{\alpha_s}{4\pi} \right) \ln \frac{M_Z}{M_\rho} + \mathcal{O}\left(\frac{\alpha \alpha_S}{\pi^2}\right)$$
$$= 1.0232(3)$$

encodes process-independent electroweak corrections not included in  $G_{F}$  Marciano and Sirlin, 1993 PRL



$$\Gamma_{K_{\ell3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{\rm EW} f_+^{K^0 \pi^-}(0) \mathcal{I}_{K\ell}^{(0)} \left(1 + \delta_{\rm EM}^{K\ell} + \delta_{\rm SU(2)}^{K\pi}\right)$$

K
$$\pi$$
 form factor at t=0:  $\langle \pi^{-}(p) | J_{W}^{\mu} | K^{0}(p) \rangle = 2p^{\mu} f_{+}^{K^{0}\pi^{-}}(0)$ 

#### **Lattice QCD inputs:**

$$N_f = 2 + 1 + 1 \quad : \quad f_+(0) = 0.9698(17)$$
$$N_f = 2 + 1 \quad : \quad f_+(0) = 0.9677(27)$$
$$N_f = 2 \quad : \quad f_+(0) = 0.9560(57)(62)$$





$$\Gamma_{K_{\ell 3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{\rm EW} |f_+^{K^0 \pi^-}(0)| \left( I_{K\ell}^{(0)} \left( 1 + \delta_{\rm EM}^{K\ell} + \delta_{\rm SU(2)}^{K\pi} \right) \right) \right)$$

#### Phase-space factor: probes the t-dependence of the form factors

Fit to the **K**<sub>13</sub> **Dalitz** plot with **dispersive parameterization**:

Mode	Update	M Maulaan
$K^{0}_{e3}$	0.15470(15)	in the 11 <sup>th</sup> International Workshop on the CKM Unitarity
$K^{+}_{e3}$	0.15915(15)	
$K^{0}_{\ \mu 3}$	0.10247(15)	Triangle, 2021
$K^{+}_{~\mu3}$	0.10553(16)	



$$\Gamma_{K_{\ell 3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{\rm EW} |f_+^{K^0 \pi^-}(0)|^2 I_{K\ell}^{(0)} \left(1 + \delta_{\rm EM}^{K\ell} + \delta_{\rm SU(2)}^{K\pi}\right)$$

**ISB correction:** Difference between the  $K^+\pi^0$  and  $K^0\pi^-$  form factor Depends on quark mass parameters:

$$m_s/\hat{m}$$
 ,  $Q^2 \equiv (m_s^2 - \hat{m}^2)/(m_d^2 - m_u^2)$ 

Lattice and phenomenology  $(\eta \rightarrow 3\pi)$  return somewhat different results:

$$\begin{pmatrix} \delta_{\mathrm{SU}(2)}^{K^{+}\pi^{0}} \\ \mathrm{attice} \end{pmatrix}_{\text{lattice}} = 0.0457(20) \quad \text{FLAG 2021 (2023 update)}$$
$$\begin{pmatrix} \delta_{\mathrm{SU}(2)}^{K^{+}\pi^{0}} \\ \mathrm{pheno} \end{pmatrix}_{\text{pheno}} = 0.0522(34) \quad \begin{array}{c} \text{Colangelo, Lanz, Leutwyler} \\ \text{and Passemar, 2018 EPJC} \end{array}$$



$$\Gamma_{K_{\ell 3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{\rm EW} |f_+^{K^0 \pi^-}(0)|^2 I_{K\ell}^{(0)} \left(1 + \delta_{\rm EM}^{K\ell} + \delta_{\rm SU(2)}^{K\pi}\right)$$

#### Long-distance electromagnetic radiative corrections (RC)



"Long-distance" : W-propagator shrinks to a point (Fermi's interaction)

# 3. Chiral Perturbation Theory

# **Chiral perturbation theory (ChPT)**

- Low-energy EFT of QCD
- Constructed from spontaneously-broken chiral symmetry
- DOFs: Pseudoscalar mesons, leptons, photon
- Chiral power-counting scheme ensure convergence
- NP QCD contained in the LECs

$$\mathcal{L} = \mathcal{L}_{\text{lepton}} + \mathcal{L}_{\gamma} + \mathcal{L}_{\text{ChPT}}$$

$$\mathcal{L}_{\text{lepton}} = \sum_{\ell} [\bar{\ell}(i\partial + eA - m_{\ell})\ell + \bar{\nu}_{\ell L}i\partial \bar{\nu}_{\ell L}]$$

$$\mathcal{L}_{\gamma} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{2\xi}(\partial \cdot A)^{2} + \frac{1}{2}M_{\gamma}^{2}A_{\mu}A^{\mu}$$

$$\mathcal{L}_{\text{ChPT}} = \mathcal{L}^{(2)} + \mathcal{L}^{(4)} + \dots$$

$$\mathcal{L}^{p^{2}} = \frac{F_{0}^{2}}{4} \langle D_{\mu}U(D^{\mu}U)^{\dagger} + \chi U^{\dagger} + U\chi^{\dagger} \rangle, \quad \mathcal{L}^{e^{2}} = ZF_{0}^{4} \langle q_{L}U^{\dagger}q_{R}U \rangle$$

#### LO Chiral Lagrangian

#### **NLO Chiral Lagrangian:**



**One-loop and bremsstrahlung diagrams** are calculated

### **Connection to the full EW theory:**

Done through a specific combination of LECs:

$$X_6^{\text{phys}} \equiv X_6^r - 4K_{12}^r$$
$$= \left(X_6^{\text{phys}}\right)_{\text{SD}} + \left(X_6^{\text{phys}}\right)_{\text{LD}}$$
$$S_{\text{EW}} = 1 - e^2 \left(X_6^{\text{phys}}\right)_{\text{SD}}$$

#### **Quantification of hadronic uncertainties:**

- Uncertainties due to neglected terms at O(e<sup>2</sup>p<sup>4</sup>)
- Uncertainties due to unknown LECs at O(e<sup>2</sup>p<sup>2</sup>)

	$\delta^{K\ell}_{ m EM}(\%)$
$K_{e3}^0$	$0.99\pm 0.19_{e^2p^4}\pm 0.11_{\rm LEC}$
$K_{e3}^{\pm}$	$0.10\pm 0.19_{e^2p^4}\pm 0.16_{\rm LEC}$
$K^0_{\mu 3}$	$1.40\pm 0.19_{e^2p^4}\pm 0.11_{\rm LEC}$
$K_{\mu 3}^{\pm}$	$0.016\pm 0.19_{e^2p^4}\pm 0.16_{\rm LEC}$

Ananthanarayan and Moussallam, JHEP 2004 Descotes-Genon and Moussallam, EPJC 2005

Uncertainty ~10-3

"Natural limitations" of ChPT precision

Cirigliano, Giannotti and Neufeld, JHEP 2008

To overcome the natural limitations:

- New theory framework to resum the most important O(e<sup>2</sup>p<sup>n</sup>) contributions
- Appropriate lattice-QCD inputs to reduce uncertainties from NP QCD

Direct lattice calculations of the full  $K_{13}$  RC: ~10 years to reach 10<sup>-3</sup> precision

Boyle et al., SnowMass 2021 Lol

# 4. Sirlin's representation

## <u>"Sirlin's representation" of the O(G<sub>F</sub> $\alpha$ ) EWRC</u>

- First constructed by Sirlin to deal with EWRC in superallowed beta decays *Sirlin, 1978 Rev.Mod.Phys*
- Re-introduced to study EWRC in general semileptonic decays

CYS, Galviz and Meißner, 2020 JHEP CYS, 2021 Particles

Foundations: Current Algebra (CA)

$$\left[J_W^{0\dagger}(\vec{x},t), J_{\rm em}^{\mu}(\vec{y},t)\right] = J_W^{\mu\dagger}(\vec{x},t)\delta^{(3)}(\vec{x}-\vec{y})$$

#### Splitting the full **photon propagator**:





EWRCs at q' $\sim$ M<sub>w</sub> at the **lepton side** are reabsorbed into **G**<sub>F</sub> measured from **muon decay** 

$$\frac{1}{\tau_{\mu}} \equiv \frac{G_F^2 m_{\mu}^5}{192\pi^3} F(x)(1+\delta_{\mu})$$
 24

#### Splitting the full **photon propagator**:

$$\frac{1}{q'^2} = \frac{1}{q'^2 - M_W^2} + \frac{M_W^2}{M_W^2 - q'^2} \frac{1}{q'^2}$$
  
"\gamma\_>" \quad \cap\color \color \cap\color \cap\col\cap\color \cap\co



"Modified Fermi's interaction":  $\mathcal{L} = \mathcal{L}_{QCD} + \mathcal{L}_{QED,\gamma_{<}} + \mathcal{L}'_{4f}$ 

$$\mathcal{L}'_{4f} = -\left\{1 - \frac{\alpha}{2\pi} \left[\ln\frac{M_W^2}{M_Z^2} + \mathcal{O}(\alpha_s)\right]\right\} \frac{G_F}{\sqrt{2}} J^{\mu} \,\bar{l}\gamma_{\mu}(1-\gamma_5)\nu + h.c.$$

#### Virtual electromagnetic RC (EMRC)





Real EMRC (Bremsstrahlung)

#### Key quantity: "Generalized Compton tensor"

$$T^{\mu\nu}(q';p',p) = \int d^4x e^{iq'\cdot x} \langle \phi_f(p') | T\{J^{\mu}_{\rm em}(x)J^{\nu}_W(0)\} | \phi_i(p) \rangle$$



The "convection term": Meister and Yennie, PR 1963

$$T_{\rm conv}^{\mu\nu}(q';p',p) = \frac{iZ_f(2p'+q')^{\mu}F^{\nu}(p',p)}{(p'+q')^2 - M_f^2} + \frac{iZ_i(2p-q')^{\mu}F^{\nu}(p',p)}{(p-q')^2 - M_i^2}$$

Simplest structure satisfying the **exact EM Ward identity**; gives the **IR-divergent structure** 

#### EMRC to the hadron form factors:



"Two-point function": contains all the UV-dependence

$$\delta F_2^{\lambda}(p',p) = -e^2 \eta \int \frac{d^4 q'}{(2\pi)^4} \left[ \frac{M_W^2}{(M_W^2 - q'^2)^2} \frac{q'^{\lambda}}{q'^2 - M_\gamma^2} - \frac{M_W^2}{M_W^2 - q'^2} \frac{q'^{\lambda}}{(q'^2 - M_\gamma^2)^2} \right] T^{\mu}_{\mu}$$

 $\delta F_3^{\lambda}(p', p)$  **"Three-point function"**: insensitive to UV, vanishes in the degenerate + forward limit

$$\sim \langle f | T \{ J_W(x) J_{em}(y) J_{em}(0) \} | i \rangle$$

## γW-box diagram:



$$\delta\mathfrak{M}_{\gamma W} = \delta\mathfrak{M}^a_{\gamma W} + \delta\mathfrak{M}^b_{\gamma W}$$

$$\delta\mathfrak{M}^{a}_{\gamma W} = \eta \frac{G_{F}e^{2}}{\sqrt{2}} L_{\lambda} \int \frac{d^{4}q'}{(2\pi)^{4}} \frac{M_{W}^{2}}{M_{W}^{2} - q'^{2}} \frac{2g^{\nu\lambda}p_{\ell}^{\mu} - g^{\mu\lambda}q'^{\nu} - g^{\nu\lambda}q'^{\mu} + g^{\mu\nu}q'^{\lambda}}{[(p_{\ell} - q')^{2} - m_{\ell}^{2}][q'^{2} - M_{\gamma}^{2}]} T_{\mu\nu} \quad \text{W/O epsilon tensor}$$

$$\delta\mathfrak{M}^{b}_{\gamma W} = -i\frac{G_{F}e^{2}}{\sqrt{2}} L_{\lambda} \int \frac{d^{4}q'}{(2\pi)^{4}} \frac{M_{W}^{2}}{M_{W}^{2} - q'^{2}} \frac{\epsilon^{\mu\nu\alpha\lambda}q'_{\alpha}}{[(p_{\ell} - q')^{2} - m_{\ell}^{2}]q'^{2}} T_{\mu\nu} \quad \text{With epsilon tensor}$$

## Sirlin's representation: $O(\alpha)$ virtual EWRC to semileptonic decays



# **5. Electroweak RC to K\_{13}**

K
$$\pi$$
 charged weak form factors  

$$F_{\mu}^{K\pi}(p',p) \equiv \langle \pi(p') | (J_{\mu}^{W})^{\dagger} | K(p) \rangle = V_{us}^{*} [f_{\pm}^{K\pi}(t)(p+p')_{\mu} + f_{-}^{K\pi}(t)(p-p')_{\mu}]$$
Virtual corrections expressed as corrections to  
form factors:  

$$f_{\pm}^{K\pi}(t) \rightarrow f_{\pm}^{K\pi}(t) + \delta f_{\pm}^{K\pi}(y,z)$$

$$y = \frac{2E_{\ell}}{M_{K}}$$

$$z = \frac{2E_{\pi}}{M_{K}}$$

Contains IR-divergences canceled by bremsstrahlung

# Calculation of $\delta f_{+}$ :

(A) "Residual integral" + the vector current contribution to  $\delta \mathfrak{M}^b_{\gamma W}$ 

$$I_{\mathfrak{A}}^{\lambda} = -e^{2} \int \frac{d^{4}q'}{(2\pi)^{4}} \frac{1}{[(p_{\ell} - q')^{2} - m_{\ell}^{2}][q'^{2} - M_{\gamma}^{2}]} \left\{ \frac{2p_{\ell} \cdot q'q'^{\lambda}}{q'^{2} - M_{\gamma}^{2}} + 2p_{\ell\mu}T^{\mu\lambda} - (p - p')_{\mu}T^{\lambda\mu} + i\Gamma^{\lambda} - i\epsilon^{\mu\nu\alpha\lambda}q'_{\alpha}(T_{\mu\nu})_{V} \right\}$$

Contribution to  $f_{\downarrow}$  is saturated by the "pole" terms" :



- Inputs:  $K/\pi$  EM form factors and charged weak form factors, well-measured in experiments
- Effectively re-summing the most important O(e<sup>2</sup>p<sup>n</sup>) corrections

# **Calculation of** $\delta f_{+}$ :

(B)Axial current contribution to  $\delta \mathfrak{M}_{\gamma W}^{b}$ 

$$I_{\mathfrak{B}}^{\lambda} = ie^2 \int \frac{d^4q'}{(2\pi)^4} \frac{M_W^2}{M_W^2 - q'^2} \frac{\epsilon^{\mu\nu\alpha\lambda}q'_{\alpha}(T_{\mu\nu})_A}{[(p_{\ell} - q')^2 - m_{\ell}^2]q'^2}$$



#### Obtained from **lattice + pQCD**



Feng, Gorchtein, Jin, Ma and CYS, 2020 PRL Ma, Feng, Gorchtein, Jin and CYS, 2021 PRD

# **Calculation of** $\delta f_{+}$ :

(c)  $\delta f_+$  contributed by the "3-pt function" Seng, Galviz and Meißner, 2020 JHEP



## **Calculation of** $\delta f_{-}$ :

suppressed in decay rate by  $m_\ell^2/M_K^2$ 



Cirigliano et al., 2002 EPJC

$$(\delta f_{-}^{K^{0}\pi^{-}})_{\rm rem}^{e^{2}p^{2}} = -\frac{\alpha}{4\pi} \left[ r_{0,1}\Lambda(s, M_{\pi}, m_{\ell}) + r_{0,2}\ln\frac{M_{\pi}^{2}}{m_{\ell}^{2}} + r_{0,3}\ln\frac{M_{\pi}^{2}}{\mu^{2}} + r_{0,4} \right]$$

$$(\delta f_{-}^{K^{+}\pi^{0}})_{\rm rem}^{e^{2}p^{2}} = -\frac{\alpha}{4\sqrt{2\pi}} \left[ r_{+,1}\Lambda(u, M_{K}, m_{\ell}) + r_{+,2}\ln\frac{M_{K}^{2}}{m_{\ell}^{2}} + r_{+,3}\ln\frac{M_{K}^{2}}{\mu^{2}} + r_{+,4} \right]$$

#### **Bremsstrahlung:**



## Final Result: (in units of 10<sup>-3</sup>)

CYS, Galviz, Gorchtein and Meißner, 2021 PLB; 2021 JHEP

CYS, Galviz, Gorchtein and Meißner, 2022 JHEP

	$\delta^{K\ell}_{ m EM}$	ChPT
$K^0e$	$11.6(2)_{inel}(1)_{lat}(1)_{NF}(2)_{e^2p^4}$	$9.9(1.9)_{e^2p^4}(1.1)_{\rm LEC}$
$K^+e$	$2.1(2)_{\rm inel}(1)_{\rm lat}(4)_{\rm NF}(1)_{e^2p^4}$	$1.0(1.9)_{e^2p^4}(1.6)_{\text{LEC}}$
$K^0\mu$	$15.4(2)_{\text{inel}}(1)_{\text{lat}}(1)_{\text{NF}}(2)_{\text{LEC}}(2)_{e^2p^4}$	$14.0(1.9)_{e^2p^4}(1.1)_{\text{LEC}}$
$K^+\mu$	$0.5(2)_{\text{inel}}(1)_{\text{lat}}(4)_{\text{NF}}(2)_{\text{LEC}}(2)_{e^2p^4}$	$0.2(1.9)_{e^2p^4}(1.6)_{\text{LEC}}$

Sources of uncertainty:

- inel: inelastic states contributions to the residual integral
- **lat**: Lattice uncertainty in the  $\gamma$ W-box diagram
- NF: Non-forward effects in the  $\gamma$ W-box diagram
- e<sup>2</sup>p<sup>4</sup>: Higher-order ChPT corrections
- LEC: Poorly-determined LECs in  $\delta f_{\underline{}}$ 
  - Consistent with pure ChPT result
  - Significant improvement of precision:  $10^{-3} \rightarrow 10^{-4}$



$$\Gamma_{K_{\ell 3}} = \frac{G_F^2 |V_{us}|^2 M_K^5 C_K^2}{192\pi^3} S_{\rm EW} |f_+^{K^0 \pi^-}(0)|^2 I_{K\ell}^{(0)} \left(1 + \delta_{\rm EM}^{K\ell} + \delta_{\rm SU(2)}^{K\pi}\right)$$

#### **Averaging over all six channels:**

	$ V_{us}f_{+}^{K^{0}\pi^{-}}(0) $	
$K_L e$	$0.21617(46)_{\exp}(10)_{I_K}(4)_{\delta_{\mathrm{EM}}}$	
$K_S e$	$0.21530(122)_{\exp}(10)_{I_K}(4)_{\delta_{\rm EM}}$	
$K^+e$	$0.21714(88)_{\exp}(10)_{I_K}(21)_{\delta_{\mathrm{SU}(2)}}(5)_{\delta_{\mathrm{EM}}}$	
$K_L \mu$	$0.21649(50)_{\exp}(16)_{I_K}(4)_{\delta_{\mathrm{EM}}}$	
$K_S \mu$	$0.21251(466)_{\exp}(16)_{I_K}(4)_{\delta_{\mathrm{EM}}}$	
$K^+\mu$	$0.21699(108)_{\exp}(16)_{I_K}(21)_{\delta_{\mathrm{SU}(2)}}(6)_{\delta_{\mathrm{EM}}}$	
Average: $Ke$	$0.21626(40)_K(3)_{\rm HO}$	
Average: $K\mu$	$0.21654(48)_K(3)_{\rm HO}$	
Average: tot	$0.21634(38)_K(3)_{\rm HO}$	

$$|V_{us}|_{K_{\ell 3}} = \begin{cases} \mathsf{N}_{f} = 2 + 1 + 1\\ 0.22308(39)_{\mathrm{lat}}(39)_{K}(3)_{\mathrm{HO}}\\ 0.22356(62)_{\mathrm{lat}}(39)_{K}(3)_{\mathrm{HO}}\\ \mathsf{N}_{f} = 2 + 1 \end{cases}$$

CYS, Galviz, Gorchtein and Meißner, 2022 JHEP

# 6. "Cabibbo angle anomaly"











# Summary

- First-row CKM unitarity offers precision tests of SM;  $\rm K_{_{13}}$  is the primary avenue to extract  $\rm V_{_{us}}.$
- EMRC to the  $\rm K_{\rm l3}$  decay rate requires non-perturbative QCD. ChPT calculation is limited by unknown higher-order contributions and LECs.
- A reformulation using ChPT + current algebra allows inputs from meson form factors + lattice QCD, results in a much-improved determination of the  $K_{13}$  RC, with central values consistent to pure ChPT result
- These new calculations helped to sharpen the "Cabibbo angle anomaly", which points towards new physics.

Thanks for your attention!