

Gauge Theory Bootstrap:

Pion amplitudes and low energy parameters

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based on [2309.12402](#) and [2403.10772](#) with [Martin Kruczenski](#)

Strongly coupled gauge theory

confinement

energy

asymptotic freedom

$$\mathcal{L} = i \sum_j^{N_f} \bar{q}_j \not{D} q_j - \sum_j^{N_f} m_q \bar{q}_j q_j - \frac{1}{4} G_a^{\mu\nu} G_{\mu\nu}^a$$

weakly coupled quarks/gluons (pQCD)

$$N_c \quad N_f \quad m_q \quad \Lambda_{\text{QCD}}$$

Strongly coupled gauge theory

confinement

energy

chiSB

asymptotic freedom

$$\mathcal{L} = \frac{f_\pi^2}{4} \left\{ \text{Tr} \left(\partial_\mu U \partial^\mu U^\dagger \right) + m_\pi^2 \text{Tr} \left(U + U^\dagger \right) \right\}$$

weakly coupled pions (EFT)

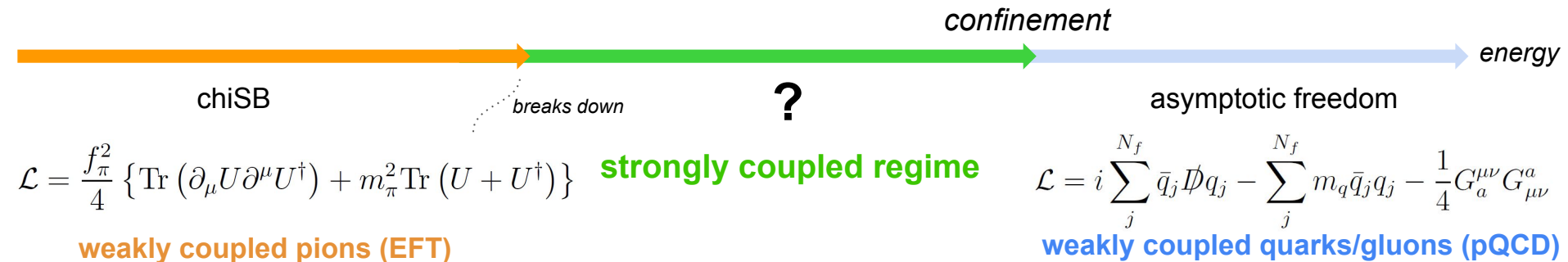
$$f_\pi \quad m_\pi$$

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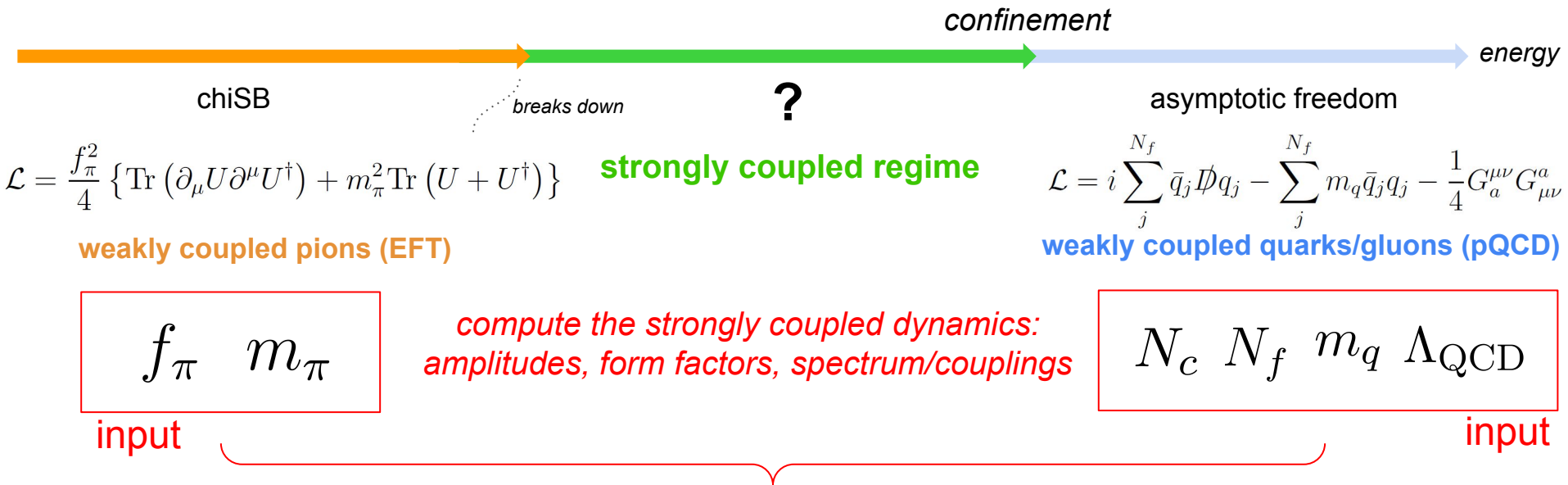
Strongly coupled gauge theory



$$f_\pi \quad m_\pi$$

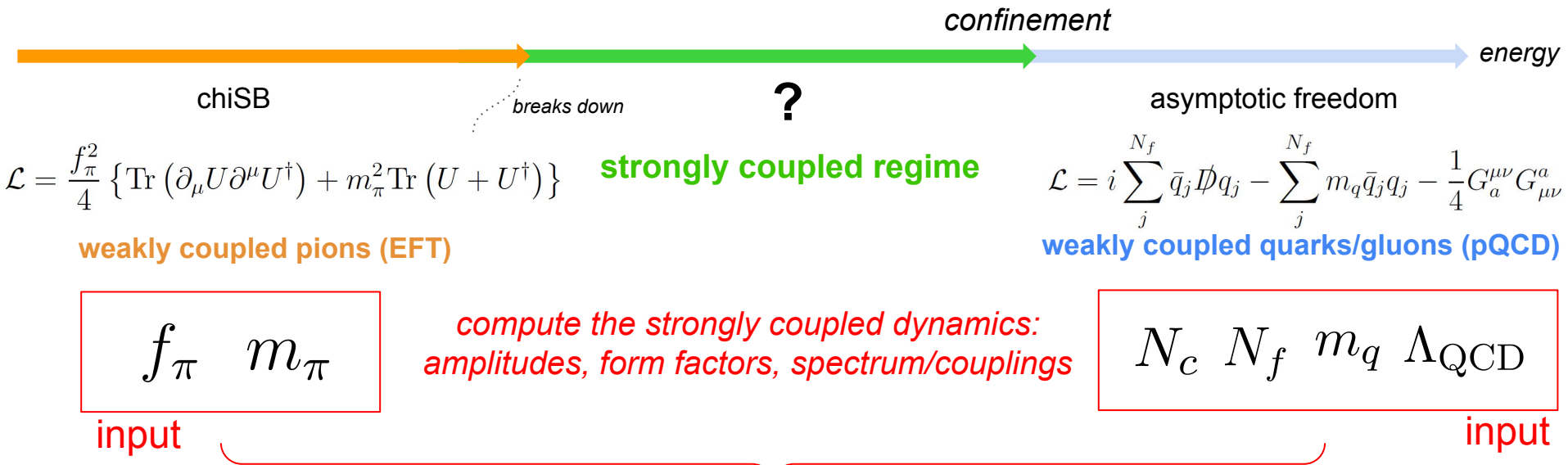
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Strongly coupled gauge theory \rightarrow Gauge Theory Bootstrap



Gauge Theory Bootstrap

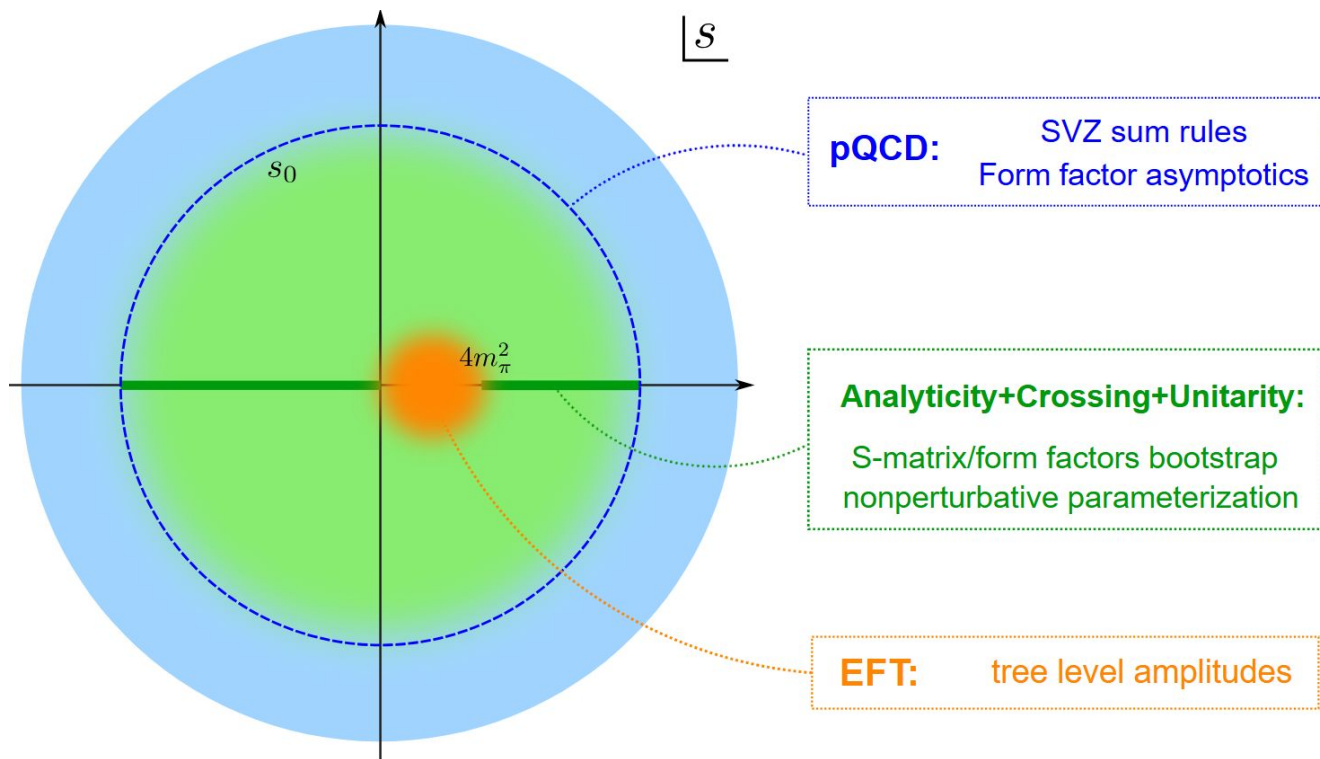
Strongly coupled gauge theory \rightarrow Gauge Theory Bootstrap



Gauge Theory Bootstrap

theoretical/numerical computation, not using experimental scattering data as input

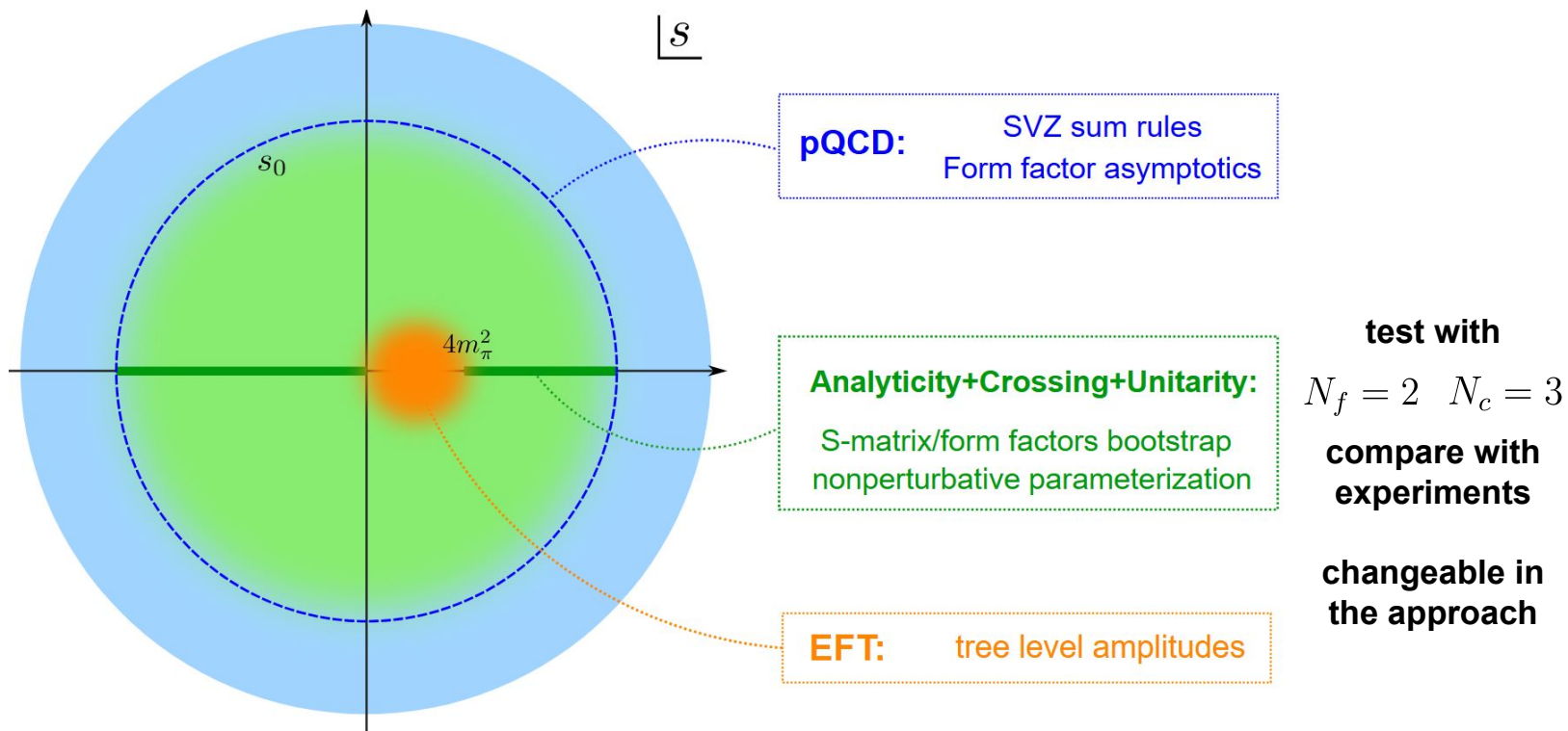
Gauge Theory Bootstrap: summary



look for amplitudes/form factors that:

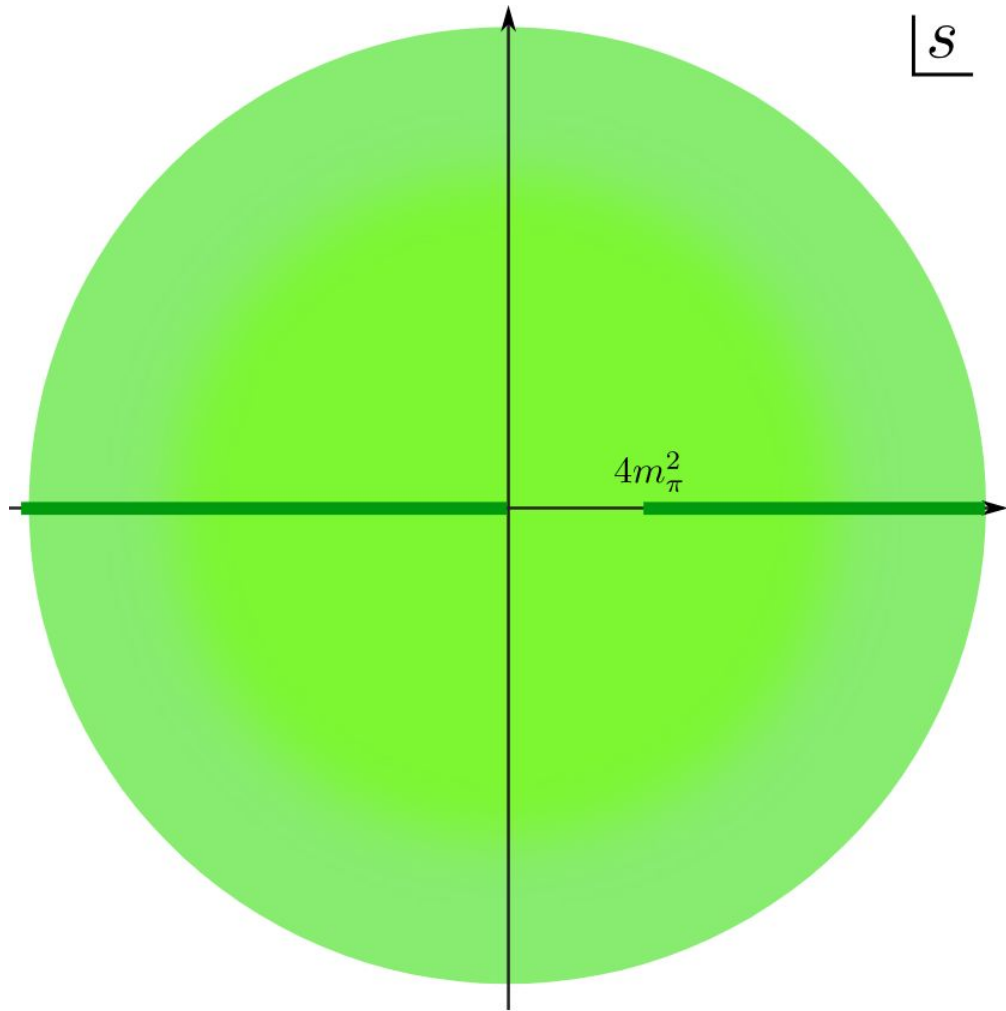
- 1, satisfy generic consistency conditions (analyticity, crossing, unitarity)
- 2, match low energy behavior (chiSB) and high energy (pQCD)

Gauge Theory Bootstrap: summary



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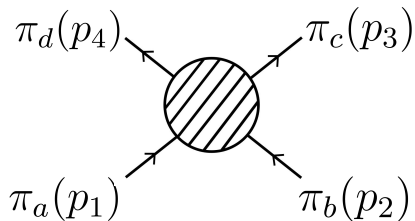
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Analyticity+Crossing+Unitarity:

S-matrix bootstrap
nonperturbative parameterization

S-matrix bootstrap parameterization



modern S-matrix bootstrap: [Paulos, Penedones, Toledo, van Rees, Vieira, 2016&2017]

$$\langle p_1, a; p_2, b | \mathbf{T} | p_3, c; p_4, d \rangle = A(s, t, u) \delta_{ab} \delta_{cd} + A(t, s, u) \delta_{ac} \delta_{bd} + A(u, t, s) \delta_{ad} \delta_{bc}$$

Crossing

$$A(s, t, u) = A(s, u, t)$$

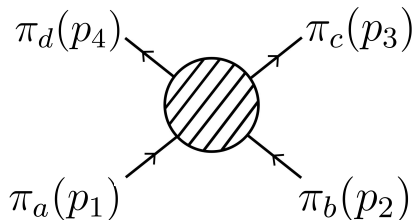
Analyticity

$$\text{cuts } s, t, u > 4$$

$$m_\pi = 1$$

$$s + t + u = 4$$

S-matrix bootstrap parameterization



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Crossing

$$A(s, t, u) = A(s, u, t)$$

Analyticity

cuts $s, t, u > 4$

$$m_\pi = 1$$

$$s + t + u = 4$$

nonperturbative parameterization encoding **Analyticity** and **Crossing**:

$$A(s, t, u) = \frac{1}{\pi^2} \int_4^\infty dx \int_4^\infty dy \left[\frac{\rho_1(x, y)}{(x-s)(y-t)} + \frac{\rho_1(x, y)}{(x-s)(y-u)} + \frac{\rho_2(x, y)}{(x-t)(y-u)} \right] + \text{subtraction terms}$$

parameters: $\{\rho_{\alpha=1,2}(x, y), \dots\}$

numerics: discretize

$\{\rho_{\alpha,ij}, \dots\}$

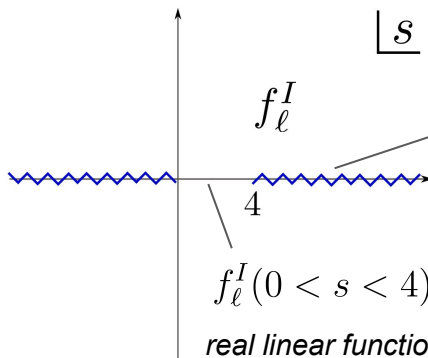
bootstrap variables

S-matrix bootstrap parameterization

$SU(2)_V$ isospin **Symmetry**

$$f_\ell^I(s) = \frac{1}{4} \int_{-1}^1 d\cos\theta P_\ell(\cos\theta) T^I(s, t)$$

analytic function of s



$$S_\ell^I(s) = 1 + i\pi \sqrt{\frac{s-4}{s}} f_\ell^I(s) \quad s > 4 \quad \forall \ell, I$$

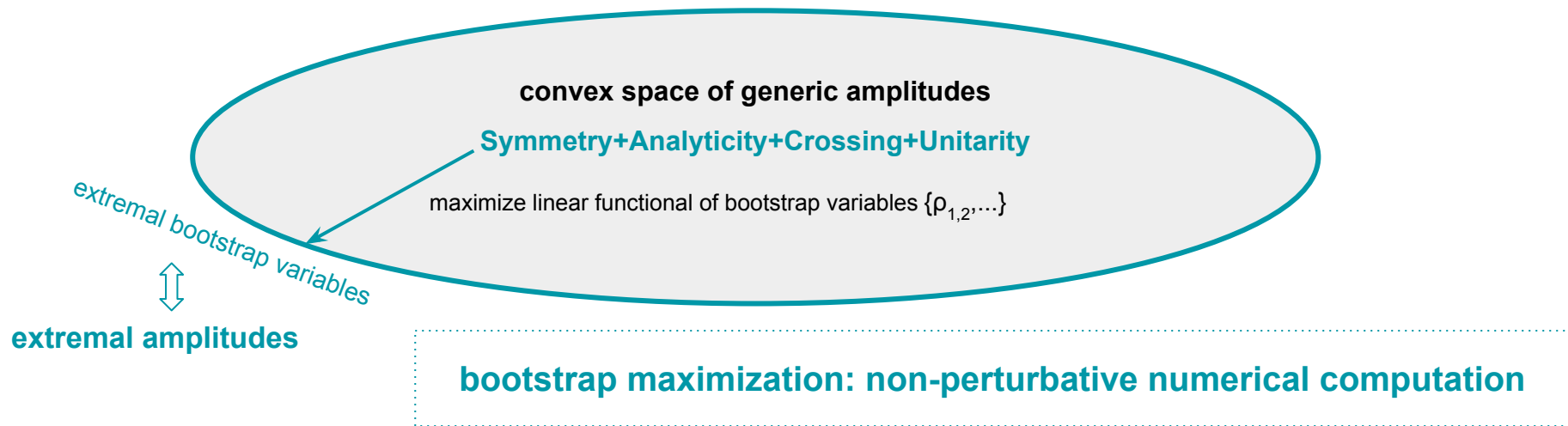
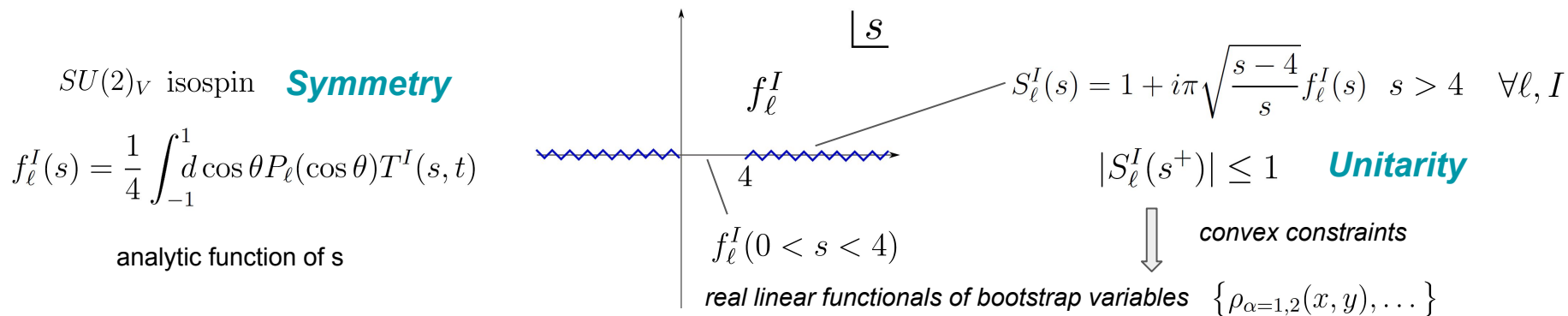
$$|S_\ell^I(s^+)| \leq 1 \quad \textbf{Unitarity}$$

convex constraints



real linear functionals of bootstrap variables $\{\rho_{\alpha=1,2}(x, y), \dots\}$

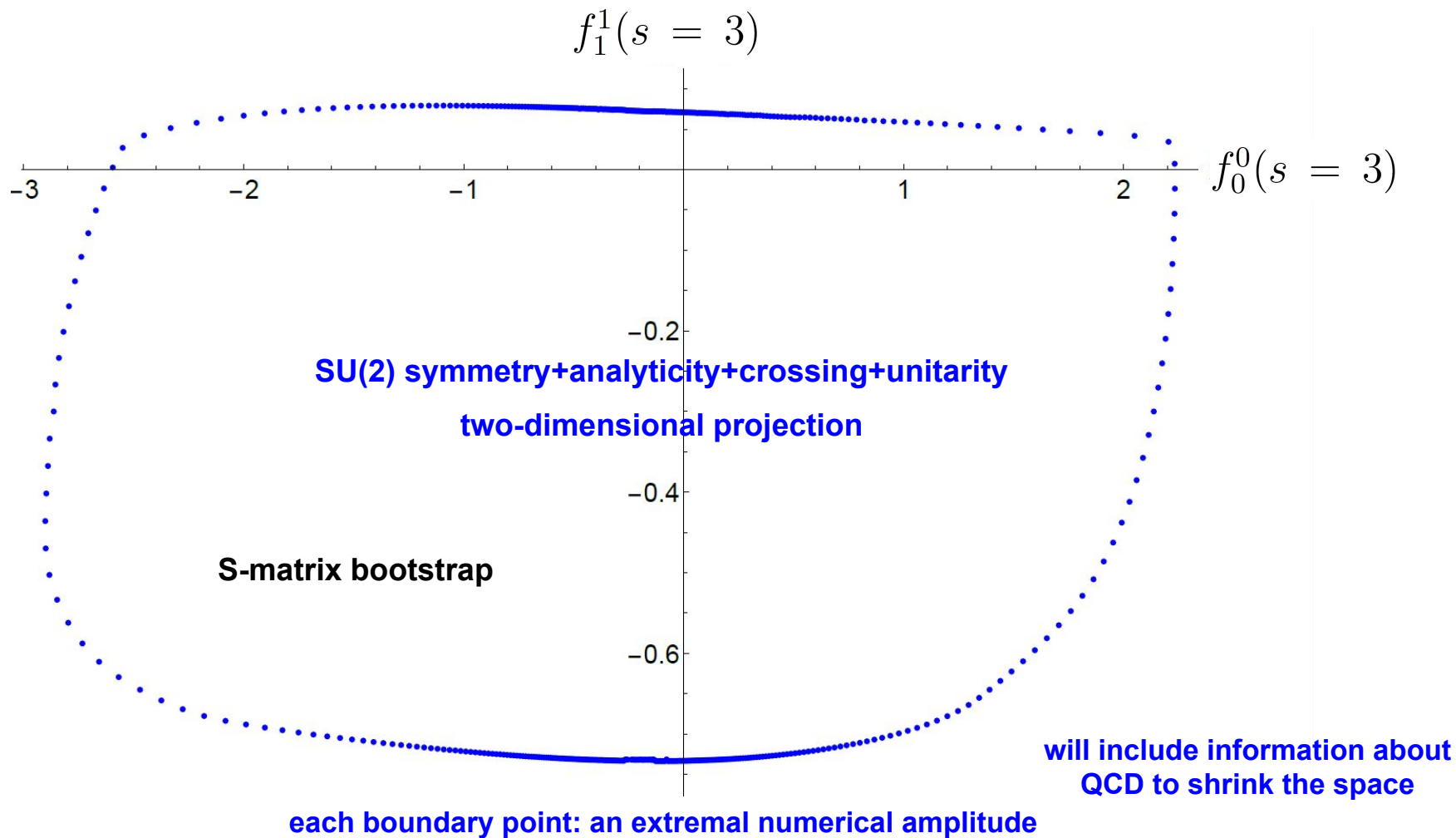
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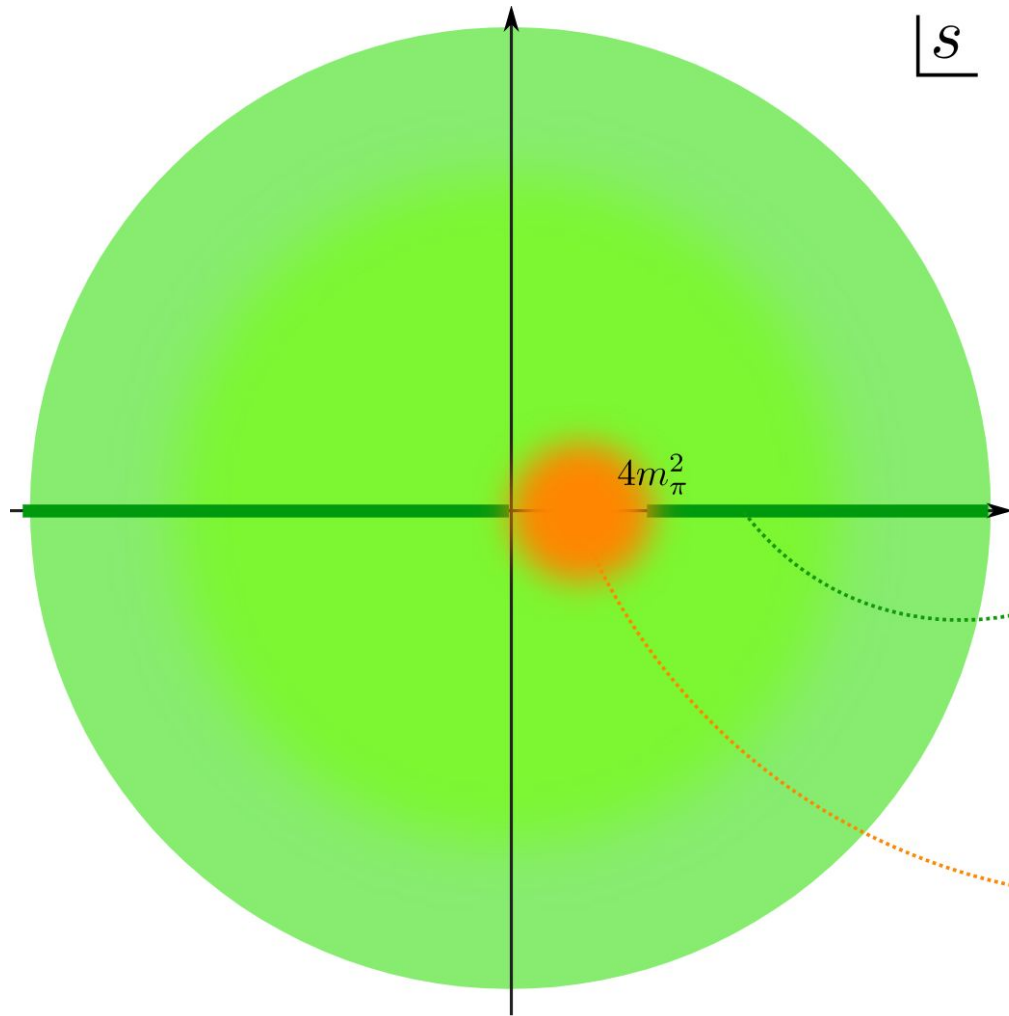


$$f_1^1(s = 3)$$

$$f_0^0(s = 3)$$

**plot a two-dimensional projection
of the space of amplitudes under
SU(2) symmetry, analyticity, crossing, unitarity**





Analyticity+Crossing+Unitarity:

S-matrix bootstrap
nonperturbative parameterization

EFT: tree level amplitudes

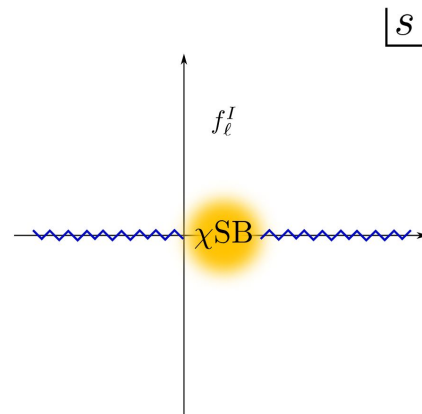
Low energy matching with EFT

interaction: $\mathcal{L}_2^{4\pi} = \frac{1}{6f_\pi^2} \left((\vec{\pi} \cdot \partial_\mu \vec{\pi})^2 - \vec{\pi}^2 (\partial_\mu \vec{\pi} \cdot \partial^\mu \vec{\pi}) \right) + \frac{m_\pi^2}{24f_\pi^2} (\vec{\pi}^2)^2$

tree-level amplitude: $A_{\text{tree}}(s, t, u) = \frac{4}{\pi} \frac{s - m_\pi^2}{32\pi f_\pi^2}$ [Weinberg, 1966]

S0: $f_{0,\text{tree}}^0(s) = \frac{2}{\pi} \frac{2s - m_\pi^2}{32\pi f_\pi^2}$ P1: $f_{1,\text{tree}}^1(s) = \frac{2}{\pi} \frac{s - 4m_\pi^2}{96\pi f_\pi^2}$ S2: $f_{0,\text{tree}}^2(s) = \frac{2}{\pi} \frac{2m_\pi^2 - s}{32\pi f_\pi^2}$

good in the unphysical region (very low energy) $0 < s < 4m_\pi^2$



Low energy matching with EFT

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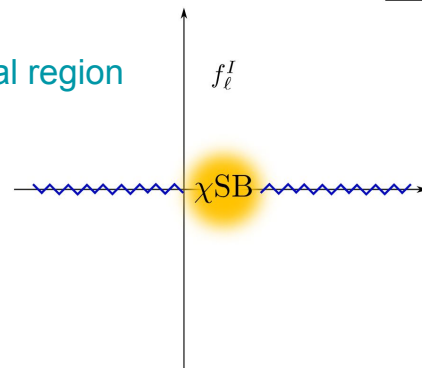
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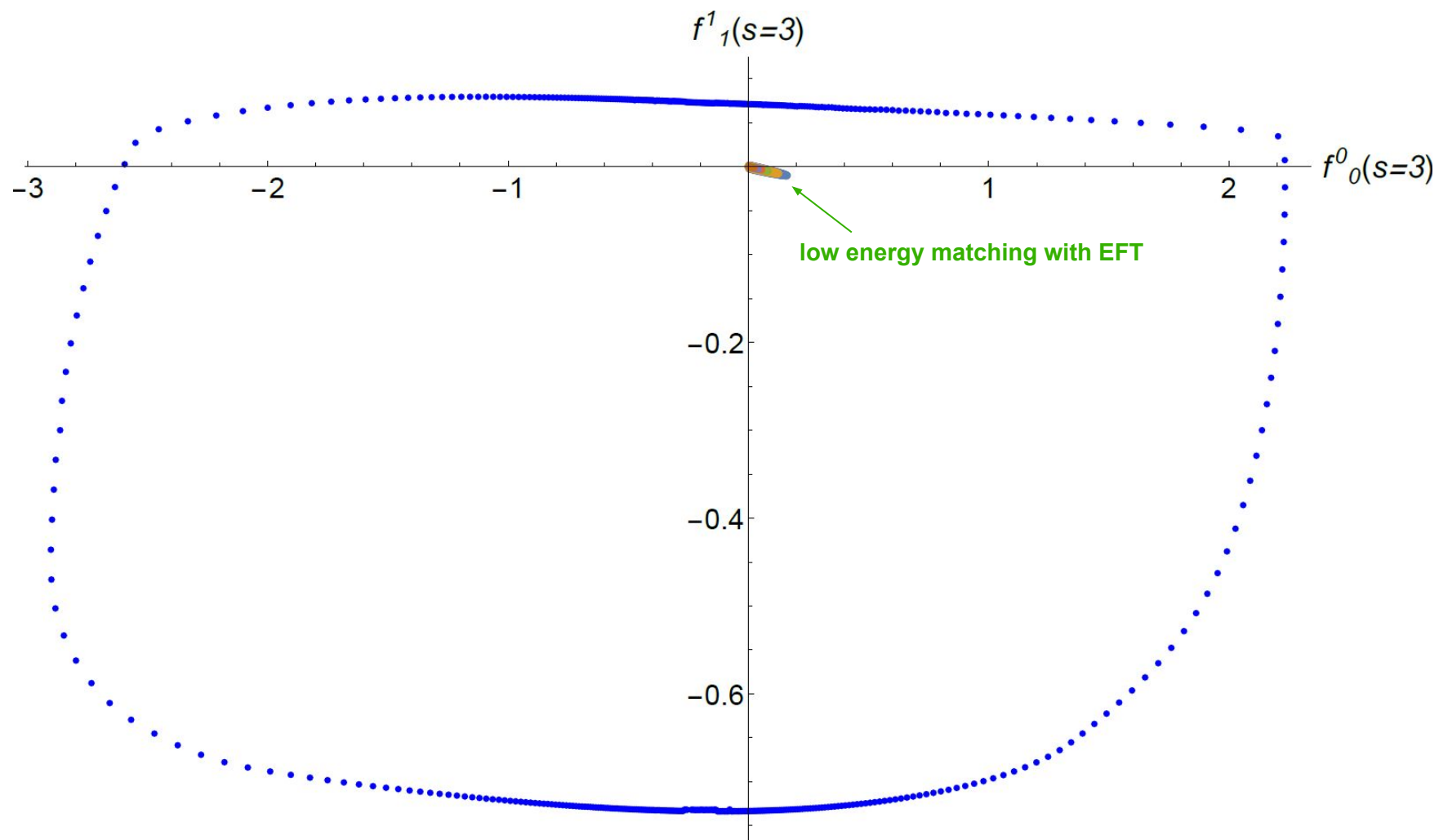
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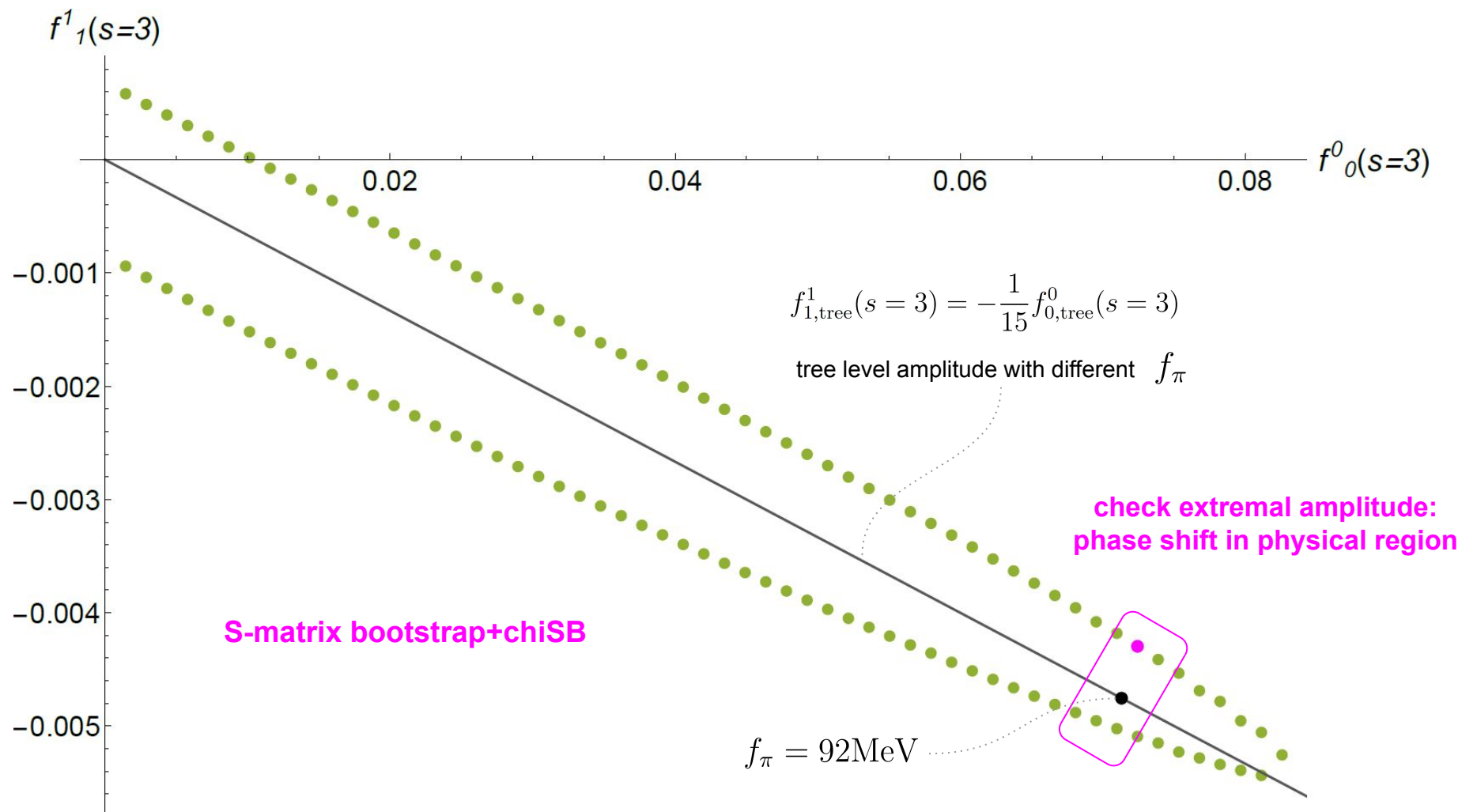
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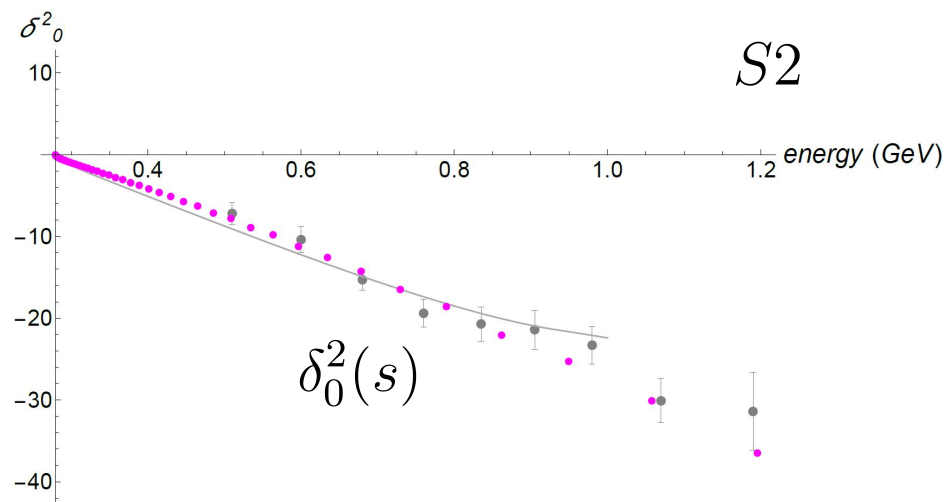
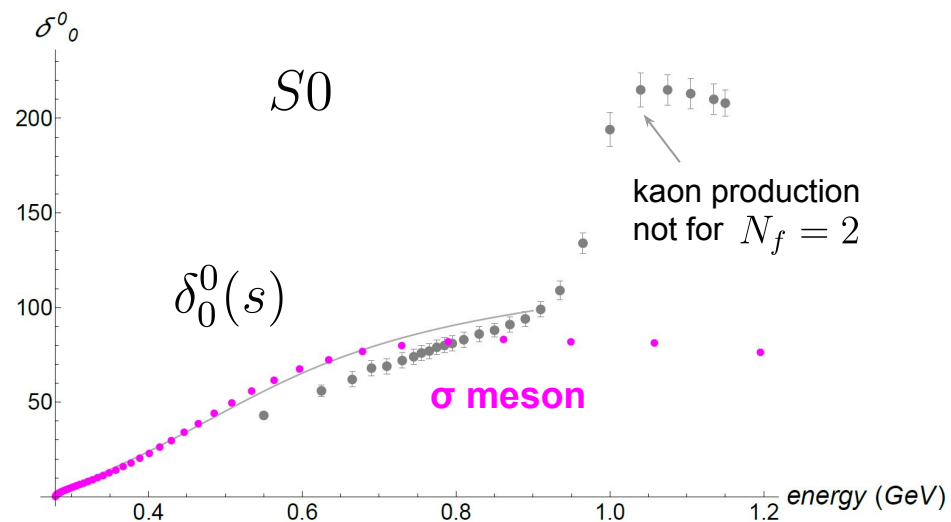
numerically: requires pw in bootstrap match tree level pw in the very low energy unphysical region

$$f_0^0(s) \simeq f_{0,\text{tree}}^0(s) \quad f_1^1(s) \simeq f_{1,\text{tree}}^1(s) \quad f_0^2(s) \simeq f_{0,\text{tree}}^2(s) \quad 0 < s < 4m_\pi^2$$



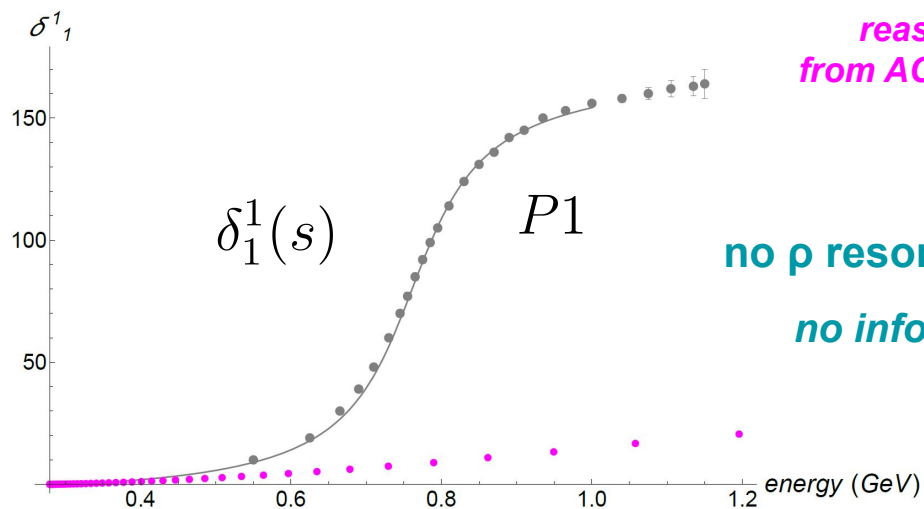






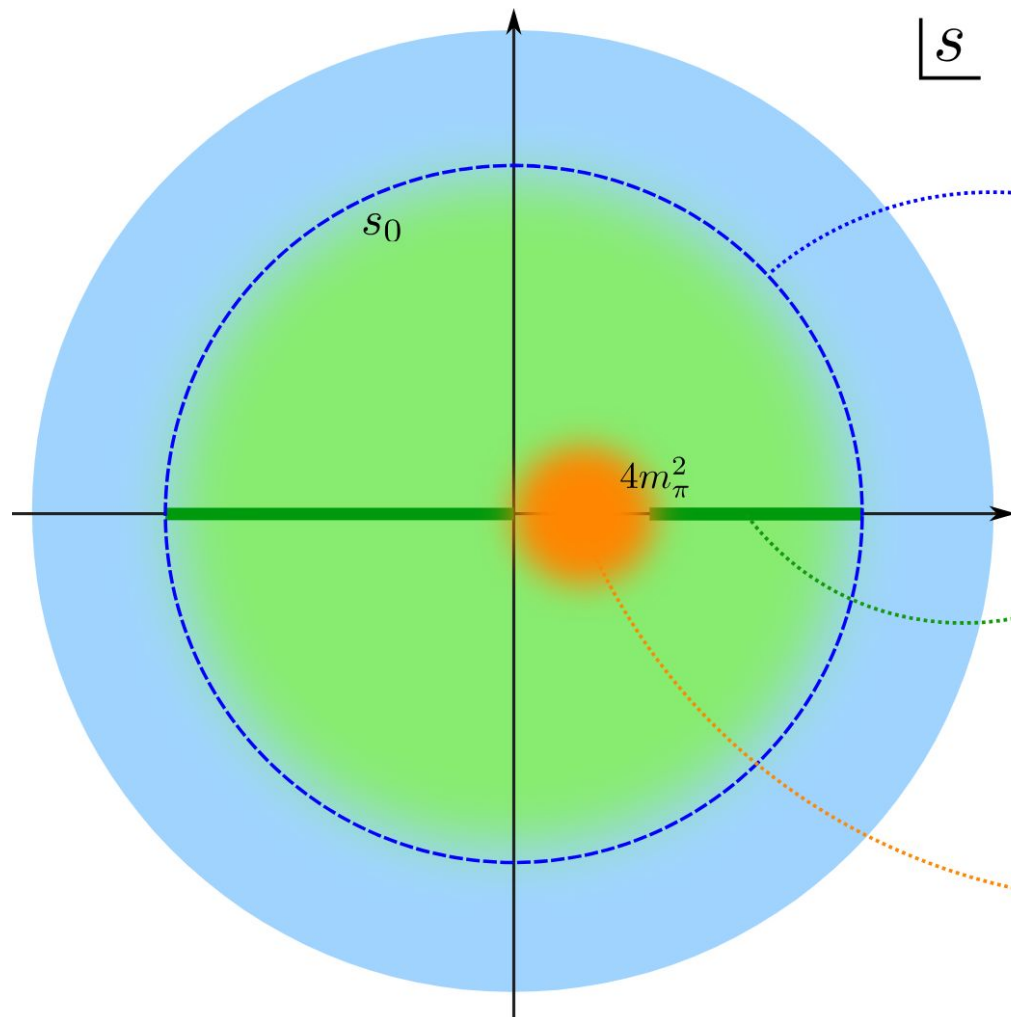
bootstrap numerical
computation of
phase shifts

experimental data (gray dots)
[Protopopescu et al, 1973]
[Losty et al, 1974]
pheno fit (gray line)
[Pelaez, Yndurain, 2005]



reasonable $S0$, $S2$ waves
from ACU + weakly coupled EFT

no ρ resonance without UV info
no information on QCD yet



s

pQCD: SVZ sum rules
Form factor asymptotics

Analyticity+Crossing+Unitarity:

S-matrix/form factors bootstrap
nonperturbative parameterization

EFT: tree level amplitudes

S-matrix/form factor bootstrap

[Karateev, Kuhn, Penedones, 2019]

$$|\psi_1\rangle = |p_1, p_2\rangle_{in}, \quad |\psi_2\rangle = |p_1, p_2\rangle_{out}, \quad |\psi_3\rangle = \int dx e^{-i(p_1+p_2)\cdot x} \mathcal{O}(x)|0\rangle$$

positive semidefinite matrix

$$\langle \psi_a | \psi_b \rangle = \begin{pmatrix} 1 & S & \mathcal{F} \\ S^* & 1 & \mathcal{F}^* \\ \mathcal{F}^* & \mathcal{F} & \rho \end{pmatrix} \succeq 0$$

S-matrix/form factor bootstrap

[Karateev, Kuhn, Penedones, 2019]

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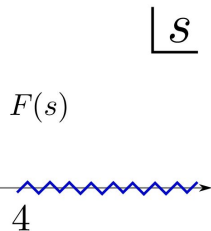
2-particle form factor:

$${}_{out} \langle p_1, p_2 | \mathcal{O}(0) | 0 \rangle = F(s)$$

spectral density: $\int \frac{d^4x}{(2\pi)^4} e^{iPx} \langle 0 | \mathcal{O}^\dagger(x) \mathcal{O}(0) | 0 \rangle = \rho(s)$
supported at $s > 4$

analytic function of s

$$F(s) = \frac{1}{\pi} \int_4^\infty dx \frac{\text{Im} F(x)}{x - s} + \text{subtractions}$$



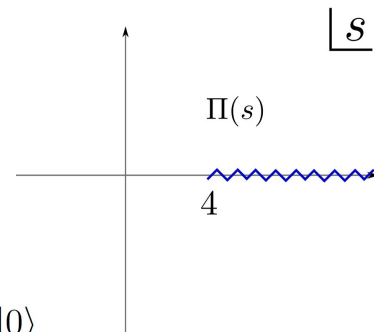
extended bootstrap variables: $\{\rho_{1,2}(x, y), \dots, \text{Im} F(x), \rho(x)\}$

Gauge theory current correlators

$$\begin{matrix} \langle \text{in} |_{P', I, \ell} & \langle \text{out} |_{P, I, \ell} & \mathcal{O}_{P, I, \ell} | 0 \rangle \\ \langle \text{out} |_{P', I, \ell} & \left(\begin{array}{ccc} 1 & S_\ell^I(s) & \mathcal{F}_\ell^I \\ S_\ell^{I*}(s) & 1 & \mathcal{F}_\ell^{I*} \\ \mathcal{F}_\ell^{I*} & \mathcal{F}_\ell^I & \rho_\ell^I(s) \end{array} \right) & \succeq 0 \quad s > 4 \quad \forall \ell, I \\ \langle 0 | \mathcal{O}_{P', I, \ell}^\dagger & & \end{matrix}$$

construct operators from gauge theory with desired quantum numbers

$$\rho_\ell^I(s) = 2 \text{Im} \Pi_\ell^I(x + i\epsilon)$$



e.g. vector (electromagnetic) current

$$\Pi_1^1(s) = i \int \frac{d^4x}{(2\pi)^4} e^{iPx} \langle 0 | \hat{T} \{ j_\sigma^\dagger(x) j_\sigma(0) \} | 0 \rangle$$

$$P1 : j_V^\mu(x) = \frac{1}{2} (\bar{u} \gamma^\mu u - \bar{d} \gamma^\mu d)$$

Gauge theory current correlators & SVZ expansion

$$\begin{matrix} \langle \text{in} |_{P', I, \ell} \\ \langle \text{out} |_{P', I, \ell} \\ \langle 0 | \mathcal{O}_{P', I, \ell}^\dagger \end{matrix} \begin{pmatrix} 1 & S_\ell^I(s) & \mathcal{F}_\ell^I \\ S_\ell^{I*}(s) & 1 & \mathcal{F}_\ell^{I*} \\ \mathcal{F}_\ell^{I*} & \mathcal{F}_\ell^I & \rho_\ell^I(s) \end{pmatrix} \begin{matrix} | \text{in} \rangle_{P, I, \ell} \\ | \text{out} \rangle_{P, I, \ell} \\ \mathcal{O}_{P, I, \ell} | 0 \rangle \end{matrix} \geq 0 \quad s > 4 \quad \forall \ell, I$$

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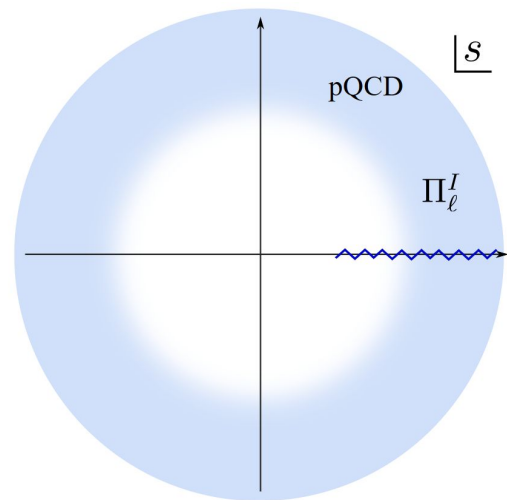
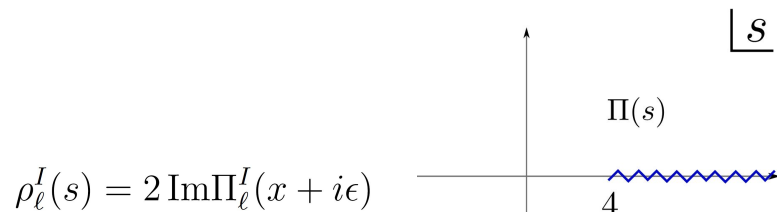
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large spacelike momenta — asymptotic free region with pQCD computation

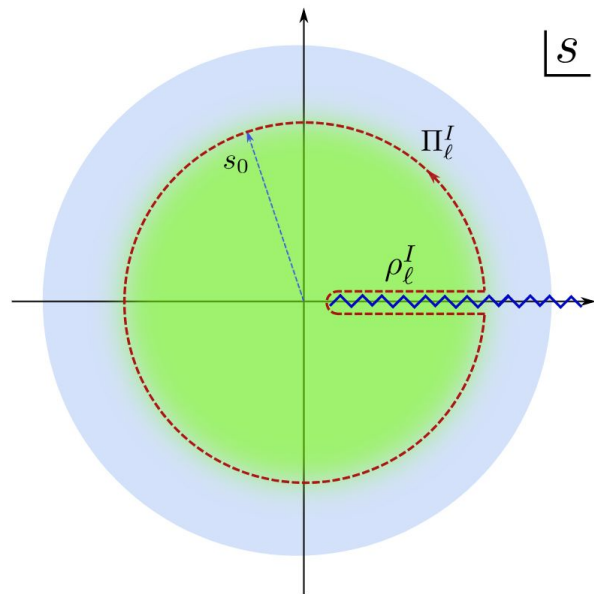
$$\Pi_1^1(s) = \frac{1}{2} \frac{1}{(2\pi)^4} \left\{ -\frac{1}{4\pi^2} \left(1 + \frac{\alpha_s}{\pi} \right) s \ln\left(-\frac{s}{\mu^2}\right) + \dots \right\}$$

pQCD computation

[Shifman, Vainshtein, Zakharov, 1979]



Finite energy sum rule



connect pQCD with bootstrap at s_0

contour integral $s^n \Pi(s)$ vanishes

$$\int_4^{s_0} \rho(x) x^n dx = -s_0^{n+1} \int_0^{2\pi} e^{i(n+1)\varphi} \Pi(s_0 e^{i\varphi}) d\varphi$$

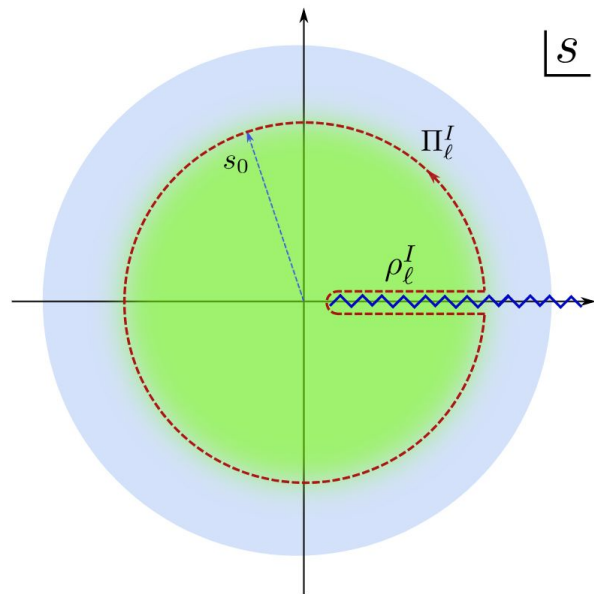
bootstrap variables

gauge theory information

linear constraints

SVZ

Finite energy sum rule



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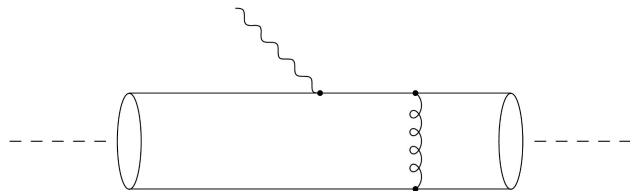
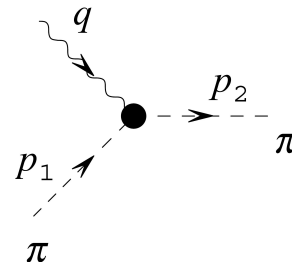
$$P1 : \frac{1}{s_0^{n+2}} \int_4^{s_0} \rho_1^1(x) x^n dx = \frac{1}{2(2\pi)^4} \left\{ \frac{1}{2\pi(n+2)} \left(1 + \frac{\alpha_s}{\pi} \right) + \dots \right\}, \quad n \geq -1$$

Asymptotic behavior of form factor from pQCD

perturbative QCD also controls asymptotic behavior of form factors

e.g. electromagnetic FF $\langle \pi(p_2) | J_{\text{em}}^\mu(0) | \pi(p_1) \rangle = (p_1^\mu + p_2^\mu) F_\pi(q^2)$

$$q = (p_2 - p_1)$$



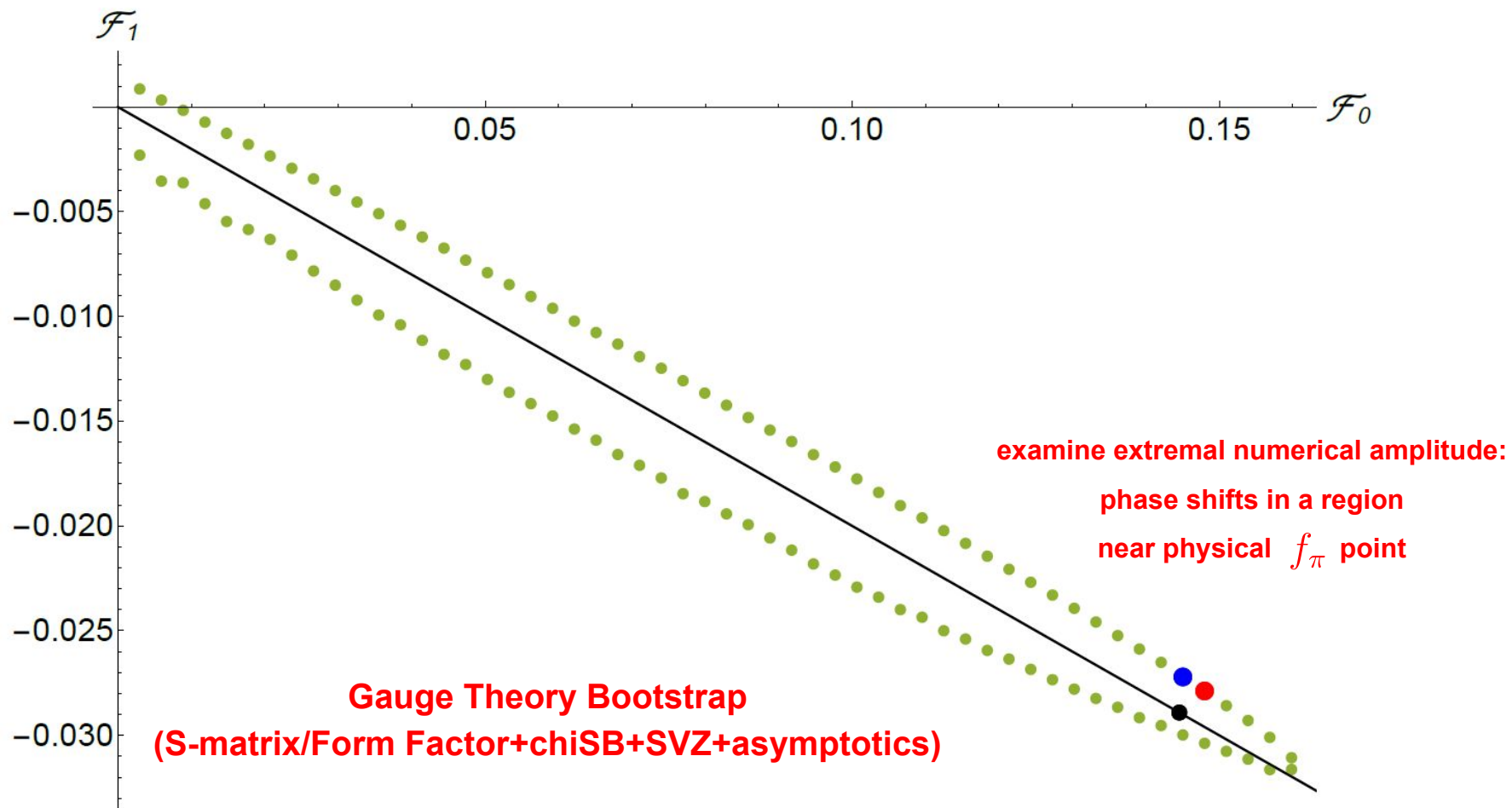
$$F_\pi(s) \simeq -\frac{16\pi\alpha_s(s)f_\pi^2}{s}$$

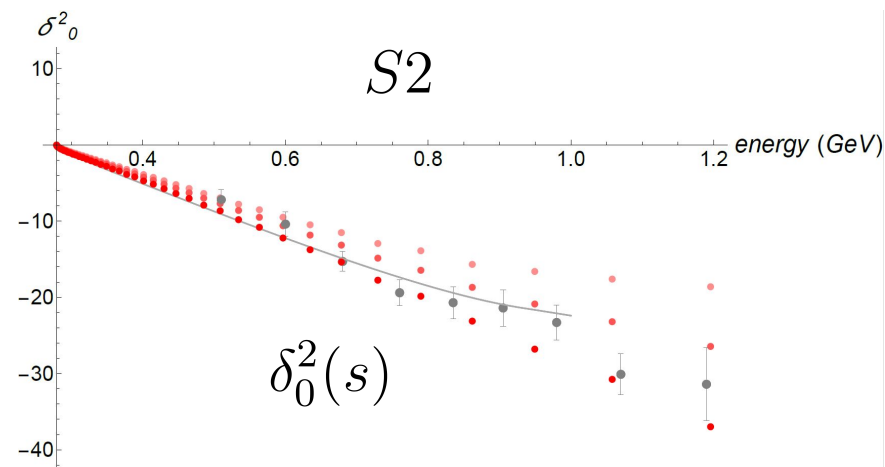
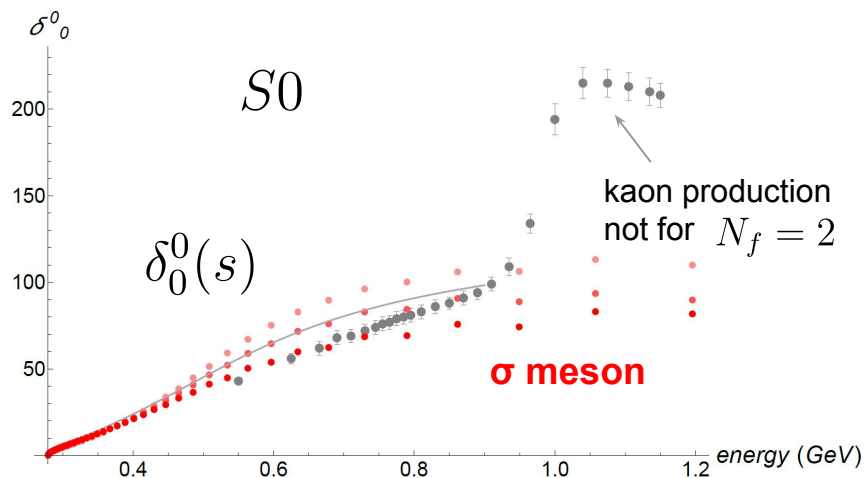
[Lepage, Brodsky, 1979]

evaluate to estimate

in practical numerical implementation
suffices to require smallness above $s = s_0$

$$||\mathcal{F}(s > s_0)|| \lesssim \epsilon$$





Gauge Theory Bootstrap

phase shifts up to 1.2 GeV

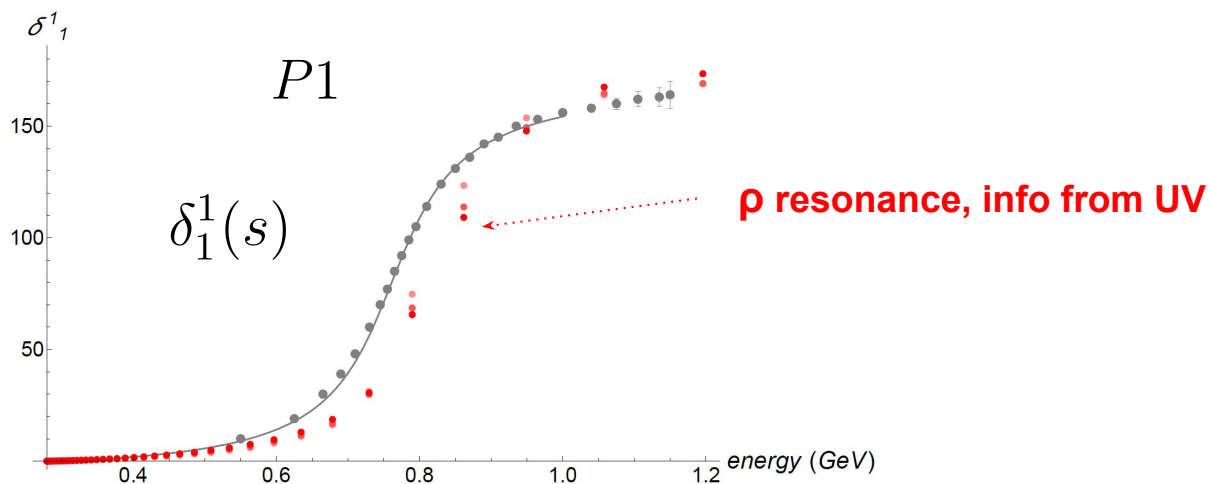
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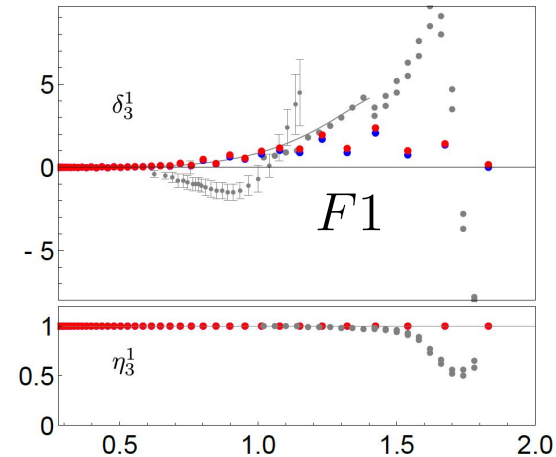
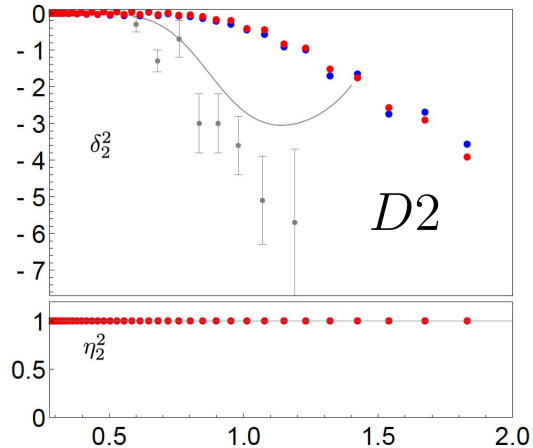
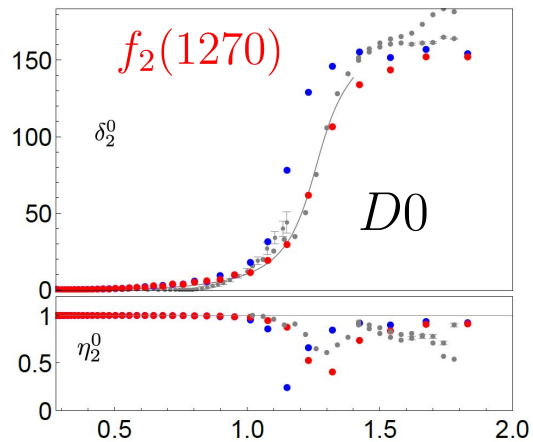
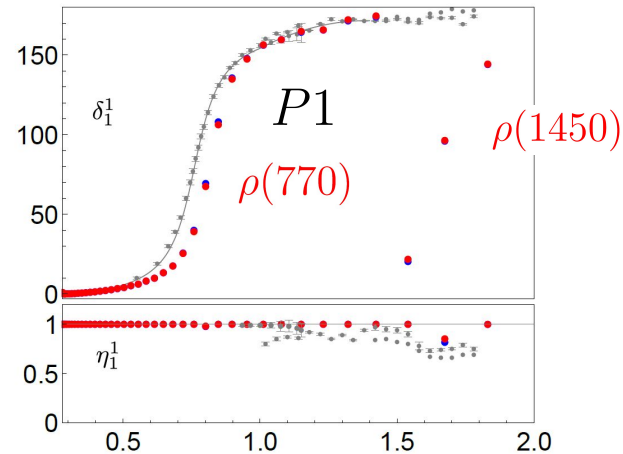
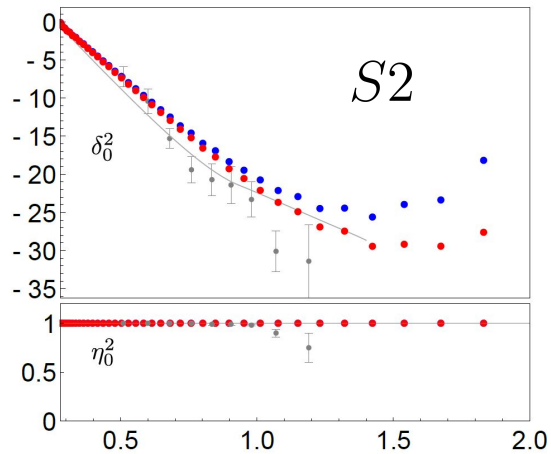
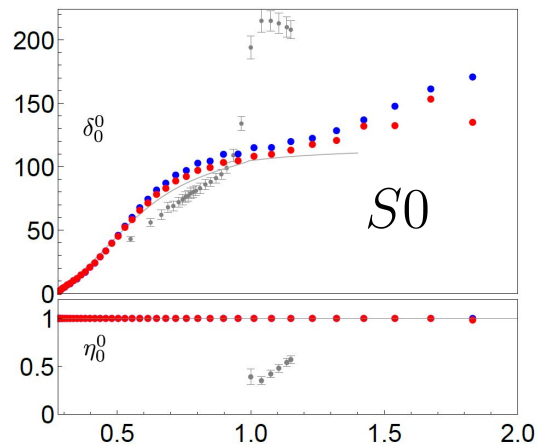
pheno fit (gray line)

[Pelaez, Yndurain, 2005]



Gauge Theory Bootstrap

phase shifts up to 2 GeV



Low energy parameters: threshold expansions

scattering lengths and effective range parameters

$$\text{Ref}_\ell^I(s) \stackrel{k \rightarrow 0}{\simeq} \frac{2m_\pi}{\pi} k^{2\ell} (a_\ell^I + b_\ell^I k^2 + \dots) \quad k = \frac{\sqrt{s - 4m_\pi^2}}{2}$$

	W	GTB	CGL	PY
$a_0^{(0)}$	0.16	0.178, 0.182	0.220 ± 0.005	0.230 ± 0.010
$a_0^{(2)}$	-0.046	-0.0369, -0.0378	-0.0444 ± 0.0010	-0.0422 ± 0.0022
$b_0^{(0)}$	0.18	0.287, 0.290	0.280 ± 0.001	0.268 ± 0.010
$b_0^{(2)}$	-0.092	-0.064, -0.066	-0.080 ± 0.001	-0.071 ± 0.004
$a_1^{(1)}$	31	28.0, 28.4	37.0 ± 0.13	$38.1 \pm 1.4 (\times 10^{-3})$
$b_1^{(1)}$	0	2.86, 3.37	5.67 ± 0.13	$4.75 \pm 0.16 (\times 10^{-3})$
$a_2^{(0)}$	0	12.6, 12.3	17.5 ± 0.3	$18.0 \pm 0.2 (\times 10^{-4})$
$a_2^{(2)}$	0	2.87, 2.81	1.70 ± 0.13	$2.2 \pm 0.2 (\times 10^{-4})$

pion charge radii

$$F_0^0(s) = F_0^0(0) \left[1 + \frac{1}{6} s \langle r^2 \rangle_S^\pi + \dots \right]$$

$$F_1^1(s) = 1 + \frac{1}{6} s \langle r^2 \rangle_V^\pi + \dots$$

	GTB	Exp. fits
$\langle r^2 \rangle_S^\pi$	0.64, 0.61	$0.61 \pm 0.04 \text{ fm}^2$
$\langle r^2 \rangle_V^\pi$	0.388, 0.381	$0.439 \pm 0.008 \text{ fm}^2$

Low energy parameters: chiral Lagrangian coefficients

calculate the chiral
Lagrangian coefficients

$$a_{D0} = \frac{1}{1440\pi^3 f_\pi^4} \left\{ \bar{l}_1 + 4\bar{l}_2 - \frac{53}{8} \right\} + \dots$$

$$a_{D2} = \frac{1}{1440\pi^3 f_\pi^4} \left\{ \bar{l}_1 + \bar{l}_2 - \frac{103}{40} \right\} + \dots$$

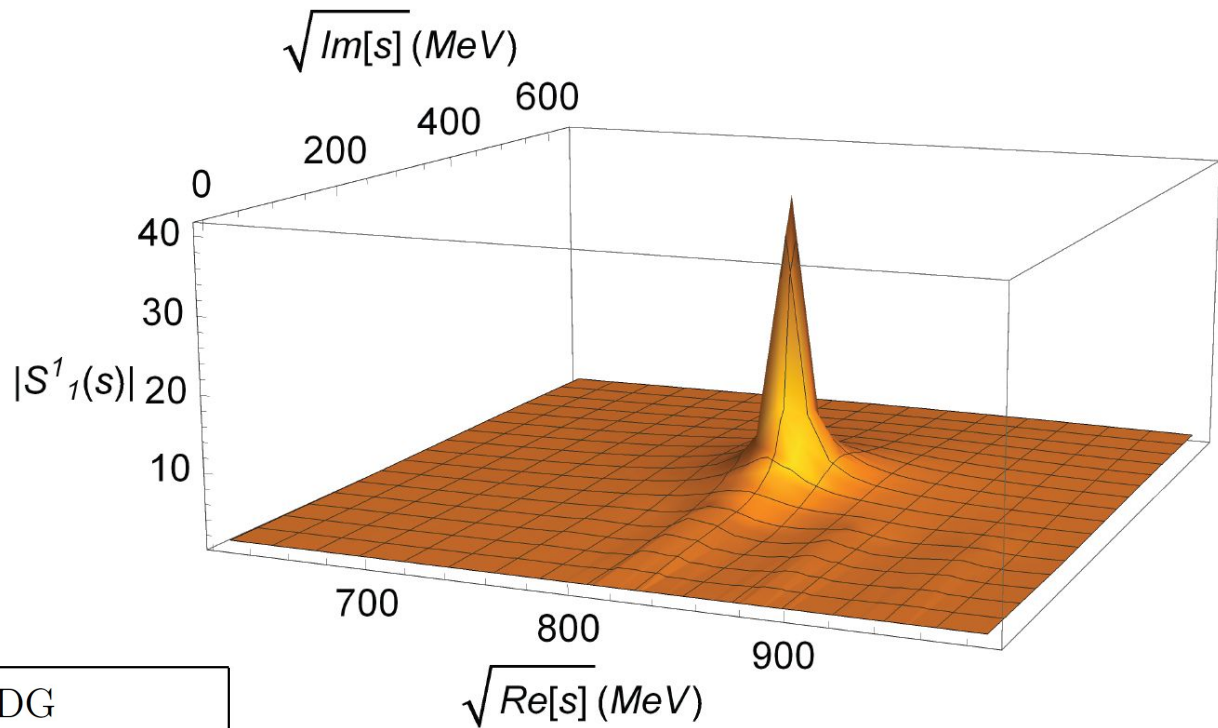
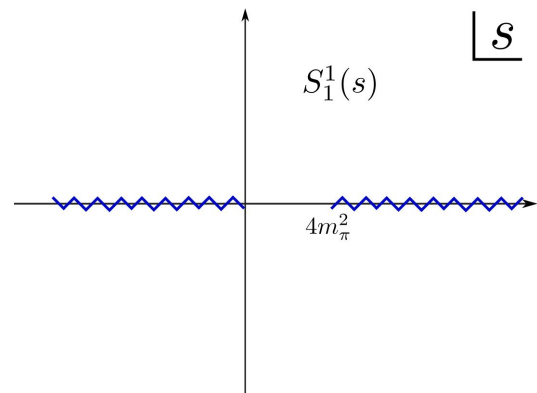
[Gasser Leutwyler, 1984]

$$F_0(s) = 1 + \frac{s}{16\pi^2 f_\pi^2} \left(\bar{l}_4 - \frac{13}{12} \right) + \dots$$

$$F_1(s) = 1 + \frac{s}{96\pi^2 f_\pi^2} (\bar{l}_6 - 1) + \dots$$

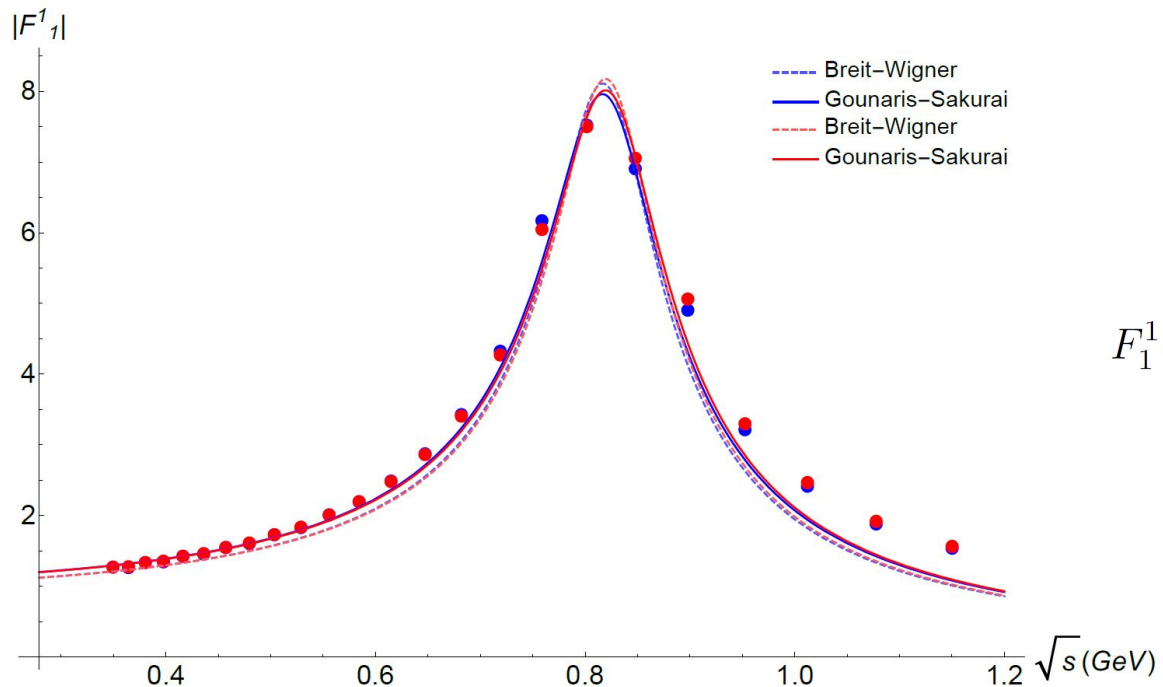
	GTB	GL	Bij	CGL
\bar{l}_1	0.92, 0.93	-2.3 ± 3.7	-1.7 ± 1.0	-0.4 ± 0.6
\bar{l}_2	4.1, 4.0	6.0 ± 1.3	6.1 ± 0.5	4.3 ± 0.1
\bar{l}_4	4.7, 4.6	4.3 ± 0.9	4.4 ± 0.3	4.4 ± 0.2
\bar{l}_6	14.3, 14.1	16.5 ± 1.1	$16.0 \pm 0.5 \pm 0.7$	

*$\rho(770)$ meson as pole on
the second sheet of $S_1^1(s)$*



	GTB	PDG
$\text{Re}(\sqrt{s_\rho})$	829, 832	$761 - 765 \pm 0.23$ MeV
$\text{Im}(\sqrt{s_\rho})$	63, 64	$71 - 74 \pm 0.8$ MeV

Vector (electromagnetic) form factor and $\rho(770)$ meson



Breit-Wigner form

$$F_1^1(s) = -\frac{m_\rho^2}{s - m_\rho^2 + im_\rho\Gamma_\rho\theta(s - 4m_\pi^2)}$$

	GTB	PDG
m_ρ	836, 839	775 ± 0.23 MeV
Γ_ρ	111, 111	149.1 ± 0.8 MeV

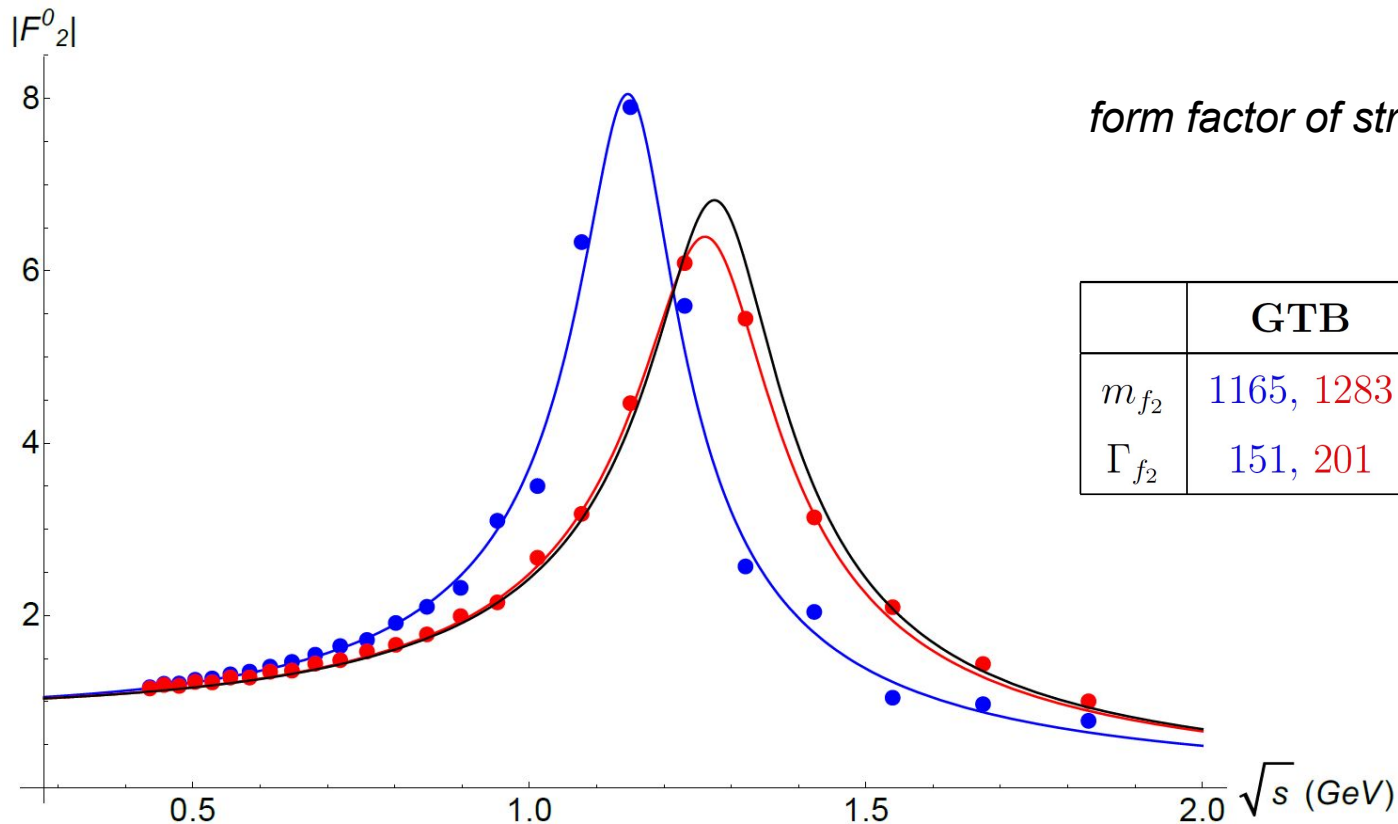
couplings

$$\Gamma_\rho = g_{\rho\pi\pi}^2 \frac{m_\rho}{48\pi} \left[1 - \frac{4m_\pi^2}{m_\rho^2} \right]^{\frac{3}{2}}$$

$$g_{\rho\pi\pi} = 4.9, \quad 4.9$$

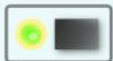
$$g_{\rho\pi\pi} = 6$$

Gravitational form factor and f_2 meson

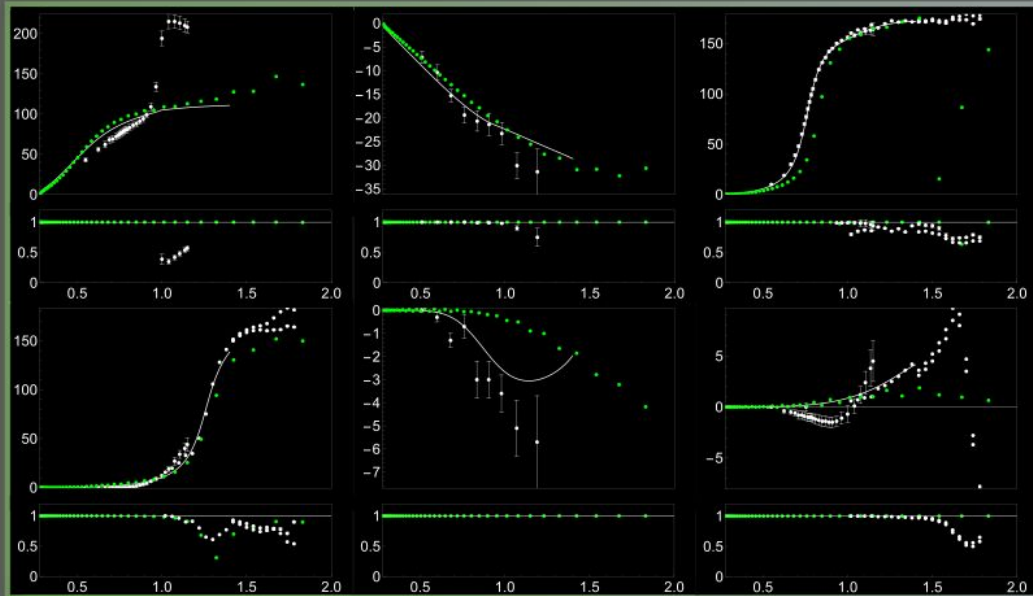


form factor of stress energy tensor

	GTB	PDG
m_{f_2}	1165, 1283	1275.4 ± 0.6 MeV
Γ_{f_2}	151, 201	186.6 ± 2.3 MeV



Gauge Theory Bootstrap



pQCD



N_f



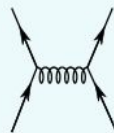
N_c



α_s



m_q



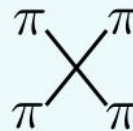
chiSB



m_π



f_π



- theoretical/numerical computation
- good track for solving QCD (gauge theories)
- improvements to make: more precise and robust
- fast: 20min on average laptop

Ancillary files ([details](#)):

- [GTB_numerics.m](#)
- [GTB_numerics.nb](#)

Thank you !